num=[0 1100 0]; den=[1 1100 10^5]; w=logspace(0,5,100); bode(num,den,w); grid
Students and working professionals will find *Circuit Analysis II with MATLAB® Applications* to be a concise and easy-to-learn text. It provides complete, clear, and detailed explanations of advanced electrical engineering concepts illustrated with numerous practical examples.

This text includes the following chapters and appendices:
- Second Order Circuits
- Resonance
- Elementary Signals
- The Laplace Transformation
- The Inverse Laplace Transformation
- Circuit Analysis with Laplace Transforms
- Frequency Response and Bode Plots
- Self and Mutual Inductances - Transformers
- One and Two Port Networks
- Three Phase Systems
- Introduction to MATLAB
- Differential Equations
- Matrices and Determinants
- Constructing Semilog Plots with Microsoft Excel
- Scaling

Each chapter contains numerous practical applications supplemented with detailed instructions for using MATLAB to obtain quick and accurate answers.

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Preface

This text is written for use in a second course in circuit analysis. The reader of this book should have the traditional undergraduate knowledge of an introductory circuit analysis material such as *Circuit Analysis I with MATLAB® Applications* by this author. Another prerequisite would be knowledge of differential equations, and in most cases, engineering students at this level have taken all required mathematics courses. It encompasses a spectrum of subjects ranging from the most abstract to the most practical, and the material can be covered in one semester or two quarters. Appendix B serves as a review of differential equations with emphasis on engineering related topics and it is recommended for readers who may need a review of this subject.

There are several textbooks on the subject that have been used for years. The material of this book is not new, and this author claims no originality of its content. This book was written to fit the needs of the average student. Moreover, it is not restricted to computer oriented circuit analysis. While it is true that there is a great demand for electrical and computer engineers, especially in the internet field, the demand also exists for power engineers to work in electric utility companies, and facility engineers to work in the industrial areas.

Chapter 1 is an introduction to second order circuits and it is essentially a sequel to first order circuits that were discussed in the last chapter of *Circuit Analysis I with MATLAB® Applications*. Chapter 2 is devoted to resonance, and Chapter 3 presents practical methods of expressing signals in terms of the elementary functions, i.e., unit step, unit ramp, and unit impulse functions. Thus, any signal can be represented in the complex frequency domain using the Laplace transformation.

Chapters 4 and 5 are introductions to the unilateral Laplace transform and Inverse Laplace transform respectively, while Chapter 6 presents several examples of analyzing electric circuits using Laplace transformation methods. Chapter 7 begins with the frequency response concept and Bode magnitude and frequency plots. Chapter 8 is devoted to transformers with an introduction to self and mutual inductances. Chapter 9 is an introduction to one- and two-terminal devices and presents several practical examples. Chapter 10 is an introduction to three-phase circuits.

It is not necessary that the reader has previous knowledge of MATLAB®. The material of this text can be learned without MATLAB. However, this author highly recommends that the reader studies this material in conjunction with the inexpensive MATLAB Student Version package that is available at most college and university bookstores. Appendix A of this text provides a practical introduction to MATLAB. As shown on the front cover of this text the magnitude and phase plots can be easily obtained with a one line MATLAB code. Moreover, MATLAB will be invaluable in later studies such as the design of analog and digital filters.
As stated above, Appendix B is a review of differential equations. Appendix C is an introduction to matrices, Appendix D provides instructions on constructing semilog templates to be used with Bode plots, and Appendix E discusses scaling methods.

In addition to numerous real-world examples, this text contains several exercises at the end of each chapter. Detailed solutions of all exercises are provided at the end of each chapter. The rationale is to encourage the reader to solve all exercises and check his effort for correct solutions and appropriate steps in obtaining the correct solution. And since this text was written to serve as a self-study or supplementary textbook, it provides the reader with a resource to test his knowledge.

The author has accumulated many additional problems for homework assignment and these are available to those instructors who adopt this text either as primary or supplementary text, and prefer to assign problems without the solutions. He also has accumulated many sample exams.

The author is indebted to the class of the Spring semester of 2001 at San Jose State University, San Jose, California, for providing several of the examples and exercises of this text.

Like any other new book, this text may contain some grammar and typographical errors. Accordingly, all feedback for errors, advice, and comments will be most welcomed and greatly appreciated.

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Chapter 1

Second Order Circuits

This chapter discusses the natural, forced and total responses in circuits containing resistors, inductors and capacitors. These circuits are characterized by linear second-order differential equations whose solutions consist of the natural and the forced responses. We will consider both DC (constant) and AC (sinusoidal) excitations.

1.1 The Response of a Second Order Circuit

A circuit containing \( n \) energy storage devices (inductors and capacitors) is said to be an \( n \text{-th order circuit} \), and the differential equation describing the circuit is an \( n \text{-th order differential equation} \). For example, if a circuit contains an inductor and a capacitor, or two capacitors or two inductors, along with other devices such as resistors, it is said to be a second-order circuit and the differential equation that describes it is a second order differential equation. It is possible, however, to describe a circuit having two energy storage devices with a set of two first-order differential equations, a circuit which has three energy storage devices with a set of three first-order differential equations and so on. These are called \( \text{state equations}^* \) but these will not be discussed here.

The response is found from the differential equation describing the circuit, and its solution is obtained as follows:

1. We write the differential or integrodifferential (nodal or mesh) equation describing the circuit. We differentiate, if necessary, to eliminate the integral.

2. We obtain the forced (steady-state) response. Since the excitation in our work here will be either a constant (DC) or sinusoidal (AC) in nature, we expect the forced response to have the same form as the excitation. We evaluate the constants of the forced response by substitution of the assumed forced response into the differential equation and equate terms of the left side with the right side. Refer to Appendix B for the general expression of the forced response (particular solution).

3. We obtain the general form of the natural response by setting the right side of the differential equation equal to zero, in other words, solve the homogeneous differential equation using the characteristic equation.

4. Add the forced and natural responses to form the complete response.

5. We evaluate the constants of the complete response from the initial conditions.

---

* State variables and state equations are discussed in Signals and Systems with MATLAB Applications, ISBN 0-9709511-3-2 by this author.
1.2 The Series RLC Circuit with DC Excitation

Let us consider the series RLC circuit of Figure 1.1 where the initial conditions are \( i_L(0) = I_0 \), \( v_C(0) = V_0 \), and \( u_o(t) \) is the unit step function. * We want to find an expression for the current \( i(t) \) for \( t > 0 \).

![Series RLC Circuit](image)

Figure 1.1. Series RLC Circuit

For this circuit

\[
Ri + L\frac{di}{dt} + \frac{1}{C}\int_0^t i\,dt + V_0 = v_S \quad t > 0
\]  

(1.1)

and by differentiation

\[
R\frac{di}{dt} + L\frac{d^2i}{dt^2} + \frac{i}{C} = \frac{dv_S}{dt}, \quad t > 0
\]

To find the forced response, we must first specify the nature of the excitation \( v_S \), that is, DC or AC. If \( v_S \) is DC (\( v_S = \text{constant} \)), the right side of (1.1) will be zero and thus the forced response component \( i_f = 0 \). If \( v_S \) is AC (\( v_S = V\cos(\omega t + \varphi) \)), the right side of (1.1) will be another sinusoid and therefore \( i_f = I\cos(\omega t + \varphi) \). Since in this section we are concerned with DC excitations, the right side will be zero and thus the total response will be just the natural response.

The natural response is found from the homogeneous equation of (1.1), that is,

\[
R\frac{di}{dt} + L\frac{d^2i}{dt^2} + \frac{i}{C} = 0
\]  

(1.2)

The characteristic equation of (1.2) is

\[
Ls^2 + Rs + \frac{1}{C} = 0
\]

* The unit step function is discussed in detail in Chapter 3. For our present discussion it will suffice to state that \( u_o(t) = 0 \) for \( t < 0 \) and \( u_o(t) = 1 \) for \( t > 0 \).
or

\[ s^2 + \frac{R}{L} s + \frac{1}{LC} = 0 \]

from which

\[ s_1, s_2 = -\frac{R}{2L} \pm \frac{\sqrt{\frac{R^2}{4L^2} - \frac{1}{LC}}}{\sqrt{4L^2}} \]  

(1.3)

We will use the following notations:

\[
\begin{align*}
\alpha_S &= \frac{R}{2L} & \omega_0 &= \frac{1}{\sqrt{LC}} & \beta_S &= \sqrt{\alpha_S^2 - \omega_0^2} & \omega_{ns} &= \sqrt{\omega_0^2 - \alpha_S^2} \\
\text{\textit{\textalpha \ or Damping}} & & \text{\textit{Resonant}} & & \text{\textit{Beta}} & & \text{\textit{\text Damped Natural}} \\
\text{Coefficient} & & \text{Frequency} & & \text{Coefficient} & & \text{Frequency} \\
\end{align*}
\]

(1.4)

where the subscript \(s\) stands for series circuit. Then, we can express (1.3) as

\[ s_1, s_2 = -\alpha_S \pm \sqrt{\alpha_S^2 - \omega_0^2} = -\alpha_S \pm \beta_S \text{ if } \alpha_S^2 > \omega_0^2 \]  

(1.5)

or

\[ s_1, s_2 = -\alpha_S \pm \sqrt{\omega_0^2 - \alpha_S^2} = -\alpha_S \pm \omega_{ns} \text{ if } \omega_0^2 > \alpha_S^2 \]  

(1.6)

Case I: If \( \alpha_S^2 > \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and unequal. This results in the over-damped natural response and has the form

\[ i_n(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t} \]  

(1.7)

Case II: If \( \alpha_S^2 = \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and equal. This results in the critically damped natural response and has the form

\[ i_n(t) = e^{-\alpha_S t} (k_1 + k_2 t) \]  

(1.8)

Case III: If \( \omega_0^2 > \alpha_S^2 \), the roots \( s_1 \) and \( s_2 \) are complex conjugates. This is known as the underdamped or oscillatory natural response and has the form

\[ i_n(t) = e^{-\alpha_S t} (k_1 \cos \omega_{ns} t + k_2 \sin \omega_{ns} t) = k_3 e^{-\alpha_S t} (\cos \omega_{ns} t + \phi) \]  

(1.9)

A typical overdamped response is shown in Figure 1.2 where it is assumed that \( i_n(0) = 0 \). This plot was created with the following MATLAB code:
A typical critically damped response is shown in Figure 1.3 where it is assumed that $i_n(0) = 0$. This plot was created with the following MATLAB code:

```matlab
t=0: 0.01: 6; ft=420.*t.*(exp(2.45.*t)); plot(t,ft); grid; xlabel('t'); ylabel('f(t)'); title('Critically Damped Response for 420.*t.*(exp(2.45.*t))')
```

![Critically Damped Response](image)

**Figure 1.3. Typical critically damped response**

A typical underdamped response is shown in Figure 1.4 where it is assumed that $i_n(0) = 0$. This plot was created with the following MATLAB code:

```matlab
t=0: 0.01: 6; ft=420.*t.*(exp(-2.45.*t)); plot(t,ft); grid; xlabel('t'); ylabel('f(t)'); title('Underdamped Response for 420.*t.*(exp(-2.45.*t))')
```

![Underdamped Response](image)

**Figure 1.4. Typical underdamped response**
1.3 Response of Series RLC Circuits with DC Excitation

Depending on the circuit constants $R$, $L$, and $C$, the total response of a series $RLC$ circuit that is excited by a DC source, may be overdamped, critically damped, or underdamped. In this section we will derive the total response of series $RLC$ circuits that are excited by DC sources.

Example 1.1

For the circuit of Figure 1.5, $i_L(0) = 5$ A, $v_C(0) = 2.5$ V, and the 0.5 Ω resistor represents the resistance of the inductor. Compute and sketch $i(t)$ for $t > 0$.

Solution:

This circuit can be represented by the integrodifferential equation
Differentiating and noting that the derivatives of the constants \( v_c(0) \) and 15 are zero, we obtain the homogeneous differential equation

\[
R \frac{di}{dt} + L \frac{d^2i}{dt^2} + \frac{i}{C} = 0
\]

or

\[
\frac{d^2i}{dt^2} + \frac{R}{L} \frac{di}{dt} + \frac{i}{LC} = 0
\]

and by substitution of the known values \( R, \ L, \) and \( C \)

\[
\frac{d^2i}{dt^2} + 500 \frac{di}{dt} + 60000i = 0
\]

The roots of the characteristic equation of (1.11) are \( s_1 = -200 \) and \( s_2 = -300 \). The total response is just the natural response and for this example it is overdamped. Therefore, from (1.7),

\[
i(t) = i_n(t) = k_1 e^{s_1t} + k_2 e^{s_2t} = k_1 e^{-200t} + k_2 e^{-300t}
\]

(1.12)

The constants \( k_1 \) and \( k_2 \) can be evaluated from the initial conditions. Thus from the first initial condition \( i_L(0) = i(0) = 5 \) A and (1.12) we get

\[
i(0) = k_1 e^0 + k_2 e^0 = 5
\]

\[
k_1 + k_2 = 5
\]

(1.13)

We need another equation in order to compute the values of \( k_1 \) and \( k_2 \). With this equation we will make use of the second initial condition, that is, \( v_c(0) = 2.5 \) V. Since \( i_c(t) = i(t) = C \frac{dv_c}{dt} \), we differentiate (1.12), we evaluate it at \( t = 0^+ \), and we equate it with this initial condition. Then,

\[
\frac{di}{dt} = -200k_1 e^{-200t} -300k_2 e^{-300t} \quad \text{and} \quad \left. \frac{di}{dt} \right|_{t = 0^+} = -200k_1 -300k_2
\]

(1.14)

Also, at \( t = 0^+ \),

\[
R i(0^+) + L \left. \frac{di}{dt} \right|_{t = 0^+} + v_c(0^+) = 15
\]
and solving for \( \frac{di}{dt} \Big|_{t=0^+} \) we get

\[
\frac{di}{dt} \Big|_{t=0^+} = \frac{15 - 0.5 \times 5 - 2.5}{10^{-3}} = 10000
\]  

(1.15)

Next, equating (1.14) with (1.15) we get:

\[
-200k_1 - 300k_2 = 10000
\]

\[-k_1 - 1.5k_2 = 50
\]  

(1.16)

Simultaneous solution of (1.13) and (1.16) yields \( k_1 = 115 \) and \( k_2 = -110 \). By substitution into (1.12) we find the total response as

\[
i(t) = i_n(t) = 115e^{-200t} - 110e^{-300t}
\]  

(1.17)

Check with MATLAB:

```matlab
syms t; % Define symbolic variable t
R=0.5; L=10^(-3); C=100*10^(-3)/6; % Circuit constants
y0=115*exp(-200*t)-110*exp(-300*t); % Let solution i(t) = y0
y1=diff(y0); % Compute the first derivative of y0, i.e., di/dt
y2=diff(y0,2); % Compute the second derivative of y0, i.e., di2/dt2
% Substitute the solution i(t), i.e., equ (1.17)
% into differential equation of (1.11) to verify
% that correct solution was obtained.
% We must also verify that the initial
% conditions are satisfied
y=y2+500*y1+60000*y0;
i0=115*exp(-200*0)-110*exp(-300*0);
vC0=-R*i0-L*(-23000*exp(-200*0)+33000*exp(-300*0))+15;
fprintf('
');
disp('Solution was entered as y0 = '); disp(y0);
disp('1st derivative of solution is y1 = '); disp(y1);
disp('2nd derivative of solution is y2 = '); disp(y2);
disp('Differential equation is satisfied since y = y2+y1+y0 = '); disp(y);
disp('1st initial condition is satisfied since at t = 0, i0 = '); disp(i0);
disp('2nd initial condition is also satisfied since vC+vL+vR=15 and vC0 = '); disp(vC0);
fprintf('
')
Solution was entered as y0 =
115*exp(-200*t)-110*exp(-300*t)
```
1st derivative of solution is \( y_1 = -23000\exp(-200t)+33000\exp(-300t) \)

2nd derivative of solution is \( y_2 = 4600000\exp(-200t)-9900000\exp(-300t) \)

Differential equation is satisfied since \( y = y_2+y_1+y_0 = 0 \)

1st initial condition is satisfied since at \( t = 0 \), \( i_0 = 5 \)

2nd initial condition is also satisfied since \( v_C+v_L+v_R=15 \) and \( v_C0 = 2.5000 \)

We will use the following MATLAB code to sketch \( i(t) \).

```matlab
t=0: 0.0001: 0.025; i1=115.*(exp(200.*t)); i2=110.*(exp(300.*t)); iT=i1-i2;...
plot(t,i1,t,i2,t,iT); grid; xlabel('t'); ylabel('i1, i2, iT'); title('Response iT for Example 1.1')
```

![Figure 1.6. Plot for i(t) of Example 1.1](image_url)

In the above example, differentiation eliminated (set equal to zero) the right side of the differential equation and thus the total response was just the natural response. A different approach however, may not set the right side equal to zero, and therefore the total response will contain both the natural and forced components. To illustrate, we will use the following approach.

The capacitor voltage, for all time \( t \), may be expressed as \( v_C(t) = \frac{1}{C} \int_{-\infty}^{t} idt \) and as before, the circuit can be represented by the integrodifferential equation

\[
Ri + L \frac{di}{dt} + \frac{1}{C} \int_{-\infty}^{t} idt = 15u_0(t)
\]

(1.18)
and since

\[ i = i_c = C \frac{dv_c}{dt} \]

we rewrite (1.18) as

\[ RC \frac{dv_c}{dt} + LC \frac{d^2v_c}{dt^2} + v_c = 15u_0(t) \quad (1.19) \]

We observe that this is a non-homogeneous differential equation whose solution will have both the natural and the forced response components. Of course, the solution of (1.19) will give us the capacitor voltage \( v_c(t) \). This presents no problem since we can obtain the current by differentiation of the expression for \( v_c(t) \).

Substitution of the given values into (1.19) yields

\[ \frac{50}{6} \times 10^{-3} \frac{dv_c}{dt} + 1 \times 10^{-3} \times \frac{100}{6} \times 10^{-3} \frac{d^2v_c}{dt^2} + v_c = 15u_0(t) \]

or

\[ \frac{d^2v_c}{dt^2} + 500 \frac{dv_c}{dt} + 60000v_c = 9 \times 10^5 u_0(t) \quad (1.20) \]

The characteristic equation of (1.20) is the same as of that of (1.11) and thus the natural response is

\[ v_{cn}(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t} = k_1 e^{-2000t} + k_2 e^{-3000t} \quad (1.21) \]

Since the right side of (1.20) is a constant, the forced response will also be a constant and we denote it as \( v_{cf} = k_3 \). By substitution into (1.20) we get

\[ 0 + 0 + 60000k_3 = 900000 \]

or

\[ v_{cf} = k_3 = 15 \quad (1.22) \]

The total solution then is the summation of (1.21) and (1.22), that is,

\[ v_c(t) = v_{cn}(t) + v_{cf} = k_1 e^{-2000t} + k_2 e^{-3000t} + 15 \quad (1.23) \]

As before, the constants \( k_1 \) and \( k_2 \) will be evaluated from the initial conditions. First, using \( v_c(0) = 2.5 \text{ V} \) and evaluating (1.23) at \( t = 0 \), we get

\[ v_c(0) = k_1 e^0 + k_2 e^0 + 15 = 2.5 \]

or
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\[ k_1 + k_2 = -12.5 \]  \( (1.24) \)

Also,

\[ i_L = i_C = C \frac{dv_C}{dt}, \quad \frac{dv_C}{dt} = \frac{i_L}{C} \quad \text{and} \quad \left. \frac{dv_C}{dt} \right|_{t=0} = i_L(0) = \frac{5}{100} \times 10^{-3} = 300 \]  \( (1.25) \)

Next, we differentiate (1.23), we evaluate it at \( t = 0 \) and equate it with (1.25). Then,

\[ \frac{dv_C}{dt} = -200k_1e^{-200t} - 300k_2e^{-300t} \quad \text{and} \quad \left. \frac{dv_C}{dt} \right|_{t=0} = -200k_1 - 300k_2 \]  \( (1.26) \)

Equating the right sides of (1.25) and (1.26) we get

\[ -200k_1 - 300k_2 = 300 \]

or

\[ -k_1 - 1.5k_2 = 1.5 \]  \( (1.27) \)

From (1.24) and (1.27), we get \( k_1 = -34.5 \) and \( k_2 = 22 \). By substitution into (1.23), we obtain the total solution as

\[ v_C(t) = (22e^{-300t} - 34.5e^{-200t} + 15)u_0(t) \]  \( (1.28) \)

Check with MATLAB:

```matlab
syms t
y0=22*exp(-300*t)-34.5*exp(-200*t)+15; % Define symbolic variable t
y1=diff(y0) % The total solution y(t)
y1 = -6600*exp(-300*t)+6900*exp(-200*t) % The first derivative of y(t)
y2=diff(y0,2) % The second derivative of y(t)
y2 = 1980000*exp(-300*t) - 1380000*exp(-200*t)
y=y2+500*y1+60000*y0 % Summation of y and its derivatives
y = 900000
```

Using the expression for \( v_C(t) \) we can find the current as

\[ i = i_L = i_C = C \frac{dv_C}{dt} = \frac{100}{6} \times 10^{-3} (6900e^{-200t} - 6600e^{-300t}) = 115e^{-200t} - 110e^{-300t} \text{ A} \]  \( (1.29) \)
Response of Series RLC Circuits with AC Excitation

We observe that (1.29) is the same as (1.17).

We will use the following MATLAB code to sketch \( i(t) \).

\[
t = 0: 0.001: 0.03; \quad v_{c1} = 22.0 \times (\exp(-300.0t)); \quad v_{c2} = -34.5 \times (\exp(-200.0t)); \quad v_{c3} = 15;...
\]

\[
v_{cT} = v_{c1} + v_{c2} + v_{c3}; \quad \text{plot}(t, v_{c1}, t, v_{c2}, t, v_{c3}, t, v_{cT}); \quad \text{grid}; \quad \text{xlabel}(t);...
\]

\[
\text{ylabel}('v_{c1}, v_{c2}, v_{c3}, v_{cT}'); \quad \text{title}('Response v_{cT} for Example 1.1')
\]

![Figure 1.7. Plot for \( v_c(t) \) of Example 1.1](image)

1.4 Response of Series RLC Circuits with AC Excitation

The total response of a series RLC circuit, which is excited by a sinusoidal source, will also consist of the natural and forced response components. As we found in the previous section, the natural response can be overdamped, or critically damped, or underdamped. The forced component will be a sinusoid of the same frequency as that of the excitation, and since it represents the AC steady-state condition, we can use phasor analysis to find it. The following example illustrates the procedure.

**Example 1.2**

For the circuit of Figure 1.8, \( i_L(0) = 5 \, \text{A} \), \( v_C(0) = 2.5 \, \text{V} \), and the 0.5 \, \Omega \) resistor represents the resistance of the inductor. Compute and sketch \( i(t) \) for \( t > 0 \).

**Solution:**

This circuit is the same as that of Example 1.1 except that the circuit is excited by a sinusoidal source; therefore it can be represented by the integrodifferential equation

\[
Ri + L \frac{di}{dt} + \frac{1}{C} \int_0^t idt + v_C(0) = 200 \cos 10000t \quad t > 0
\]  

(1.30)
whose solution consists of the summation of the natural and forced responses. We know its natural response from the previous example. We start with
\[
    i(t) = i_n(t) + i_f(t) = k_1 e^{-2000t} + k_2 e^{-3000t} + i_f(t)
\]  
(1.31)

where the constants \( k_1 \) and \( k_2 \) will be evaluated from the initial conditions after \( i_f(t) \) has been found. The steady state (or forced) response will have the form \( i_f(t) = k_3 \cos(10,000t + \theta) \) in the time domain (\( t \)-domain) and has the form \( k_3 \angle \theta \) in the frequency domain (\( j\omega \)-domain).

To find \( i_f(t) \) we will use the phasor analysis relation \( I = V/Z \) where \( I \) is the phasor current, \( V \) is the phasor voltage, and \( Z \) is the impedance of the phasor circuit which, as we know, is
\[
    Z = R + j(\omega L - 1/\omega C) = \sqrt{R^2 + (\omega L - 1/\omega C)^2} \angle \sin^{-1}(\omega L - 1/\omega C)/R
\]  
(1.32)

The inductive and capacitive reactances are
\[
X_L = \omega L = 10^4 \times 10^{-3} = 10 \ \Omega
\]
and
\[
X_C = \frac{1}{\omega C} = \frac{1}{10^4 \times (100/6)10^{-3}} = 6 \times 10^{-3} \ \Omega
\]

Then,
\[
R^2 = (0.5)^2 = 0.25 \quad \text{and} \quad (\omega L - 1/\omega C)^2 = (10 - 6 \times 10^{-3})^2 = 99.88
\]

Also,
\[
\tan^{-1}(\omega L - 1/\omega C)/R = \tan^{-1}\left(\frac{10 - 6 \times 10^{-3}}{0.5}\right) = \tan^{-1}\left(\frac{9.994}{0.5}\right)
\]

and this yields \( \theta = 1.52 \text{ rads} = 87.15^\circ \). Then, by substitution into (1.32),
\[
Z = \sqrt{0.25 + 99.88} \angle 0^\circ = 10 \angle 87.15^\circ
\]
and thus
\[
I = \frac{V}{Z} = \frac{200 \angle 0^\circ}{10 \angle 87.15^\circ} = 20 \angle -87.15^\circ \Leftrightarrow 20 \cos(10000t - 87.15^\circ) = i_f(t)
\]
The total solution is
\[ i(t) = i_n(t) + i_f(t) = k_1 e^{-200t} + k_2 e^{-300t} + 20 \cos(10000t - 87.15^\circ) \] (1.33)

The constants \( k_1 \) and \( k_2 \) are evaluated from the initial conditions. From (1.33) and the first initial condition \( i_L(0) = 5 \text{ A} \) we get
\[ i(0) = k_1 e^0 + k_2 e^0 + 20 \cos(-87.15^\circ) = 5 \]
\[ i(0) = k_1 + k_2 + 20 \times 0.05 = 5 \]
\[ k_1 + k_2 = 4 \] (1.34)

We need another equation in order to compute the values of \( k_1 \) and \( k_2 \). This equation will make use of the second initial condition, that is, \( v_C(0) = 2.5 \text{ V} \). Since \( i_C(t) = i(t) = C \frac{dv_C}{dt} \), we differentiate (1.33), we evaluate it at \( t = 0 \), and we equate it with this initial condition. Then,
\[ \frac{di}{dt} = -200k_1 e^{-200t} - 300k_2 e^{-300t} - 2 \times 10^5 \sin(10000t - 87.15^\circ) \] (1.35)

and at \( t = 0^+ \),
\[ \left. \frac{di}{dt} \right|_{t=0} = -200k_1 - 300k_2 - 2 \times 10^6 \sin(-87.15^\circ) = -200k_1 - 300k_2 + 2 \times 10^5 \] (1.36)

Also, at \( t = 0^+ \)
\[ Ri(0^+) + L \left. \frac{di}{dt} \right|_{t=0^+} + v_C(0^+) = 200 \cos(0) = 200 \]
and solving for \( \left. \frac{di}{dt} \right|_{t=0^+} \) we get
\[ \left. \frac{di}{dt} \right|_{t=0^+} = \frac{200 - 0.5 \times 5 - 2.5}{10^{-3}} = 195000 \] (1.37)

Next, equating (1.36) with (1.37) we get
\[ -200k_1 - 300k_2 = -5000 \]
or
\[ k_1 + 1.5k_2 = 25 \] (1.38)
Simultaneous solution of (1.34) and (1.38) yields \( k_1 = -38 \) and \( k_2 = 42 \). Then, by substitution into (1.31), the total response is

\[
i(t) = -38e^{-200t} + 42e^{-300t} + 20\cos(10000t - 87.15\degree) \text{ A} 
\]  

(1.39)

The plot is shown in Figure 1.9 and was created with the following MATLAB code:

\[
t = 0: 0.005: 0.20; i1 = -38.*\exp(-200.*t); i2 = 42.*\exp(-300.*t);... \\
i3 = 20.*\cos(10000.*t - 87.15.*\pi./180); iT = i1 + i2 + i3; plot(t,i1,t,i2,t,i3,t,iT); grid; xlabel('t'); ylabel('i1, i2, i3, iT'); title('Response iT for Example 1.2')
\]

Figure 1.9. Plot for \( i(t) \) of Example 1.2

### 1.5 The Parallel GLC Circuit

Consider the circuit of Figure 1.10 where the initial conditions are \( i_L(0) = I_0, \ v_C(0) = V_0, \) and \( u_0(t) \) is the unit step function. We want to find an expression for the voltage \( v(t) \) for \( t > 0 \).

![Parallel RLC circuit](image)

For this circuit
By differentiation,

\[ \frac{dv}{dt} + \frac{v}{L} = 0 \quad t > 0 \tag{1.40} \]

To find the forced response, we must first specify the nature of the excitation \( i_s \), that is DC or AC.

If \( i_s \) is DC (\( v_s \) = constant), the right side of (1.40) will be zero and thus the forced response component \( v_f = 0 \). If \( i_s \) is AC (\( i_s = I \cos(\omega t + \theta) \)), the right side of (1.40) will be another sinusoid and therefore \( v_f = V \cos(\omega t + \phi) \). Since in this section we are concerned with DC excitations, the right side will be zero and thus the total response will be just the natural response.

The natural response is found from the homogeneous equation of (1.40), that is,

\[ C \frac{dv}{dt}^2 + G \frac{dv}{dt} + \frac{v}{L} = 0 \tag{1.41} \]

whose characteristic equation is

\[ Cs^2 + Gs + \frac{1}{L} = 0 \]

or

\[ s^2 + \frac{G}{C}s + \frac{i}{LC} = 0 \]

from which

\[ s_1, s_2 = -\frac{G}{2C} \pm \sqrt{\frac{G^2}{4C^2} - \frac{1}{LC}} \tag{1.42} \]

and with the following notations,

\[ \alpha_p = \frac{G}{2C} \quad \omega_0 = \frac{1}{\sqrt{LC}} \quad \beta_p = \sqrt{\alpha_p^2 - \omega_0^2} \quad \omega_{np} = \sqrt{\omega_0^2 - \alpha_p^2} \tag{1.43} \]

where the subscript \( p \) stands for parallel circuit, we can express (1.42) as
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\[ s_1, s_2 = -\alpha_p \pm \sqrt{\alpha_p^2 - \omega_0^2} = -\alpha_p \pm \beta_p \quad \text{if} \quad \alpha_p^2 > \omega_0^2 \quad (1.44) \]

or

\[ s_1, s_2 = -\alpha_p \pm \sqrt{\omega_0^2 - \alpha_p^2} = -\alpha_p \pm \omega_n \omega_p \quad \text{if} \quad \omega_0^2 > \alpha_p^2 \quad (1.45) \]

Note: From (1.4) and (1.43) we observe that \( \alpha_s \neq \alpha_p \)

As in a series circuit, the natural response \( v_n(t) \) can be overdamped, critically damped, or underdamped.

**Case I:** If \( \alpha_p^2 > \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and unequal. This results in the overdamped natural response and has the form

\[ v_n(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t} \quad (1.46) \]

**Case II:** If \( \alpha_p^2 = \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and equal. This results in the critically damped natural response and has the form

\[ v_n(t) = e^{-\alpha_p t} (k_1 + k_2 t) \quad (1.47) \]

**Case III:** If \( \omega_0^2 > \alpha_p^2 \), the roots \( s_1 \) and \( s_2 \) are complex conjugates. This results in the underdamped or oscillatory natural response and has the form

\[ v_n(t) = e^{-\alpha_p t} (k_1 \cos \omega_n t + k_2 \sin \omega_n t + \varphi) \quad (1.48) \]

### 1.6  Response of Parallel GLC Circuits with DC Excitation

Depending on the circuit constants \( G \) (or \( R \)), \( L \), and \( C \), the natural response of a parallel \( GLC \) circuit may be overdamped, critically damped or underdamped. In this section we will derive the total response of a parallel \( GLC \) circuit which is excited by a DC source using the following example.

**Example 1.3**

For the circuit of Figure 1.11, \( i_L(0) = 2 \text{ A} \) and \( v_C(0) = 5 \text{ V} \). Compute and sketch \( v(t) \) for \( t > 0 \).
Solution:

We could write the integrodifferential equation that describes the given circuit, differentiate, and find the roots of the characteristic equation from the homogeneous differential equation as we did in the previous section. However, we will skip these steps and start with

\[ v(t) = v_f(t) + v_n(t) \]  

(1.49)

and when steady-state conditions have been reached we will have \( v = v_L = L(di/dt) = 0 \), \( v_f = 0 \) and \( v(t) = v_n(t) \).

To find out whether the natural response is overdamped, critically damped, or oscillatory, we need to compute the values of \( \alpha_p \) and \( \omega_0 \) using (1.43) and the values of \( s_1 \) and \( s_2 \) using (1.44) or (1.45). Then will use (1.46), or (1.47), or (1.48) as appropriate. For this example,

\[ \alpha_p = \frac{G}{2C} = \frac{1}{2RC} = \frac{1}{2 \times 32 \times 1/640} = 10 \]

or

\[ \alpha_p^2 = 100 \]

and

\[ \omega_0^2 = \frac{1}{LC} = \frac{1}{10 \times 1/640} = 64 \]

Then

\[ s_1, s_2 = -\alpha_p \pm \sqrt{\alpha_p^2 - \omega_0^2} = -10 \pm 6 \]

or \( s_1 = -4 \) and \( s_2 = -16 \). Therefore, the natural response is overdamped and from (1.46) we get

\[ v(t) = v_n(t) = k_1 e^{s_1t} + k_2 e^{s_2t} = k_1 e^{-4t} + k_2 e^{-16t} \]  

(1.50)

and the constants \( k_1 \) and \( k_2 \) will be evaluated from the initial conditions.

From the initial condition \( v_C(0) = v(0) = 5 \) V and (1.50) we get
\[ v(0) = k_1 e^0 + k_2 e^0 = 5 \]

or
\[ k_1 + k_2 = 5 \]  
(1.51)

The second equation that is needed for the computation of the values of \( k_1 \) and \( k_2 \) is found from another initial condition, that is, \( i_L(0) = 2 \, \text{A} \). Since \( i_C(t) = C \frac{dv}{dt} = C \frac{dy}{dt} \), we differentiate (1.50), evaluate it at \( t = 0^+ \), and we equate it with this initial condition. Then,

\[ \frac{dv}{dt} = -4k_1 e^{-4t} - 16k_2 e^{-16t} \text{ and } \left. \frac{dv}{dt} \right|_{t = 0^+} = -4k_1 - 16k_2 \]  
(1.52)

Also, at \( t = 0^+ \)

\[ \frac{l}{R} v(0^+) + i_L(0^+) + \left. C \frac{dv}{dt} \right|_{t = 0^+} = 10 \]

and solving for \( \left. \frac{dv}{dt} \right|_{t = 0^+} \) we get

\[ \left. \frac{dv}{dt} \right|_{t = 0^+} = \frac{10 - 5/32 - 2}{1/640} = 502 \]  
(1.53)

Next, equating (1.52) with (1.53) we get

\[ -4k_1 - 16k_2 = 502 \]

or

\[ -2k_1 - 8k_2 = 251 \]  
(1.54)

Simultaneous solution of (1.51) and (1.54) yields \( k_1 = 291/6 \), \( k_2 = -261/6 \), and by substitution into (1.50) we get the total response as

\[ v(t) = v_c(t) = \frac{291}{6} e^{-4t} - \frac{261}{6} e^{-16t} = \frac{1}{6} (291 e^{-4t} - 261 e^{-16t}) \, V \]  
(1.55)

Check with MATLAB:

\begin{verbatim}
> syms t % Define symbolic variable t
> y0=291*exp(-4*t)/6-261*exp(-16*t)/6; % Let solution v(t) = y0
> y1=diff(y0) % Compute and display first derivative
> y1 =
> -194*exp(-4*t)+696*exp(-16*t)
\end{verbatim}
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\[ y_2 = \text{diff}(y0,2) \quad \% \text{Compute and display second derivative} \]

\[ y_2 = 776 \exp(-4t) - 11136 \exp(-16t) \]

\[ y = y_2/640 + y_1/32 + y_0/10 \quad \% \text{Verify that (1.40) is satisfied} \]

\[ y = 0 \]

The plot is shown in Figure 1.12 where we have used the following MATLAB code:

\[
t=0:0.01:1; \ v1=(291./6).*(exp(4.*t)); \ v2=-(261./6).*(exp(-16.*t));
\]

\[
\vT = v1 + v2; \ \text{plot}(t,v1,t,v2,t,vT); \ \text{grid}; \ \text{xlabel('t')};
\]

\[
\text{ylabel('v1, v2, vT')}; \ \text{title('Response vT for Example 1.3')}
\]

![Figure 1.12. Plot for \( v(t) \) of Example 1.3](image)

From the plot of Figure 1.12, we observe that \( v(t) \) attains its maximum value somewhere in the interval 0.10 and 0.12 sec., and the maximum voltage is approximately 24 V. If we desire to compute precisely the maximum voltage and the exact time it occurs, we can find the derivative of (1.55), set it equal to zero, and solve for \( t \). Thus,

\[
\frac{dv}{dt} \bigg|_{t=0} = -1164e^{-4t} + 4176e^{-16t} = 0 \quad (1.56)
\]

Division of (1.56) by \( e^{-16t} \) yields

\[-1164e^{12t} + 4176 = 0 \]

or
\[ e^{12t} = \frac{348}{97} \]

or

\[ 12t = \ln\left(\frac{348}{97}\right) = 1.2775 \]

and

\[ t = t_{\text{max}} = \frac{1.2775}{12} = 0.106 \text{ s} \]

By substitution into (1.55)

\[ v_{\text{max}} = \frac{1}{6} \left(291e^{-4.106t} - 261e^{-16.106t}\right) = 23.76 \text{ V} \quad (1.57) \]

A useful quantity, especially in electronic circuit analysis, is the *settling time*, denoted as \( t_S \), and it is defined as the time required for the voltage to drop to 1% of its maximum value. Therefore, \( t_S \) is an indication of the time it takes for \( v(t) \) to damp-out, meaning to decrease the amplitude of \( v(t) \) to approximately zero. For this example, \( 0.01 \times 23.76 = 0.2376 \text{ V} \), and we can find \( t_S \) by substitution into (1.55). Then,

\[ 0.01v_{\text{max}} = 0.2376 = \frac{1}{6} \left(291e^{-4t} - 261e^{-16t}\right) \quad (1.58) \]

and we need to solve for the time \( t \). To simplify the computation, we neglect the second term inside the parentheses of (1.58) since this component of the voltage damps out much faster than the other component. This expression then simplifies to

\[ 0.2376 = \frac{1}{6} \left(291e^{-4t_S}\right) \]

or

\[ -4t_S = \ln(0.005) = (-5.32) \]

or

\[ t_S = 1.33 \text{ s} \quad (1.59) \]

**Example 1.4**

For the circuit of Figure 1.13, \( i_L(0) = 2 \text{ A} \) and \( v_C(0) = 5 \text{ V} \), and the resistor is to be adjusted so that the natural response will be critically damped. Compute and sketch \( v(t) \) for \( t > 0 \).
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Solution:

Since the natural response is to be critically damped, we must have \( \omega_0^2 = 64 \) because the L and C values are the same as in the previous example. Please refer to (1.43). We must also have

\[
\alpha_p = \frac{G}{2C} = \frac{1}{2RC} = \omega_0 = \frac{\sqrt{T}}{\sqrt{LC}} = 8
\]

or

\[
\frac{1}{R} = 8 \times \frac{2}{640} = \frac{1}{40}
\]

or \( R = 40 \, \Omega \) and thus \( s_1 = s_2 = -\alpha_p = -8 \). The natural response will have the form

\[
v(t) = v_n(t) = e^{-\alpha_p t}(k_1 + k_2t) \quad \text{or} \quad v(t) = v_n(t) = e^{-st}(k_1 + k_2t)
\]

(1.60)

Using the initial condition \( v_C(0) = 5 \, V \) and evaluating (1.60) at \( t = 0 \), we get

\[
v(0) = e^0(k_1 + k_20) = 5
\]

or

\[
k_1 = 5
\]

(1.61)

and (1.60) simplifies to

\[
v(t) = e^{-st}(5 + k_2t)
\]

(1.62)

As before, we need to compute the derivative \( dv/dt \) in order to apply the second initial condition and find the value of the constant \( k_2 \).

We obtain the derivative using MATLAB as follows:

```matlab
syms t k2; v0=exp(-8*t)*(5+k2*t); v1=diff(v0); % v1 is 1st derivative of v0
v1 = -8*exp(-8*t) * (5+k2*t) + exp(-8*t) * k2
```

Figure 1.13. Circuit for Example 1.4
Then,
\[ \frac{dv}{dt} = -8e^{-8t}(5 + k_2t) + k_2e^{-8t} \]
and
\[ \left. \frac{dv}{dt} \right|_{t=0} = -40 + k_2 \]  \hspace{1cm} (1.63)

Also, \( i_C = \frac{dv}{dt} \) or \( \frac{dv}{dt} = \frac{i_C}{C} \) and

\[ \left. \frac{dv}{dt} \right|_{t=0^+} = \frac{i_C(0^+)}{C} = \frac{I_S(i_C(0^+) - i_L(0^+))}{C} \]  \hspace{1cm} (1.64)

or

\[ \left. \frac{dv}{dt} \right|_{t=0} = \frac{I_S - v_C(0) + R - i_L(0)}{C} = \frac{10 - 5/40 - 2}{1/640} = \frac{7.875}{1/640} = 5040 \]  \hspace{1cm} (1.65)

Equating (1.63) with (1.65) and solving for \( k_2 \) we get

\[ -40 + k_2 = 5040 \]

or

\[ k_2 = 5080 \]  \hspace{1cm} (1.66)

and by substitution into (1.62), we obtain the total solution as

\[ v(t) = e^{-8t}(5 + 5080t) \]  \hspace{1cm} (1.67)

Check with MATLAB:

```matlab
syms t; y0=exp(-8*t)*(5+5080*t); y1=diff(y0); % Compute 1st derivative
y1 = -8*exp(-8*t)*(5+5080*t)+5080*exp(-8*t)
y2=diff(y0,2); % Compute 2nd derivative
y2 = 64*exp(-8*t)*(5+5080*t)-81280*exp(-8*t)
y=y2/640+y1/40+y0/10; % Verify differential equation, see (1.40)
y = 0
```

The plot is shown in Figure 1.14 where we have used the following MATLAB code:
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t = 0: 0.01: 1; vt = exp(-8.*t).*(5+5080.*t); plot(t,vt); grid; xlabel('t'); ylabel('vt'); title('Response vt for Example 1.4')

\[ t = 0: 0.01: 1; vt = \exp(-8t) \cdot (5+5080t); \text{plot}(t,vt); \text{grid}; \text{xlabel}('t'); \text{ylabel}('vt'); \text{title}('Response vt for Example 1.4') \]

Figure 1.14. Plot for \( v(t) \) of Example 1.4

By inspection of (1.67), we see that at \( t = 0 \), \( v(t) = 5 \text{ V} \) and thus the second initial condition is satisfied. We can verify that the first initial condition is also satisfied by differentiation of (1.67). We can also show that \( v(t) \) approaches zero as \( t \) approaches infinity with L'Hôpital's rule as follows:

\[
\lim_{t \to \infty} v(t) = \lim_{t \to \infty} e^{-8t} (5 + 5080t) = \lim_{t \to \infty} \frac{5 + 5080t}{e^{8t}} = \lim_{t \to \infty} \frac{d(5 + 5080t)/dt}{d(e^{8t})/dt} = \lim_{t \to \infty} \frac{5080}{8e^{8t}} = 0 \tag{1.68}
\]

Example 1.5

For the circuit of Figure 1.15, \( i_L(0) = 2 \text{ A} \) and \( v_C(0) = 5 \text{ V} \). Compute and sketch \( v(t) \) for \( t > 0 \).

\[
\text{Solution:}
\]

This is the same circuit as the that of the two previous examples except that the resistance has been increased to 50 Ω. For this example,
\[ \alpha_p = \frac{G}{2C} = \frac{1}{2RC} = \frac{1}{2 \times 50 \times 1/640} = 6.4 \]

or

\[ \alpha_p^2 = 40.96 \]

and as before,

\[ \omega_n^2 = \frac{1}{L C} = \frac{1}{10 \times 1/640} = 64 \]

Also, \( \omega_n^2 > \alpha_p^2 \). Therefore, the natural response is underdamped with natural frequency

\[ \omega_{np} = \sqrt{\omega_n^2 - \alpha_p^2} = \sqrt{64 - 40.96} = \sqrt{23.04} = 4.8 \]

Since \( v_f = 0 \), the total response is just the natural response. Then, from (1.48),

\[ v(t) = v_n(t) = ke^{-\alpha_p t} \cos(\omega_{np}t + \varphi) = ke^{-6.4t} \cos(4.8t + \varphi) \] (1.69)

and the constants \( k \) and \( \varphi \) will be evaluated from the initial conditions.

From the initial condition \( v_C(0) = v(0) = 5 \text{ V} \) and (1.69) we get

\[ v(0) = ke^0 \cos(0 + \varphi) = 5 \]

or

\[ k \cos \varphi = 5 \] (1.70)

To evaluate the constants \( k \) and \( \varphi \) we differentiate (1.69), we evaluate it at \( t = 0^+ \), and we equate these two expressions. Using MATLAB we get:

```matlab
syms t k phi; y0=k*exp(-6.4*t)*cos(4.8*t+phi); y1=diff(y0)
y1 = -32/5*k*exp(-32/5*t)*cos(24/5*t+phi)-24/5*k*exp(-32/5*t)*sin(24/5*t+phi)
pretty(y1)
```

Thus,

\[ \frac{dv}{dt} = -6.4ke^{-6.4t}\cos(4.8t + \varphi) - 4.8ke^{-6.4t}\sin(4.8t + \varphi) \] (1.71)

and
Response of Parallel GLC Circuits with DC Excitation

\[ \frac{dv}{dt} \bigg| _{t=0} = -6.4k \cos \varphi - 4.8k \sin \varphi \]

By substitution of (1.70), the above expression simplifies to

\[ \frac{dv}{dt} \bigg| _{t=0} = -32 - 4.8k \sin \varphi \] (1.72)

Also, \( i_c = C \frac{dv}{dt} \) or \( \frac{dv}{dt} = \frac{i_c}{C} \) and

\[ \frac{dv}{dt} \bigg| _{t=0^+} = \frac{i_c(0^+)}{C} = \frac{I_S - i_L(0^+)}{C} - i_L(0^+) \]

or

\[ \frac{dv}{dt} \bigg| _{t=0^-} = \frac{I_S - v_C(0^-)}{C} = \frac{10 - 5/50 - 2}{1/640} = 7.9 \times 640 = 5056 \] (1.73)

Equating (1.72) with (1.73) we get

\[ -32 - 4.8k \sin \varphi = 5056 \]

or

\[ k \sin \varphi = -1060 \] (1.74)

The phase angle \( \varphi \) can be found by dividing (1.74) by (1.70). Then,

\[ \frac{k \sin \varphi}{k \cos \varphi} = \tan \varphi = \frac{-1060}{5} = -212 \]

or

\[ \varphi = \tan^{-1}(-212) = -1.566 \text{ radians} = -89.73 \text{ degrees} \]

The value of the constant \( k \) is found from (1.70) as

\[ k \cos(-1.566) = 5 \]

or

\[ k = \frac{5}{\cos(-1.566)} = 1042 \]

and by substitution into (1.69), the total solution is

\[ v(t) = 1042 e^{-6.4t} \cos(4.8t - 89.73^\circ) \] (1.75)

The plot is shown in Figure 1.16 where we have used the following MATLAB code:
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\[ t=0: 0.005: 1.5; \ vt=10.42.*\exp(-6.4.*t).*\cos(4.8.*t-89.73.*\pi./180);... \]
\[ \text{plot}(t,\vt); \ \text{grid}; \ \text{xlabel}(t); \ \text{ylabel}(\vt); \ \text{title}(\text{Response } v(t) \text{ for Example 1.5}) \]

\[ \begin{array}{|c|c|}
\hline
 t (sec) & v (V) \\
\hline
 \text{Maximum} & 0.13 & 266.71 \\
\hline
 \text{Minimum} & 0.79 & -4.05 \\
\hline
\end{array} \]

We can also use a spreadsheet to plot (1.75). From the columns of that spreadsheet we can read the following maximum and minimum values and the times these occur.

Alternately, we can find the maxima and minima by differentiating the response of (1.75) and setting it equal to zero.

1.7 Response of Parallel GLC Circuits with AC Excitation

The total response of a parallel GLC (or RLC) circuit that is excited by a sinusoidal source also consists of the natural and forced response components. The natural response will be overdamped, critically damped, or underdamped. The forced component will be a sinusoid of the same frequency as that of the excitation, and since it represents the AC steady-state condition, we can use phasor analysis to find the forced response. We will derive the total response of a parallel GLC (or RLC) circuit which is excited by an AC source with the following example.

Example 1.6

For the circuit of Figure 1.17, \( i_L(0) = 2 \ \text{A} \) and \( v_C(0) = 5 \ \text{V} \). Compute and sketch \( v(t) \) for \( t > 0 \).
Response of Parallel GLC Circuits with AC Excitation

Solution:

This is the same circuit as the previous example where the DC source has been replaced by an AC source. The total response will consist of the natural response \( v_n(t) \) which we already know from the previous example, and the forced response \( v_f(t) \) which is the AC steady-state response, will be found by phasor analysis.

The \( t \)-domain to \( j\omega \)-domain transformation yields

\[
i_s(t) = 20 \sin(6400t + 90^\circ) \Rightarrow I = 20 \angle 0^\circ
\]

The admittance \( Y \) is

\[
Y = G + j\left(\omega C - \frac{1}{\omega L}\right) = \sqrt{G^2 + \left(\omega C - \frac{1}{\omega L}\right)^2} \tan^{-1}\left(\frac{\omega C - \frac{1}{\omega L}}{G}\right)
\]

where

\[
G = \frac{1}{R} = \frac{1}{50}, \quad \omega C = 6400 \times \frac{1}{6400} = 10 \quad \text{and} \quad \frac{1}{\omega L} = \frac{1}{6400 \times 10} = \frac{1}{6400}
\]

and thus

\[
Y = \sqrt{\left(\frac{1}{50}\right)^2 + \left(10 - \frac{1}{6400}\right)^2} \tan^{-1}\left(\frac{10 - \frac{1}{6400}}{\frac{1}{50}}\right) = 10 \angle 89.72^\circ
\]

Now, we find the phasor voltage \( V \) as

\[
V = I = \frac{20 \angle 0^\circ}{10 \angle 89.72^\circ} = 2 \angle -89.72^\circ
\]

and \( j\omega \)-domain to \( t \)-domain transformation yields

\[
V = 2 \angle -89.72^\circ \Rightarrow v_f(t) = 2\cos(6400t - 89.72^\circ)
\]

The total solution is

\[
v(t) = v_n(t) + v_f(t) = ke^{-6.4t}\cos(4.8t + \varphi) + 2\cos(6400t - 89.72^\circ) \quad (1.76)
\]
Now, we need to evaluate the constants \( k \) and \( \varphi \).

With the initial condition \( v_C(0) = 5 \ V \) (1.76) becomes

\[
v(0) = v_C(0) = ke^\varphi \cos \varphi + 2 \cos(-89.72^\circ) = 5
\]

or

\[
k \cos \varphi \approx 5 \quad (1.77)
\]

To make use of the second initial condition, we differentiate (1.76) using MATLAB as follows, and then we evaluate it at \( t = 0 \).

```plaintext
syms t k phi; y0=k*exp(-6.4*t)*cos(4.8*t+phi)+2*cos(6400*t-1.5688);
y1=diff(y0); % Differentiate v(t) of (1.76)
y1 =
-32/5*k*exp(-32/5*t)*cos(24/5*t+phi)-24/5*k*exp(-32/5*t)*sin(24/5*t+phi)-12800*sin(6400*t-1961/1250)
```

or

\[
\frac{dv}{dt} = -6.4ke^{-6.4t} \cos(4.8t + \varphi) - 4.8ke^{-6.4t} \sin(4.8t + \varphi) - 12800 \sin(6400t - 1.5688)
\]

and

\[
\frac{dv}{dt} \bigg|_{t=0} = -6.4k \cos \varphi - 4.8k \sin \varphi - 12800 \sin(-1.5688) = -6.4k \cos \varphi - 4.8k \sin \varphi + 12800 \quad (1.78)
\]

With (1.77) we get

\[
\frac{dv}{dt} \bigg|_{t=0} = -32 - 4.8k \sin \varphi + 12800 \approx -4.8k \sin \varphi + 12832 \quad (1.79)
\]

Also, \( i_C = C \frac{dv}{dt} \) or \( \frac{dv}{dt} = \frac{i_C}{C} \) and

\[
\frac{dv}{dt} \bigg|_{t=0^+} = \frac{i_C(0^+)}{C} = \frac{i_S(0^+) - i_L(0^+) - i_R(0^+)}{C}
\]

or

\[
\frac{dv}{dt} \bigg|_{t=0} = \frac{i_S(0^+) - v_C(0)/R - i_L(0)}{C} = \frac{20-5/50-2}{1/640} = 11456 \quad (1.80)
\]

Equating (1.79) with (1.80) and solving for \( k \) we get

\[-4.8k \sin \varphi + 12832 = 11456\]
or

\[ k \sin \varphi = 287 \]  \hspace{1cm} (1.81)

Then with (1.77) and (1.81),

\[ \frac{k \sin \varphi}{k \cos \varphi} = \tan \varphi = \frac{287}{5} = 57.4 \]

or

\[ \varphi = 1.53 \text{ rad} = 89^\circ \]

The value of the constant \( k \) is found from (1.77), that is,

\[ k = \frac{5}{\cos 89^\circ} = 279.4 \]

By substitution into (1.76), we obtain the total solution as

\[ v(t) = 279.4e^{-6.4t} \cos(4.8t + 89^\circ) + 2 \cos(6400t - 89.72^\circ) \]  \hspace{1cm} (1.82)

With MATLAB we get the plot shown in Figure 1.18.

![Figure 1.18. Plot for \( v(t) \) of Example 1.6](image)

### 1.8 Other Second Order Circuits

Second order circuits are not restricted to series \( RLC \) and parallel \( GLC \) circuits. Other second order circuits include amplifiers and filters. It is beyond the scope of this text to analyze such circuits in detail. In this section we will use the following example to illustrate the transient analysis of a second order active low-pass filter.
Example 1.7

The circuit of Figure 1.19 a known as a Multiple Feed Back (MFB) active low-pass filter. For this circuit, the initial conditions are \( v_{C1} = v_{C2} = 0 \). Compute and sketch \( v_{out}(t) \) for \( t > 0 \).

**Figure 1.19. Circuit for Example 1.7**

\[
\begin{align*}
\text{At node } v_1: \\
\frac{v_1 - v_{in}}{R_1} + C_1 \frac{dv_1}{dt} + \frac{v_1 - v_{out}}{R_2} + \frac{v_1 - v_2}{R_3} &= 0 \quad t > 0 \\
\text{At node } v_2: \\
\frac{v_2 - v_1}{R_3} &= C_2 \frac{dv_{out}}{dt} \\
\end{align*}
\]

We observe that \( v_2 = 0 \) (virtual ground).

Collecting like terms and rearranging (1.83) and (1.84) we get

\[
\left( \frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3} \right) v_1 + C_1 \frac{dv_1}{dt} - \frac{1}{R_2} v_{out} = \frac{1}{R_1} v_{in} \\
\]

and

\[
v_1 = -R_3 C_2 \frac{dv_{out}}{dt} \\
\]

Differentiation of (1.86) yields

\[
\frac{dv_1}{dt} = -R_3 C_2 \frac{dv_{out}^2}{dt^2} \\
\]
and by substitution of given numerical values into (1.85) through (1.87), we get

\[
\left(\frac{1}{2 \times 10^5} + \frac{1}{4 \times 10^3} + \frac{1}{5 \times 10^4}\right) v_1 + 25 \times 10^{-9} \frac{dv_1}{dt} - \frac{1}{4 \times 10^{-3}} v_{out} = \frac{1}{2 \times 10^{-5}} v_{in}
\]

or

\[
(0.05 \times 10^{-3}) v_1 + 25 \times 10^{-9} \frac{dv_1}{dt} - (0.25 \times 10^{-4}) v_{out} = (0.5 \times 10^{-5}) v_{in}
\]  

(1.88)

\[
v_1 = -5 \times 10^{-4} \frac{dv_{out}}{dt}
\]

(1.89)

\[
\frac{dv_1}{dt} = -5 \times 10^{-4} \frac{d^2 v_{out}}{dt^2}
\]

(1.90)

Next, substitution of (1.89) and (1.90) into (1.88) yields

\[
0.05 \times 10^{-3} \left(-5 \times 10^{-4} \frac{dv_{out}}{dt}\right) + 25 \times 10^{-9} \left(-5 \times 10^{-4}\right) \frac{d^2 v_{out}}{dt^2} - (0.25 \times 10^{-4}) v_{out} = (0.5 \times 10^{-5}) v_{in}
\]

\[
-125 \times 10^{-13} \frac{d^2 v_{out}}{dt^2} - 0.25 \times 10^{-2} \frac{dv_{out}}{dt} - (0.25 \times 10^{-4}) v_{out} = 10^{-4} v_{in}
\]

(1.91)

Division by \(-125 \times 10^{-13}\) yields

\[
\frac{d^2 v_{out}}{dt^2} + 2 \times 10^3 \frac{dv_{out}}{dt} + 2 \times 10^6 v_{out} = (-1.6 \times 10^5) v_{in}
\]

or

\[
\frac{d^2 v_{out}}{dt^2} + 2 \times 10^3 \frac{dv_{out}}{dt} + 2 \times 10^6 v_{out} = -10^6 \cos 6280t
\]

(1.92)

We use MATLAB to find the roots of the characteristic equation of (1.92).

```matlab
syms s; y0=solve('s^2+2*10^3*s+2*10^6')
y0 =
[ -1000+1000*i]
[ -1000-1000*i]
```

that is,
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\[ s_1, s_2 = -\alpha \pm j\beta = -1000 \pm j1000 = 1000(-1 \pm j1) \]

We cannot classify the given circuit as series or parallel and therefore, we should not use the damping ratio \(\alpha_s\) or \(\alpha_p\). Instead, for the natural response \(v_n(t)\) we will use the general expression

\[ v_n(t) = Ae^{s_1t} + Be^{s_2t} = e^{-\alpha t}(k_1 \cos \beta t + k_2 \sin \beta t) \quad (1.93) \]

where

\[ s_1, s_2 = -\alpha \pm j\beta = -1000 \pm j1000 \]

Therefore, the natural response is oscillatory and has the form

\[ v_n(t) = e^{-1000t}(k_1 \cos 1000t + k_2 \sin 1000t) \quad (1.94) \]

Since the right side of (1.92) is a sinusoid, the forced response has the form

\[ v_f(t) = k_3 \cos 6280t + k_4 \sin 6280t \quad (1.95) \]

Of course, for the derivation of the forced response we could use phasor analysis but we must first derive an expression for the impedance or admittance because the expressions we’ve used earlier are valid for series and parallel circuits only.

The coefficients \(k_3\) and \(k_4\) will be found by substitution of (1.95) into (1.92) and then by equating like terms. Using MATLAB we get:

```matlab
syms t k3 k4; y0=k3*cos(6280*t)+k4*sin(6280*t); y1=diff(y0)
y1 = -6280*k3*sin(6280*t)+6280*k4*cos(6280*t)
y2=diff(y0,2)
y2 = -39438400*k3*cos(6280*t)-39438400*k4*sin(6280*t)
y=y2+2*10^6*y1+2*10^4*y0
y = -37438400*k3*cos(6280*t)-37438400*k4*sin(6280*t)-12560000*k3*sin(6280*t)+12560000*k4*cos(6280*t)
```

Equating like terms with (1.92) we get

\[ (-37438400 \cdot k_3 + 12560000 \cdot k_4) \cos 6280t = -10^6 \cos 6280t \]
\[ (-12560000 \cdot k_3 - 37438400 \cdot k_4) \sin 6280t = 0 \quad (1.96) \]

Simultaneous solution of the equations of (1.96) is done with MATLAB.
syms k3 k4; eq1=-37438400*k3+12560000*k4+10^6;...
eq2=-12560000*k3-37438400*k4+0; y=solve(eq1,eq2)
y =
  k3: [1x1 sym]
  k4: [1x1 sym]
y.k3
ans =
  0.0240
y.k4
ans =
  -0.0081

that is, \( k_3 = 0.024 \) and \( k_4 = -0.008 \). Then, by substitution into (1.95)

\[
v_f(t) = 0.024 \cos 6280t - 0.008 \sin 6280t
\]

(1.97)

The total response is

\[
v_{\text{out}}(t) = v_n(t) + v_f(t) = e^{-1000t}(k_1 \cos 1000t + k_2 \sin 1000t)
+ 0.024 \cos 6280t - 0.008 \sin 6280t
\]

(1.98)

We will use the initial conditions \( v_{c1} = v_{c2} = 0 \) to evaluate \( k_1 \) and \( k_2 \). We observe that \( v_{c2} = v_{\text{out}} \) and at \( t = 0 \) relation (1.98) becomes

\[
v_{\text{out}}(0) = e^0(k_1 \cos 0 + 0) + 0.024 \cos 0 - 0 = 0
\]

or \( k_1 = -0.024 \) and thus (1.98) simplifies to

\[
v_{\text{out}}(t) = e^{-1000t}(-0.024 \cos 1000t + k_2 \sin 1000t)
+ 0.024 \cos 6280t - 0.008 \sin 6280t
\]

(1.99)

To evaluate the constant \( k_2 \), we make use of the initial condition \( v_{c1}(0) = 0 \). We observe that \( v_{c1} = v_i \) and by KCL at node \( v_i \) we have:

\[
\frac{v_i - v_2}{R_3} + C_2 \frac{dv_{\text{out}}}{dt} = 0
\]

or

\[
\frac{v_i - 0}{5 \times 10^4} = -10^{-8} \frac{dv_{\text{out}}}{dt}
\]
or

\[ v_i = -5 \times 10^{-4} \frac{dv_{out}}{dt} \]

and since \( v_C(0) = v_i(0) = 0 \), it follows that

\[ \left. \frac{dv_{out}}{dt} \right|_{t=0} = 0 \quad (1.100) \]

The last step in finding the constant \( k_2 \) is to differentiate (1.99), evaluate it at \( t = 0 \), and equate it with (1.100). This is done with MATLAB as follows:

\[
\begin{align*}
    y_0 &= \exp(-1000t)*(-0.024*\cos(1000t)+k2*\sin(1000t))... \\
        &\quad +0.024*\cos(6280t)-0.008*\sin(6280t); \\
    y_1 &= \text{diff}(y0)
\end{align*}
\]

\[
\begin{align*}
    y_1 &= -1000*\exp(-1000t)*(-3/125*\cos(1000t)+k2*\sin(1000t)) + \exp(-1000t)*(24*\sin(1000t)+1000*k2*\cos(1000t)) - 3768/25*\sin(6280t)-1256/25*\cos(6280t)
\end{align*}
\]

or

\[
\begin{align*}
    \frac{dv_{out}}{dt} &= -1000e^{-1000t}\left(\frac{-3}{125}\cos1000t+k2\sin1000t\right) \\
    &\quad + e^{-1000t}(24\sin1000t+1000k2\cos1000t) \\
    &\quad - \frac{3768}{25}\sin(6280t)-\frac{1256}{25}\cos(6280t)
\end{align*}
\]

and

\[
\left. \frac{dv_{out}}{dt} \right|_{t=0} = -1000\left(\frac{-3}{125}\right) + 1000k_2 - \frac{1256}{25} \quad (1.101)
\]

Simplifying and equating (1.100) with (1.101) we get

\[ 1000k_2 - 26.24 = 0 \]

or

\[ k_2 = 0.026 \]

and by substitution into (1.99),

\[ v_{out}(t) = e^{-1000t}(-0.024\cos1000t + 0.026\sin1000t) \]

\[ + 0.024\cos6280t-0.008\sin6280t \quad (1.102) \]
We use Excel to sketch $v_{out}(t)$. In Column A we enter several values of time $t$ and in Column B $v_{out}(t)$. The plot is shown in Figure 1.20.

![Figure 1.20. Plot for Example 1.7](image-url)
1.9 Summary

- Circuits that contain energy storing devices can be described by integrodifferential equations and upon differentiation can be simplified to differential equations with constant coefficients.

- A second order circuit contains two energy storing devices. Thus, an RLC circuit is a second order circuit.

- The total response is the summation of the natural and forced responses.

- If the differential equation describing a series RLC circuit that is excited by a constant (DC) voltage source is written in terms of the current, the forced response is zero and thus the total response is just the natural response.

- If the differential equation describing a parallel RLC circuit that is excited by a constant (DC) current source is written in terms of the voltage, the forced response is zero and thus the total response is just the natural response.

- If a circuit is excited by a sinusoidal (AC) source, the forced response is never zero.

- The natural response of a second order circuit may be overdamped, critically damped, or underdamped depending on the values of the circuit constants.

- For a series RLC circuit, the roots \( s_1 \) and \( s_2 \) are found from

\[

s_1, s_2 = -\alpha_S \pm \sqrt{\alpha_S^2 - \omega_0^2} = -\alpha_S \pm \beta_S \quad \text{if} \quad \alpha_S^2 > \omega_0^2
\]

or

\[

s_1, s_2 = -\alpha_S \pm \sqrt{\omega_0^2 - \alpha_S^2} = -\alpha_S \pm \omega_{ns} \quad \text{if} \quad \omega_0^2 > \alpha_S^2
\]

where

\[

\alpha_S = \frac{R}{2L} \quad \omega_0 = \frac{1}{\sqrt{LC}} \quad \beta_S = \sqrt{\alpha_S^2 - \omega_0^2} \quad \omega_{ns} = \sqrt{\omega_0^2 - \alpha_S^2}
\]

If \( \alpha_S^2 > \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and unequal. This results in the overdamped natural response and has the form

\[
i_n(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t}
\]

If \( \alpha_S^2 = \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and equal. This results in the critically damped natural response and has the form

\[
i_n(t) = e^{-\alpha_S t}(k_1 + k_2 t)
\]
If \( \omega_0^2 > \alpha_S^2 \), the roots \( s_1 \) and \( s_2 \) are complex conjugates. This is known as the underdamped or oscillatory natural response and has the form

\[
i_n(t) = e^{-\alpha_S t} (k_1 \cos \omega_{nS} t + k_2 \sin \omega_{nS} t) = k_3 e^{-\alpha_S t} (\cos \omega_{nS} t + \varphi)
\]

- For a parallel GLC circuit, the roots \( s_1 \) and \( s_2 \) are found from

\[
s_1, s_2 = -\alpha_p \pm \sqrt{\alpha_p^2 - \omega_0^2} = -\alpha_p \pm \beta_p \quad \text{if} \quad \alpha_p^2 > \omega_0^2
\]

or

\[
s_1, s_2 = -\alpha_p \pm \sqrt{\omega_0^2 - \alpha_p^2} = -\alpha_p \pm \omega_n \quad \text{if} \quad \omega_0^2 > \alpha_p^2
\]

where

\[
\alpha_p = \frac{G}{2C} \quad \omega_0 = \frac{1}{\sqrt{LC}} \quad \beta_p = \sqrt{\omega_p^2 - \omega_0^2} \quad \omega_n = \sqrt{\omega_0^2 - \omega_p^2}
\]

If \( \alpha_p^2 > \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and unequal. This results in the overdamped natural response and has the form

\[
v_n(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t}
\]

If \( \alpha_p^2 = \omega_0^2 \), the roots \( s_1 \) and \( s_2 \) are real, negative, and equal. This results in the critically damped natural response and has the form

\[
v_n(t) = e^{-\alpha_p t} (k_1 + k_2 t)
\]

If \( \omega_0^2 > \alpha_p^2 \), the roots \( s_1 \) and \( s_2 \) are complex conjugates. This results in the underdamped or oscillatory natural response and has the form

\[
v_n(t) = e^{-\alpha_p t} (k_1 \cos \omega_{nP} t + k_2 \sin \omega_{nP} t) = k_3 e^{-\alpha_p t} (\cos \omega_{nP} t + \varphi)
\]

- If a second order circuit is neither series nor parallel, the natural response if found from

\[
y_n = k_1 e^{s_1 t} + k_2 e^{s_2 t}
\]

or

\[
y_n = (k_1 + k_2 t) e^{s_1 t}
\]

or

\[
y_n = e^{-\alpha t} (k_3 \cos \beta t + k_4 \sin \beta t) = e^{-\alpha t} k_5 \cos (\beta t + \varphi)
\]

depending on the roots of the characteristic equation being real and unequal, real and equal, or complex conjugates respectively.
1.10 Exercises

1. For the circuit of Figure 1.21, it is known that $v_c(0^-) = 0$ and $i_L(0^-) = 0$. Compute and sketch $v_c(t)$ and $i_L(t)$ for $t > 0$.

![Figure 1.21. Circuit for Exercise 1](image)

2. For the circuit of Figure 1.22, it is known that $v_c(0^-) = 0$ and $i_L(0^-) = 0$. Compute and sketch $v_c(t)$ and $i_L(t)$ for $t > 0$.

![Figure 1.22. Circuit for Exercise 2](image)

3. In the circuit of Figure 1.23, the switch $S$ has been closed for a very long time and opens at $t = 0$. Compute $v_c(t)$ for $t > 0$.

![Figure 1.23. Circuit for Exercise 3](image)
4. In the circuit of Figure 1.24, the switch $S$ has been closed for a very long time and opens at $t = 0$. Compute $v_c(t)$ for $t > 0$.

![Figure 1.24. Circuit for Exercise 4](image1)

5. In the circuit of Figure 1.25, the switch $S$ has been in position $A$ for closed for a very long time and it is placed in position $B$ at $t = 0$. Find the value of $R$ that will cause the circuit to become critically damped and then compute $v_c(t)$ and $i_L(t)$ for $t > 0$.

![Figure 1.25. Circuit for Exercise 5](image2)

6. In the circuit of Figure 1.26, the switch $S$ has been closed for a very long time and opens at $t = 0$. Compute $v_{AB}(t)$ for $t > 0$.

![Figure 1.26. Circuit for Exercise 6](image3)
1.11 Solutions to Exercises

Dear Reader:

The remaining pages on this chapter contain the solutions to the exercises.

You must, for your benefit, make an honest effort to find the solutions to the exercises without first looking at the solutions that follow. It is recommended that first you go through and work out those you feel that you know. For the exercises that you are uncertain, review this chapter and try again. Refer to the solutions as a last resort and rework those exercises at a later date.

You should follow this practice with the rest of the exercises of this book.
1.

and since \( i = i_C = C \frac{dv_C}{dt} \), the above becomes

\[
RC \frac{dv_C}{dt} + LC \frac{d^2v_C}{dt^2} + v_C = 100
\]

\[
\frac{d^2v_C}{dt^2} + \frac{R}{L} \frac{dv_C}{dt} + \frac{1}{LC} v_C = \frac{100}{LC}
\]

\[
\frac{d^2v_C}{dt^2} + \frac{10}{0.2} \frac{dv_C}{dt} + \frac{1}{0.2 \times 8 \times 10^{-3}} v_C = \frac{100}{0.2 \times 8 \times 10^{-3}}
\]

\[
\frac{d^2v_C}{dt^2} + 50 \frac{dv_C}{dt} + 625 v_C = 62500
\]

From the characteristic equation

\[
s^2 + 50s + 625 = 0
\]

we get \( s_1 = s_2 = -25 \) (critical damping) and \( \alpha_s = R/2L = 25 \)

The total solution is

\[
v_C(t) = v_{cf} + v_{cn} = 100 + e^{-\alpha_s t}(k_1 + k_2t) = 100 + e^{-25t}(k_1 + k_2t) \quad (1)
\]

With the first initial condition \( v_C(0^-) = 0 \) the above expression becomes \( 0 = 100 + e^0(k_1 + 0) \) or \( k_1 = -100 \) and by substitution into (1) we get

\[
v_C(t) = 100 + e^{-25t}(k_2t - 100) \quad (2)
\]

To evaluate \( k_2 \) we make use of the second initial condition \( i_L(0^-) = 0 \) and since \( i_L = i_C \), and \( i = i_C = C(dv_C)/(dt) \), we differentiate (2) using the following MATLAB code:
syms t k2
v0=100+exp(-25*t)*(k2*t-100); v1=diff(v0)

\[ v1 = -25\exp(-25t) \times (k2 \times t - 100) + \exp(-25t) \times k2 \]

Thus,

\[ \frac{dv_C}{dt} = k_2 e^{-25t} - 25e^{-25t}(k_2 t - 100) \]

and

\[ \left. \frac{dv_C}{dt} \right|_{t=0} = k_2 + 2500 \quad (3) \]

Also, \( \frac{dv_C}{dt} = \frac{i_C}{C} = \frac{i_L}{C} \) and at \( t = 0 \)

\[ \left. \frac{dv_C}{dt} \right|_{t=0} = \frac{i_L(0^-)}{C} = 0 \quad (4) \]

From (3) and (4) \( k_2 + 2500 = 0 \) or \( k_2 = -2500 \) and by substitution into (2)

\[ v_C(t) = 100 - e^{-25t}(2500t + 100) \quad (5) \]

We find \( i_L(t) = i_C(t) \) by differentiating (5) and multiplication by \( C \). Using MATLAB we get:

\[ \text{syms t} \]
\[ C=8*10^{-3}; \]
\[ \text{i0=C*(100-exp(-25*t)*(100+2500*t)); iL=diff(i0)} \]

\[ iL = 1/5\exp(-25t) \times (100+2500t) - 20\exp(-25t) \]

Thus,

\[ i_L(t) = i_C(t) = 0.2e^{-25t}(100 + 2500t) - 20e^{-25t} \]

The plots for \( v_C(t) \) and \( i_L(t) \) are shown on the next page.
2.

The general form of the differential equation that describes this circuit is same as in Exercise 1, that is,

$$\frac{d^2v_C}{dt^2} + \frac{R}{L} \frac{dv_C}{dt} + \frac{1}{LC}v_C = \frac{100}{LC} \quad t > 0$$

$$\frac{d^2v_C}{dt^2} + 0.8 \frac{dv_C}{dt} + 9.16v_C = 916$$

From the characteristic equation $s^2 + 0.8s + 9.16 = 0$ and the MATLAB code below
s = [1 0.8 9.16]; roots(s)
ans =
-0.4000 + 3.0000i
-0.4000 - 3.0000i

we find that \( s_1 = -0.4 + j3 \) and \( s_2 = -0.4 - j3 \). Therefore, the total solution is

\[
v_c(t) = v_{cf} + v_{cn} = 100 + ke^{-\alpha_s t} \cos(\omega_{ns} t + \varphi)
\]

where \( \alpha_s = R/2L = 0.4 \)

and

\[
\omega_{ns} = \sqrt{\omega_0^2 - \alpha_s^2} = \sqrt{1/LC - R^2/4L^2} = \sqrt{9.16 - 0.16} = 3
\]

Thus,

\[
v_c(t) = 100 + ke^{-0.4t} \cos(3t + \varphi) \quad (1)
\]

and with the initial condition \( v_c(0^-) = 0 \) we get \( 0 = 100 + k \cos(0 + \varphi) \) or

\[
k \cos \varphi = -100 \quad (2)
\]

To evaluate \( k \) and \( \varphi \) we differentiate (1) with MATLAB and evaluate it at \( t = 0 \).

syms t k phi; v0 = 100 + k*exp(-0.4*t)*cos(3*t+phi); v1 = diff(v0)

\[
v1 = -2/5*k*exp(-2/5*t)*cos(3*t+phi) - 3*k*exp(-2/5*t)*sin(3*t+phi)
\]

Thus,

\[
\frac{dv_c}{dt} = -0.4k e^{-0.4t} \cos(3t + \varphi) - 3k e^{-0.4t} \sin(3t + \varphi)
\]

\[
\left. \frac{dv_c}{dt} \right|_{t=0} = -0.4k \cos \varphi - 3k \sin \varphi
\]

and with (2)

\[
\left. \frac{dv_c}{dt} \right|_{t=0} = 40 - 3k \sin \varphi \quad (3)
\]

Also, \( \frac{dv_c}{dt} = \frac{i_C}{C} = \frac{i_L}{C} \) and at \( t = 0 \)

\[
\left. \frac{dv_c}{dt} \right|_{t=0} = \left. \frac{i_L(0^-)}{C} \right|_{t=0} = 0 \quad (4)
\]
From (3) and (4)

\[ 3 k \sin \varphi = 40 \quad (5) \]

and from (2) and (5)

\[ \frac{3 k \sin \varphi}{k \cos \varphi} = \frac{40}{-100} \]

\[ 3 \tan \varphi = -0.4 \]

\[ \varphi = \tan^{-1}(-0.4/3) = -0.1326 \text{ rad} = -7.6^\circ \]

The value of \( k \) can be found from either (2) or (5). From (2)

\[ k \cos(-0.1236) = -100 \]

\[ k = \frac{-100}{\cos(-0.1236)} = -100.8 \]

and by substitution into (1)

\[ v_C(t) = 100 - 100.8 e^{-0.4t} \cos(3t - 7.6^\circ) \quad (6) \]

Since \( i_L(t) = i_C(t) = C(dv_C/dt) \), we use MATLAB to differentiate (6).

```matlab
syms t; vC=100-100.8*exp(-0.4*t)*cos(3*t-0.1326); C=0.02183; iL=C*diff(vC)
iL = 137529/156250*exp(-2/5*t)*cos(3*t-663/5000)+412587/62500*exp(-2/5*t)*sin(3*t-663/5000)
137529/156250, 412587/62500
ans =
0.8802
ans =
6.6014
i_L(t) = 0.88 e^{-0.4t} \cos(3t - 7.6^\circ) + 6.6 e^{-0.4t} \sin(3t - 7.6^\circ) \]
```

The plots for \( v_C(t) \) and \( i_L(t) \) are shown on the next page.
At the circuit is as shown below.

At this time the inductor behaves as a short and the capacitor as an open. Then,

\[ i_L(0^-) = \frac{100}{100 + 400} = I_0 = 0.2 \text{ A} \]

and this establishes the first initial condition as \( I_0 = 0.2 \text{ A} \). Also,

\[ v_C(0^-) = v_{400\Omega} = 400 \times i_L(0^-) = 400 \times 0.2 = V_0 = 80 \text{ V} \]

and this establishes the first initial condition as \( V_0 = 80 \text{ V} \).

For \( t > 0 \) the circuit is as shown below.
The general form of the differential equation that describes this circuit is same as in Exercise 1, that is,

\[
\frac{d^2v_c}{dt^2} + \frac{R}{L} \frac{dv_c}{dt} + \frac{1}{LC}v_c = \frac{100}{LC} \quad t > 0
\]

\[
\frac{d^2v_c}{dt^2} + 5 \frac{dv_c}{dt} + 6v_c = 600
\]

From the characteristic equation \( s^2 + 5s + 6 = 0 \) we find that \( s_1 = -2 \) and \( s_2 = -3 \) and the total response for the capacitor voltage is

\[
v_c(t) = v_{cf} + v_{cn} = 100 + k_1 e^{s_1 t} + k_2 e^{s_2 t} = 100 + k_1 e^{-2t} + k_2 e^{-3t} \quad (1)
\]

Using the initial condition \( V_0 = 80 \text{ V} \) we get

\[
v_c(0^-) = V_0 = 80 \text{ V} = 100 + k_1 e^0 + k_2 e^0
\]

or

\[
k_1 + k_2 = -20 \quad (2)
\]

Differentiation of (1) and evaluation at \( t = 0 \) yields

\[
\left. \frac{dv_c}{dt} \right|_{t=0} = -2k_1 - 3k_2 \quad (3)
\]

Also, \( \frac{dv_c}{dt} = \frac{i_C}{C} = \frac{i_L}{C} \) and at \( t = 0 \)

\[
\left. \frac{dv_c}{dt} \right|_{t=0} = \frac{i_L(0^-)}{C} = \frac{0.2}{1/120} = 24 \quad (4)
\]

Equating (3) and (4) we get

\[
-2k_1 - 3k_2 = 24 \quad (5)
\]

and simultaneous solution of (2) and (5) yields \( k_1 = -36 \) and \( k_2 = 16 \).
By substitution into (1) we find the total solution as

\[ v_c(t) = v_{cf} + v_{cn} = 100 - 36e^{-2t} + 16e^{-3t} \]

This is the same circuit as in Exercise 3 where the DC voltage source has been replaced by an AC source that is being applied at \( t = 0^\circ \). No initial conditions were given so we will assume that \( i_L(0^-) = 0 \) and \( v_c(0^-) = 0 \). Also, the circuit constants are the same and thus the natural response has the form \( v_{cn} = k_1e^{-2t} + k_2e^{-3t} \).

We will find the forced (steady-state) response using phasor circuit analysis where \( \omega = 1 \), \( j\omega L = j20 \), \( -j/\omega C = -j120 \), and \( 100\cos \theta \Rightarrow 100\angle 0^\circ \). The phasor circuit is shown below.

Using the voltage division expression we get

\[ V_c = \frac{-j120}{100 + j20 - j120} \frac{100\angle 0^\circ}{100\angle 0^\circ} = \frac{-j120}{100 + j100} \frac{120\angle -90^\circ \times 100\angle 0^\circ}{100\sqrt{2}\angle 45^\circ} = 60\sqrt{2}\angle -135^\circ \]

and in the \( t \)-domain \( v_{cf} = 60\sqrt{2}\cos(t - 135^\circ) \). Therefore, the total response is

\[ v_c(t) = 60\sqrt{2}\cos(t - 135^\circ) + k_1e^{-2t} + k_2e^{-3t} \quad (1) \]

Using the initial condition \( v_c(0^-) = 0 \) and (1) we get
and since \( \cos(-135^\circ) = -\sqrt{2}/2 \) the above expression reduces to

\[ k_1 + k_2 = 60 \quad (2) \]

Differentiating (1) we get

\[ \frac{dv_C}{dt} = 60\sqrt{2}\sin(t + 45^\circ) - 2k_1e^{-2t} - 3k_2e^{-3t} \]

and

\[ \left. \frac{dv_C}{dt} \right|_{t=0} = 60\sqrt{2}\sin(45^\circ) - 2k_1 - 3k_2 \quad (3) \]

Also, \( \frac{dv_C}{dt} = \frac{i_C}{C} = \frac{i_L}{C} \) and at \( t = 0 \)

\[ \left. \frac{dv_C}{dt} \right|_{t=0} = \frac{i_L(0^-)}{C} = 0 \quad (4) \]

Equating (3) and (4) we get

\[ 2k_1 + 3k_2 = 60 \quad (5) \]

Simultaneous solution of (2) and (5) yields \( k_1 = 120 \) and \( k_2 = -60 \). Then, by substitution into (1)

\[ v_C(t) = 60\sqrt{2}\cos(t - 135^\circ) + 120e^{-2t} - 60e^{-3t} \]

5.

We must first find the value of \( R \) before we can establish initial conditions for \( i_L(0^-) = 0 \) and \( v_C(0^-) = 0 \). The condition for critical damping is \( \sqrt{\alpha^2 - \omega_0^2} = 0 \) where \( \alpha = G/2C = 1/2RC \)
and \( \omega_0^2 = 1/LC \). Then, \( \alpha_P^2 = \left( \frac{1}{2R' \times 1/12} \right)^2 = \omega_0^2 = \frac{1}{3 \times 1/12} \) where \( R' = R + 2 \ \Omega \). Therefore, \( \left( \frac{12}{2(R+2)} \right)^2 = 4 \), or \( \left( \frac{6}{R+2} \right)^2 = 4 \), or \( (R+2)^2 = 36/4 = 9 \), or \( R+2 = 3 \) and thus \( R = 1 \).

At \( t = 0^- \) the circuit is as shown below.

From the circuit above

\[
v_C(0^-) = v_6 \Omega = \frac{6}{3+1+6} \times 12 = 7.2 \ \text{V}
\]

and

\[
i_L(0^-) = \frac{v_6 \Omega}{6} = \frac{7.2}{6} = 1.2 \ \text{A}
\]

At \( t = 0^+ \) the circuit is as shown below.

Since the circuit is critically damped, the solution has the form

\[
v_C(t) = e^{-\alpha_P t}(k_1 + k_2 t)
\]

where \( \alpha_P = \left( \frac{1}{2(1+2) \times 1/12} \right) = 2 \) and thus

\[
v_C(t) = e^{-2t}(k_1 + k_2 t) \quad (1)
\]
With the initial condition \( v_C(0^-) = 7.2 \text{ V} \) relation (1) becomes \( 7.2 = e^0(k_1 + 0) \) or \( k_1 = 7.2 \text{ V} \) and (1) simplifies to

\[
v_C(t) = e^{-2t}(7.2 + k_2 t) \quad (2)
\]

Differentiating (2) we get

\[
\frac{dv_C}{dt} = k_2 e^{-2t} - 2e^{-2t}(7.2 + k_2 t)
\]

and

\[
\left. \frac{dv_C}{dt} \right|_{t=0} = k_2 - 2(7.2 + 0) = k_2 - 14.4 \quad (3)
\]

Also, \( \frac{dv_C}{dt} = \frac{i_C}{C} \) and at \( t = 0 \)

\[
\left. \frac{dv_C}{dt} \right|_{t=0} = \frac{i_C(0)}{C} = \frac{0}{C} = 0 \quad (4)
\]

because at \( t = 0 \) the capacitor is an open circuit.

Equating (3) and (4) we get \( k_2 - 14.4 = 0 \) or \( k_2 = 14.4 \) and by substitution into (2)

\[
v_C(t) = e^{-2t}(7.2 + 14.4t) = 7.2e^{-2t}(2t + 1)
\]

We find \( i_L(t) \) from \( i_R(t) + i_C(t) + i_L(t) = 0 \) or \( i_L(t) = -i_C(t) - i_R(t) \) where \( i_C(t) = C\frac{dv_C}{dt} \) and \( i_R(t) = v_R(t)/(1 + 2) = v_C(t)/3 \). Then,

\[
i_L(t) = \frac{-1}{12}(-14.4e^{-2t}(2t + 1) + 14.4e^{-2t}) - \frac{7.2}{3}e^{-2t}(2t + 1) = -2.4e^{-2t}(t + 1)
\]

6.

At \( t = 0^- \) the circuit is as shown below where \( i_L(0^-) = 12/2 = 6 \text{ A} \), \( v_C(0^-) = 12 \text{ V} \), and thus the initial conditions have been established.
For \( t > 0 \) the circuit is as shown below.

For this circuit

\[
(R_1 + R_2)i_L + v_C + L \frac{di_L}{dt} = 0
\]

and with \( i_L = i_C = C(dv_C/dt) \) the above relation can be written as

\[
(R_1 + R_2)C \frac{dv_C}{dt} + LC \frac{d^2v_C}{dt^2} + v_C = 0
\]

\[
\frac{d^2v_C}{dt^2} + \frac{(R_1 + R_2)}{L} \frac{dv_C}{dt} + \frac{1}{LC}v_C = 0
\]

\[
\frac{d^2v_C}{dt^2} + \frac{3}{2} \frac{dv_C}{dt} + 2v_C = 0
\]

The characteristic equation of the last expression above yields \( s_1 = -1 \) and \( s_2 = -2 \) and thus

\[
v_C(t) = k_1 e^{-t} + k_2 e^{-2t} \quad (1)
\]

With the initial condition \( v_C(0^-) = 12 \text{ V} \) and (1) we get

\[
k_1 + k_2 = 12 \quad (2)
\]

Differentiating (1) we get

\[
\frac{dv_C}{dt} = -k_1 e^{-t} - 2k_2 e^{-2t}
\]

and

\[
\left. \frac{dv_C}{dt} \right|_{t=0} = -k_1 - 2k_2 \quad (3)
\]

Also, \( \frac{dv_C}{dt} = \frac{i_C}{C} = \frac{i_L}{C} \) and at \( t = 0 \)
\[
\left. \frac{dv_C}{dt} \right|_{t=0} = \frac{i_L(0)}{C} = \frac{6}{1/4} = 24 \quad (4)
\]

From (3) and (4)

\[ -k_1 - 2k_2 = 24 \quad (5) \]

and from (2) and (5) \( k_1 = 48 \) and \( k_2 = -36 \). By substitution into (1) we get

\[ v_C(t) = 48e^{-t} - 36e^{-2t} \]

Then,

\[ v_{AB} = v_L(t) - v_C(t) = L \frac{di_L}{dt} - v_C(t) = LC \frac{d^2i_C}{dt^2} - v_C(t) \]

\[ = 0.5 \left( \frac{d^2}{dt^2}(48e^{-t} - 36e^{-2t}) \right) - 48e^{-t} - 36e^{-2t} \]

\[ = 0.5(48e^{-t} - 144e^{-2t}) - 48e^{-t} - 36e^{-2t} \]

\[ = -24e^{-t} - 108e^{-2t} = -24(e^{-t} + 4.5e^{-2t}) \]
This chapter defines series and parallel resonance. The quality factor $Q$ is then defined in terms of the series and parallel resonant frequencies. The half-power frequencies and bandwidth are also defined in terms of the resonant frequency.

2.1 Series Resonance

Consider phasor series $RLC$ circuit of Figure 2.1.

\[
\text{Impedance } Z = \frac{\text{Phasor Voltage}}{\text{Phasor Current}} = \frac{V_S}{I} = R + j\omega L + \frac{1}{j\omega C} = R + j(\omega L - \frac{1}{\omega C}) \quad (2.1)
\]

or

\[
Z = \sqrt{R^2 + (\omega L - 1/\omega C)^2} \angle \tan^{-1}(\omega L - 1/\omega C)/R \quad (2.2)
\]

Therefore, the magnitude and phase angle of the impedance are:

\[
|Z| = \sqrt{R^2 + (\omega L - 1/\omega C)^2} \quad (2.3)
\]

and

\[
\theta_Z = \tan^{-1}(\omega L - 1/\omega C)/R \quad (2.4)
\]

The components of $|Z|$ are shown on the plot of Figure 2.2.
The frequency at which the capacitive reactance \( X_C = \frac{1}{\omega C} \) and the inductive reactance \( X_L = \omega L \) are equal is called the resonant frequency. The resonant frequency is denoted as \( \omega_0 \) or \( f_0 \) and these can be expressed in terms of the inductance \( L \) and capacitance \( C \) by equating the reactances, that is,

\[
\omega_0 L = \frac{1}{\omega_0 C}
\]

\[
\omega_0^2 = \frac{1}{LC}
\]

\[
\omega_0 = \frac{1}{\sqrt{LC}} \quad \text{(2.5)}
\]

and

\[
f_0 = \frac{1}{2\pi \sqrt{LC}} \quad \text{(2.6)}
\]

We observe that at resonance \( Z_0 = R \) where \( Z_0 \) denotes the impedance value at resonance, and \( \theta_Z = 0 \). In our subsequent discussion the subscript zero will be used to indicate that the circuit variables are at resonance.

**Example 2.1**

For the circuit shown in Figure 2.3, compute \( I_0, \omega_0, C, V_{R0}, |V_{L0}|, \) and \( |V_{C0}| \). Then, draw a phasor diagram showing \( V_{R0}, |V_{L0}|, \) and \( |V_{C0}| \).
Series Resonance

![Circuit Diagram]

Figure 2.3. Circuit for Example 2.1

**Solution:**

At resonance,

\[ jX_L = -jX_C \]

and thus

\[ Z_0 = R = 1.2 \Omega \]

Then,

\[ I_0 = \frac{120 \text{ V}}{1.2 \Omega} = 100 \text{ A} \]

Since

\[ X_{L0} = \omega_0 L = 10 \Omega \]

it follows that

\[ \omega_0 = \frac{10}{L} = \frac{10}{0.2 \times 10^{-3}} = 50000 \text{ rad/s} \]

Therefore,

\[ X_{C0} = X_{L0} = 10 = \frac{1}{\omega_0 C} \]

or

\[ C = \frac{1}{10 \times 50000} = 2 \mu F \]

Now,

\[ V_{R0} = RI_0 = 1.2 \times 100 = 120 \text{ V} \]

\[ |V_{L0}| = \omega_0 L I_0 = 50000 \times 0.2 \times 10^{-3} \times 100 = 1000 \]

and

\[ |V_{C0}| = \frac{1}{\omega_0 C I_0} = \frac{1}{50000 \times 2 \times 10^{-6}} \times 100 = 1000 \text{ V} \]

The phasor diagram showing \( V_{R0} \), \( |V_{L0}| \), and \( |V_{C0}| \) is shown in Figure 2.4.
Figure 2.4. Phasor diagram for Example 2.1

Figure 2.4 reveals that $|V_{L0}| = |V_{C0}| = 1000 \text{ V}$ and these voltages are much higher than the applied voltage of $120 \text{ V}$. This illustrates the useful property of resonant circuits to develop high voltages across capacitors and inductors.

### 2.2 Quality Factor $Q_{0S}$ in Series Resonance

The quality factor * is an important parameter in resonant circuits. Its definition is derived from the following relations:

At resonance,

$$\omega_0L = \frac{I}{\omega_0C}$$

and

$$I_0 = \frac{|V_s|}{R}$$

Then

$$|V_{L0}| = \omega_0L I_0 = \omega_0L \frac{|V_s|}{R} = \frac{\omega_0L}{R} |V_s| \quad (2.7)$$

and

$$|V_{C0}| = \frac{1}{\omega_0C} I_0 = \frac{1}{\omega_0C} \frac{|V_s|}{R} = \frac{1}{\omega_0RC} |V_s| \quad (2.8)$$

At series resonance the left sides of (2.7) and (2.8) are equal and therefore,

$$\frac{\omega_0L}{R} = \frac{1}{\omega_0RC}$$

* We denote the quality factor for series resonant circuits as $Q_{0S}$, and the quality factor for parallel resonant circuits as $Q_{0P}$. 

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Then, by definition

\[
Q_{0S} = \frac{\omega_0 L}{R} = \frac{1}{\omega_0 RC}
\]

\[\text{(2.9)}\]

**Quality Factor at Series Resonance**

In a practical circuit, the resistance \( R \) in the definition of \( Q_{0S} \) above, represents the resistance of the inductor and thus the quality factor \( Q_{0S} \) is a measure of the energy storage property of the inductance \( L \) in relation to the energy dissipation property of the resistance \( R \) of that inductance.

In terms of \( Q_{0S} \), the magnitude of the voltages across the inductor and capacitor are

\[
|V_{L0}| = |V_{C0}| = Q_{0S}|V_S|
\]

\[\text{(2.10)}\]

and therefore, we say that there is a “resonant” rise in the voltage across the reactive devices and it is equal to the \( Q_{0S} \) times the applied voltage. Thus in Example 2.1,

\[
Q_{0S} = \frac{|V_{L0}|}{|V_S|} = \frac{|V_{C0}|}{|V_S|} = \frac{1000}{120} = \frac{25}{3}
\]

The quality factor \( Q \) is also a measure of frequency selectivity. Thus, we say that a circuit with a high \( Q \) has a high selectivity, whereas a low \( Q \) circuit has low selectivity. The high frequency selectivity is more desirable in parallel circuits as we will see in the next section.

Figure 2.5 shows the relative response versus \( \omega \) for \( Q = 25, 50, \text{ and } 100 \) where we observe that highest \( Q \) provides the best frequency selectivity, i.e., higher rejection of signal components outside the bandwidth \( BW = \omega_2 - \omega_1 \) which is the difference in the 3 dB frequencies.

---

**Figure 2.5. Selectivity curves with \( Q = 25, 50, \text{ and } 100 \)**

---

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We will see later that

\[
Q = \frac{\omega_0}{\omega_2 - \omega_1} = \frac{\text{Resonant Frequency}}{\text{Bandwidth}}
\]  

(2.11)

We also observe from (2.9) that selectivity depends on \( R \) and this dependence is shown on the plot of Figure 2.6.

If we keep one reactive device, say \( L \), constant while varying \( C \), the relative response “shifts” as shown in Figure 2.7, but the general shape does not change.

2.3 Parallel Resonance

Parallel resonance (antiresonance) applies to parallel circuits such as that shown in Figure 2.8.
The admittance $Y$ of this circuit is given by

$$\text{Admittance } = Y = \frac{\text{Phasor Current}}{\text{Phasor Voltage}} = \frac{I_S}{V} = G + j\omega C + \frac{1}{j\omega L} = G + j\left(\omega C \frac{1}{\omega L}\right)$$

or

$$Y = \sqrt{G^2 + (\omega C - \frac{1}{\omega L})^2} \angle \tan^{-1}(\omega C - \frac{1}{\omega L}) / G \tag{2.12}$$

Therefore, the magnitude and phase angle of the admittance $Y$ are:

$$|Y| = \sqrt{G^2 + (\omega C - \frac{1}{\omega L})^2} \tag{2.13}$$

and

$$\theta_Y = \tan^{-1}(\frac{\omega C - \frac{1}{\omega L}}{G}) \tag{2.14}$$

The frequency at which the inductive suscetance $B_L = 1/\omega L$ and the capacitive suscetance $B_C = \omega C$ are equal is, again, called the resonant frequency and it is also denoted as $\omega_0$. We can find $\omega_0$ in terms of $L$ and $C$ as before.

Since

$$\omega_0 C \frac{1}{\omega_0 L}$$

then,

$$\omega_0 C = \frac{1}{\omega_0 L}$$

and

$$\omega_0 = \frac{1}{\sqrt{LC}} \tag{2.15}$$

as before. The components of $|Y|$ are shown on the plot of Figure 2.2.
We observe that at this parallel resonant frequency,

\[ Y_0 = G \]  
(2.16)

and

\[ \theta_Y = 0 \]  
(2.17)

**Example 2.2**

For the circuit of Figure 2.10, \( i_S(t) = 10 \cos 5000t \) mA. Compute \( i_G(t) \), \( i_L(t) \), and \( i_C(t) \).

![Circuit diagram](image)

*Figure 2.10. Circuit for Example 2.2*

**Solution:**

The capacitive and inductive susceptances are

\[ B_C = \omega C = 5000 \times 4 \times 10^{-6} = 0.02 \ \Omega^{-1} \]

and

\[ B_L = \frac{1}{\omega L} = \frac{1}{5000 \times 10 \times 10^{-3}} = 0.02 \ \Omega^{-1} \]

and since \( B_L = B_C \), the given circuit operates at parallel resonance with \( \omega_0 = 5000 \ rad/s \). Then,
Quality Factor $Q_{0p}$ in Parallel Resonance

\[ Y_0 = G = 0.01 \, \Omega^{-1} \]

and

\[ i_G(t) = i_S(t) = 10 \cos(5000t) \, mA \]

Next, to compute $i_L(t)$ and $i_C(t)$, we must first find $v_0(t)$. For this example,

\[ v_0(t) = \frac{i_G(t)}{G} = \frac{10 \cos(5000t) \, mA}{0.01 \, \Omega^{-1}} = 1000 \cos(5000t) \, mV = \cos(5000t) \, V \]

In phasor form,

\[ v_0(t) = \cos(5000t) \, V \iff V_0 = 1 \angle 0^\circ \]

Now,

\[ I_{L0} = (-jB_L) V_0 = (1 \angle -90^\circ)(0.02)(1 \angle 0^\circ) = 0.02 \angle -90^\circ \, A \]

and in the $t$-domain,

\[ I_{L0} = 0.02 \angle -90^\circ \, A \iff i_{L0}(t) = 0.02 \cos(5000t - 90^\circ) \, A \]

or

\[ i_{L0}(t) = 20 \sin(5000t) \, mA \]

Similarly,

\[ I_{C0} = jB_C V_0 = (1 \angle 90^\circ)(0.02)(1 \angle 0^\circ) = 0.02 \angle 90^\circ \, A \]

and in the $t$-domain,

\[ I_{C0} = 0.02 \angle 90^\circ \, A \iff i_{C0}(t) = 0.02 \cos(5000t + 90^\circ) \, A \]

or

\[ i_{C0}(t) = -20 \sin(5000t) \, mA \]

We observe that $i_{L0}(t) + i_{C0}(t) = 0$ as expected.

### 2.4 Quality Factor $Q_{0p}$ in Parallel Resonance

At parallel resonance,

\[ \omega_0 C = \frac{I}{\omega_0 L} \]

and

\[ V_0 = \frac{|I_S|}{G} \]

Then,

\[ |I_{C0}| = \omega_0 CV_0 = \omega_0 C \frac{|I_S|}{G} = \frac{\omega_0 C}{G} |I_S| \]

Also,
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\[ |I_{L0}| = \frac{I}{\omega_0 L} V_0 = \frac{I}{\omega_0 G} \left| V_S \right| = \frac{I}{\omega_0 G |I_S|} \]  
\hspace{1cm} (2.19)

At parallel resonance the left sides of (2.18) and (2.19) are equal and therefore,

\[ \frac{\omega_0 C}{G} = \frac{I}{\omega_0 G L} \]

Now, by definition

\[ Q_{0p} = \frac{\omega_0 C}{G} = \frac{I}{\omega_0 G L} \]  
\hspace{1cm} (2.20)

Quality Factor at Parallel Resonance

The above expressions indicate that at parallel resonance, it is possible to develop high currents through the capacitors and inductors. This was found to be true in Example 2.2.

2.5 General Definition of \( Q \)

The general (and best) definition of \( Q \) is

\[ Q = 2\pi \frac{\text{Maximum Energy Stored}}{\text{Energy Dissipated per Cycle}} \]  
\hspace{1cm} (2.21)

Essentially, the resonant frequency is the frequency at which the inductor gives up energy just as fast as the capacitor requires it during one quarter cycle, and absorbs energy just as fast as it is released by the capacitor during the next quarter cycle. This can be seen from Figure 2.11 where at the instant of maximum current the energy is all stored in the inductance, and at the instant of zero current all the energy is stored in the capacitor.

\[ W_L & W_C \text{ in Series RLC Circuit} \]

\[ \text{Energy (J)} \]

\[ W_L \]

\[ W_C \]

\[ v_C \]

\[ i_L \]

\[ \omega \]

Figure 2.11. Waveforms for \( W_L \) and \( W_C \) at resonance
2.6 Energy in L and C at Resonance

For a series RLC circuit we let

\[ i = I_p \cos \omega t = C \frac{dv_C}{dt} \]

Then,

\[ v_C = \frac{I_p}{\omega C} \sin \omega t \]

Also,

\[ W_L = \frac{1}{2} LI_p^2 = \frac{1}{2} LI_p^2 \cos^2 \omega t \quad (2.22) \]

and

\[ W_C = \frac{1}{2} CV_p^2 = \frac{1}{2} \frac{1}{\omega^2 C} \sin^2 \omega t \quad (2.23) \]

Therefore, by (2.22) and (2.23), the total energy \( W_T \) at any instant is

\[ W_T = W_L + W_C = \frac{1}{2} I_p^2 \left[ L \cos^2 \omega t + \frac{1}{\omega^2 C} \sin^2 \omega t \right] \quad (2.24) \]

and this expression is true for any series circuit, that is, the circuit need not be at resonance. However, at resonance,

\[ \omega_0 L = \frac{1}{\omega_0 C} \]

or

\[ L = \frac{1}{\omega_0^2 C} \]

By substitution into (2.24),

\[ W_T = \frac{1}{2} I_p^2 \left[ L \cos^2 \omega_0 t + L \sin^2 \omega_0 t \right] = \frac{1}{2} I_p^2 L = \frac{1}{2} I_p^2 \frac{1}{\omega_0^2 C} \quad (2.25) \]

and (2.25) shows that the total energy \( W_T \) is dependent only on the circuit constants \( L, C \) and resonant frequency, but it is independent of time.

Next, using the general definition of \( Q \) we get:

\[ Q_{os} = 2\pi \frac{\text{Maximum Energy Stored}}{\text{Energy Dissipated per Cycle}} = 2\pi \frac{(1/2)I_p^2 L}{(1/2)I_p^2 R/f_0} = 2\pi \frac{f_0 L}{R} \]
or

\[ Q_{0S} = \frac{\omega_0 L}{R} \]  \hspace{1cm} (2.26)

and we observe that (2.26) is the same as (2.9). Similarly,

\[ Q_{0S} = 2\pi \frac{\text{Maximum Energy Stored}}{\text{Energy Dissipated per Cycle}} = 2\pi \frac{(1/2)I_p^2(1/\omega_0^2C)}{(1/2)I_p^2R/f_0} = 2\pi \frac{f_0}{\omega_0^2RC} \]

or

\[ Q_{0S} = 2\pi \frac{\omega_0}{\omega_0^2RC} = \frac{1}{\omega_0RC} \]  \hspace{1cm} (2.27)

and this is also the same as (2.9).

Following the same procedure for a simple \textit{GLC} (or \textit{RLC}) parallel circuit we can show that:

\[ Q_{0P} = \frac{\omega_0 C}{G} = \frac{1}{\omega_0LG} \]  \hspace{1cm} (2.28)

and this is the same as (2.20).

2.7 Half-Power Frequencies - Bandwidth

Parallel resonance is by far more important and practical than series resonance and therefore, the remaining discussion will be on parallel \textit{GLC} (or \textit{RLC}) circuits.

The plot of Figure 2.12 shows the magnitude of the voltage response versus radian frequency for a typical parallel \textit{RLC} circuit.

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{parallel_rlc_bandwidth.png}
\caption{Relative voltage vs. radian frequency in a parallel RLC circuit}
\end{figure}
By definition, the half-power frequencies $\omega_1$ and $\omega_2$ in Figure 2.12 are the frequencies at which the magnitude of the input admittance of a parallel resonant circuit, is greater than the magnitude at resonance by a factor of $\sqrt{2}$, or equivalently, the frequencies at which the magnitude of the input impedance of a parallel resonant circuit, is less than the magnitude at resonance by a factor of $\sqrt{2}$ as shown above. We observe also, that $\omega_1$ and $\omega_2$ are not exactly equidistant from $\omega_0$. However, it is convenient to assume that they are equidistant, and unless otherwise stated, this assumption will be followed in the subsequent discussion.

We call $\omega_1$ the lower half-power point, and $\omega_2$ the upper half-power point. The difference $\omega_2 - \omega_1$ is the half-power bandwidth $BW$, that is,

$$BW = BW = \omega_2 - \omega_1$$ (2.29)

The names half-power frequencies and half-power bandwidth arise from the fact that the power at these frequencies drop to 0.5 since $(\sqrt{2}/2)^2 = 0.5$.

The bandwidth $BW$ can also be expressed in terms of the quality factor $Q$ as follows:

Consider the admittance

$$Y = G + j\left(\omega C - \frac{1}{\omega L}\right)$$

Multiplying the $j$ term by $G\left(\frac{\omega_0}{\omega_0 G}\right)$, we get

$$Y = G + jG\left(\frac{\omega_0 C}{\omega_0 G} - \frac{\omega_0}{\omega_0 LG}\right)$$

Recalling that for parallel resonance

$$Q_{0p} = \frac{\omega_0 C}{G} = \frac{1}{\omega_0 LG}$$

by substitution we get

$$Y = G\left[1 + jQ_{0p}\left(\frac{\omega}{\omega_0} - \frac{\omega_0}{\omega}\right)\right]$$ (2.30)

and if $\omega = \omega_0$, then

$$Y = G$$
Next, we want to find the bandwidth $\omega_2 - \omega_1$ in terms of the quality factor $Q_{0P}$. At the half-power points, the magnitude of the admittance is $(\sqrt{2}/2)|Y_p|$ and, if we use the half-power points as reference, then to obtain the admittance value of

$$||Y_{max}|| = \sqrt{2}G$$

we must set

$$Q_{0P}\left(\frac{\omega_2}{\omega_0} - \frac{\omega_0}{\omega_2}\right) = 1$$

for $\omega = \omega_2$.

We must also set

$$Q_{0P}\left(\frac{\omega_1}{\omega_0} - \frac{\omega_0}{\omega_1}\right) = -1$$

for $\omega = \omega_1$.

Recalling that $|(1 \pm j)| = \sqrt{2}$ and solving the above expressions for $\omega_1$ and $\omega_2$, we get

$$\omega_2 = \left[\sqrt{1 + \frac{1}{2Q_{0P}^2}} + \frac{1}{2Q_{0P}}\right]$$

(2.31)

and

$$\omega_1 = \left[\sqrt{1 + \frac{1}{2Q_{0P}^2}} - \frac{1}{2Q_{0P}}\right]$$

(2.32)

Subtraction of (2.32) from (2.31) yields

$$BW = \omega_2 - \omega_1 = \frac{\omega_0}{Q_{0P}}$$

(2.33)

or

$$BW = f_2 - f_1 = \frac{f_0}{Q_{0P}}$$

(2.34)

As mentioned earlier, $\omega_1$ and $\omega_2$ are not equidistant from $\omega_0$. In fact, the resonant frequency $\omega_0$ is the geometric mean * of $\omega_1$ and $\omega_2$, that is,

$$\omega_0 = \sqrt{\omega_1 \omega_2}$$

(2.35)

* The geometric mean of n positive numbers $a_1, a_2,..., a_n$ is the nth root of the product $a_1 \cdot a_2 \cdot \ldots a_n$
This can be shown by multiplication of the two expressions in (2.31) and (2.32) and substitution into (2.33).

**Example 2.3**

For the network of Figure 2.13, find:

a. $\omega_0$

b. $Q_{0P}$

c. $BW$

d. $\omega_1$

e. $\omega_2$

![Figure 2.13. Network for Example 2.3](image)

**Solution:**

a. 

$$\omega_0^2 = \frac{1}{LC} = \frac{1}{1 \times 10^{-3} \times 0.4 \times 10^{-6}} = 25 \times 10^8$$

or

$$\omega_0 = 50000 \text{ r/s}$$

b. 

$$Q_{0P} = \frac{\omega_0 C}{G} = \frac{5 \times 10^4 \times 0.4 \times 10^{-6}}{10^{-3}} = 20$$

c. 

$$BW = \frac{\omega_0}{Q_{0P}} = \frac{50000}{20} = 2500 = \text{rad/s}$$

d. 

$$\omega_1 = \omega_0 - \frac{BW}{2} = 50000 - 1250 = 48750 \text{ rad/s}$$

e. 

$$\omega_2 = \omega_0 + \frac{BW}{2} = 50000 + 1250 = 51250 \text{ rad/s}$$
2.8 A Practical Parallel Resonant Circuit

In our previous discussion, we assumed that the inductors are ideal, but a real inductor has some resistance. The circuit shown in Figure 2.14 is a practical parallel resonant circuit.

![Figure 2.14. A practical parallel resonant circuit](image)

To derive an expression for its resonant frequency, we proceed as follows:

The resonant frequency is independent of the conductance $G$ and, for simplicity, it is omitted from the network of Figure 2.14. We will therefore, find an expression for the network of Figure 2.15.

![Figure 2.15. Simplified network for derivation of the resonant frequency](image)

For the network of Figure 2.15,

$$ I_L = \frac{V}{R + j\omega L} = \frac{(R-j\omega L)}{R^2 + (\omega L)^2}V $$

and

$$ I_C = \frac{V}{1/(j\omega C)} = (j\omega C)V $$

where

$$ Re\{I_L\} = \frac{R}{R^2 + (\omega L)^2}V $$

and

$$ Im\{I_L\} = \frac{-\omega L}{R^2 + (\omega L)^2}V $$

Also,

$$ Re\{I_C\} = 0 $$

and

$$ Im\{I_C\} = (\omega C)V $$
Then,

\[ I_T = I_L + I_C = \left[ \text{Re}\{I_L\} + \text{Im}\{I_L\} \right] V + \left[ \text{Re}\{I_C\} + \text{Im}\{I_C\} \right] V \]

\[ = \left[ \text{Re}\{I_L\} + \text{Re}\{I_C\} + \text{Im}\{I_L\} + \text{Im}\{I_C\} \right] V \]

\[ = \left[ \text{Re}\{I_T\} + \text{Im}\{I_T\} \right] V \quad (2.36) \]

Now, at resonance, the imaginary component of \( I_T \) must be zero, that is,

\[ \text{Im}\{I_T\} = \text{Im}\{I_L\} + \text{Im}\{I_C\} = \left( \omega_0 C - \frac{\omega_0 L}{R^2 + (\omega_0 L)^2} \right) V = 0 \]

and solving for \( \omega_0 \) we get

\[ \omega_0 = \sqrt{\frac{1}{LC} - \frac{R^2}{L^2}} \quad (2.37) \]

or

\[ f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{L^2}} \quad (2.38) \]

We observe that for \( R = 0 \), (2.37) reduces to \( \omega_0 = \frac{1}{\sqrt{LC}} \) as before.

2.9 Radio and Television Receivers

When a radio or TV receiver is tuned to a particular station or channel, it is set to operate at the resonant frequency of that station or channel. As we have seen, a parallel circuit has high impedance (low admittance) at its resonant frequency. Therefore, it attenuates signals at all frequencies except the resonant frequency.

We have also seen that one particular inductor and one particular capacitor will resonate to one frequency only. Varying either the inductance or the capacitance of the tuned circuit, will change the resonant frequency. Generally, the inductance is kept constant and the capacitor value is changed as we select different stations or channels.

The block diagram of Figure 2.16 is a typical AM (Amplitude Modulation) radio receiver.
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Figure 2.16. Block diagram of a typical AM radio receiver

The antenna picks up signals from several stations and these are fed into the Radio Frequency (RF) Amplifier which improves the Signal-to-Noise (S/N) ratio. The RF amplifier also serves as a preselector. This preselection suppresses the image-frequency interference as explained below.

When we tune to a station of, say 740 KHz, we are setting the RF circuit to 740 KHz and at the same time the local oscillator is set at 740 KHz + 456 KHz = 1196 KHz. This is accomplished by the capacitor in the RF amplifier which is also ganged to the local oscillator. These two signals, one of 740 KHz and the other of 1196 KHz, are fed into the mixer whose output into the Intermediate Frequency (IF) amplifier is 456 KHz; this is the difference between these two frequencies (1196 KHz−740 KHz = 456 KHz).

The IF amplifier is always set at 456 KHz and therefore if the antenna picks another signal from another station, say 850 KHz, it would be mixed with the local oscillator to produce a frequency of 1196 KHz−850 KHz = 346 KHz but since the IF amplifier is set at 456 KHz, the unwanted 850 KHz signal will not be amplified. Of course, in order to hear the signal at 850 KHz the radio receiver must be retuned to that frequency and the local oscillator frequency will be changed to 850 KHz + 456 KHz = 1306 KHz so that the difference of these frequencies will be again 456 KHz.

Now let us assume that we select a station at 600 KHz. Then, the local oscillator will be set to 600 KHz + 456 KHz = 1056 KHz so that the IF signal will again be 456 KHz. Now, let us suppose that a powerful nearby station broadcasts at 1512 KHz and this signal is picked up by the mixer circuit. The difference between this signal and the local oscillator will also be 456 KHz 1512 KHz−1056 KHz = 456 KHz. The IF amplifier will then amplify both signals and the result will be a strong interference so that the radio speaker will produce unintelligent sounds. This interference is called image-frequency interference and it is reduced by the RF amplifier before entering the mixer circuit and for this reason the RF amplifier is said to act as a preselector.

The function of the detector circuit is to convert the IF signal which contains both the carrier and the desired signal to an audio signal and this signal is amplified by the Audio Frequency (AF) Amplifier whose output appears at the radio speaker.
Example 2.4

A radio receiver with a parallel \(GLC\) circuit whose inductance is \(L = 0.5 \, mH\) is tuned to a radio station transmitting at 810 KHz frequency.

a. What is the value of the capacitor of this circuit at this resonant frequency?

b. What is the value of conductance \(G\) if \(Q_{0p} = 75\)?

c. If a nearby radio station transmits at 740 KHz and both signals picked up by the antenna have the same current amplitude \(I\) (\(\mu A\)), what is the ratio of the voltage at 810 KHz to the voltage at 740 KHz?

Solution:

a. 

\[
\omega_0^2 = \frac{1}{LC}
\]

or

\[
f_0^2 = \frac{1}{4\pi^2 LC}
\]

Then,

\[
C = \frac{1}{4\pi^2 0.5 \times 10^{-3} \times (810 \times 10^3)^2} = 77.2 \, pF
\]

b. 

\[
Q_{0p} = \frac{\omega_0 C}{G}
\]

or

\[
G = \frac{2\pi f_0 C}{Q_{0p}} = \frac{2\pi \times 8.1 \times 10^5 \times 77.2 \times 10^{-12}}{75} = 5.4 \, \mu\Omega^{-1}
\]

c. 

\[
|V_{810 \, KHz}| = \frac{I}{|Y_{810 \, KHz}|} = \frac{I}{Y_0} = \frac{I}{G} = \frac{I}{5.24 \times 10^{-6}}
\]  \hspace{1cm} (2.39)

Also,

\[
|V_{740 \, KHz}| = \frac{I}{|Y_{740 \, KHz}|}
\]

where

\[
|Y_{740 \, KHz}| = \sqrt{G^2 + (\omega C - \frac{I}{\omega L})^2}
\]
or
\[ |Y_{740 \text{ KHz}}| = \left(5.24 \times 10^{-6}\right)^2 + \left(2\pi \times 740 \times 10^3 \times 77.2 \times 10^{-12} - \frac{1}{2\pi \times 740 \times 10^3 \times 0.5 \times 10^{-3}}\right)^2 \]

or
\[ |Y_{740 \text{ KHz}}| = 71.2 \ \mu\Omega^{-1} \]

and
\[ |V_{740 \text{ KHz}}| = \frac{I}{71.2 \times 10^{-6}} \quad (2.40) \]

Then from (2.39) and (2.40),
\[ \frac{|V_{810 \text{ KHz}}|}{|V_{740 \text{ KHz}}|} = \frac{I/5.24 \times 10^{-6}}{1/71.2 \times 10^{-6}} = \frac{71.2 \times 10^{-6}}{5.24 \times 10^{-6}} = 13.6 \quad (2.41) \]

that is, the voltage developed across the parallel circuit when it is tuned at \( f = 810 \text{ KHz} \) is 13.6 times larger than the voltage developed at \( f = 740 \text{ KHz} \).

2.10 Summary

- In a series \( RLC \) circuit, the frequency at which the capacitive reactance \( X_C = 1/\omega C \) and the inductive reactance \( X_L = \omega L \) are equal, is called the resonant frequency.

- The resonant frequency is denoted as \( \omega_0 \) or \( f_0 \) where

\[ \omega_0 = \frac{1}{\sqrt{LC}} \]

and

\[ f_0 = \frac{1}{2\pi\sqrt{LC}} \]

- The quality factor \( Q_{0S} \) at series resonance is defined as

\[ Q_{0S} = \frac{\omega_0 L}{R} = \frac{1}{\omega_0 RC} \]

- In a parallel \( GLC \) circuit, the frequency at which the inductive susceptance \( B_L = 1/\omega L \) and the capacitive susceptance \( B_C = \omega C \) are equal is, again, called the resonant frequency and it is also denoted as \( \omega_0 \). As in a series \( RLC \) circuit, the resonant frequency is
The quality factor $Q_{0P}$ at parallel resonance is defined as

$$Q_{0P} = \frac{\omega_0 C}{G} = \frac{1}{\omega_0 GL}$$

The general definition of $Q$ is

$$Q = 2\pi \frac{\text{Maximum Energy Stored}}{\text{Energy Dissipated per Cycle}}$$

In a parallel $RLC$ circuit, the half-power frequencies $\omega_1$ and $\omega_2$ are the frequencies at which the magnitude of the input admittance of a parallel resonant circuit, is greater than the magnitude at resonance by a factor of $\sqrt{2}$, or equivalently, the frequencies at which the magnitude of the input impedance of a parallel resonant circuit, is less than the magnitude at resonance by a factor of $\sqrt{2}$.

We call $\omega_1$ the lower half-power point, and $\omega_2$ the upper half-power point. The difference $\omega_2 - \omega_1$ is the half-power bandwidth $BW$, that is,

$$\text{Bandwidth} = BW = \omega_2 - \omega_1$$

The bandwidth $BW$ can also be expressed in terms of the quality factor $Q$ as

$$BW = \omega_2 - \omega_1 = \frac{\omega_0}{Q_{0P}}$$

or

$$BW = f_2 - f_1 = \frac{f_0}{Q_{0P}}$$
2.11 Exercises

1. A series RLC circuit is resonant at \( f_0 = 1 \text{ MHz} \) with \( Z_0 = 100 \Omega \) and its half-power bandwidth is \( BW = 20 \text{ KHz} \). Find \( R \), \( L \), and \( C \) for this circuit.

2. For the network of Figure 2.17, the impedance \( Z_1 \) is variable, \( Z_2 = 3 + j4 \) and \( Z_3 = 4 - j3 \). To what value should \( Z_1 \) be adjusted so that the network will operate at resonant frequency?

![Figure 2.17. Network for Exercise 2](image)

3. For the circuit of Figure 2.18 with the capacitance \( C \) adjusted to 1 \( \mu \text{F} \), the half-power frequencies are \( f_1 = 925 \text{ KHz} \) and \( f_2 = 1075 \text{ KHz} \).
   a. Compute the approximate resonant frequency.
   b. Compute the exact resonant frequency.
   c. Using the approximate value of the resonant frequency, compute the values of \( Q_{op} \), \( G \), and \( L \).

![Figure 2.18. Circuit for Exercise 3](image)

4. The GLC circuit of Figure 2.19, is resonant at \( f_0 = 500 \text{ KHz} \) with \( V_0 = 20 \text{ V} \) and its half-power bandwidth is \( BW = 20 \text{ KHz} \).
   a. Compute \( L \), \( C \), and \( I_0 \) for this circuit.
   b. Compute the magnitude of the admittances \( |Y_1| \) and \( |Y_2| \) corresponding to the half-power frequencies \( f_1 \) and \( f_2 \). Use MATLAB to plot \( |Y| \) in the \( 100 \text{ KHz} \leq f \leq 1000 \text{ KHz} \) range.

![Figure 2.19. Circuit for Exercise 4](image)
5. For the circuit of Figure 2.20, \( v_s = 170 \cos \omega t \) and \( Q_0 = 50 \). Find:
   
   a. \( \omega_0 \)
   
   b. \( BW \)
   
   c. \( \omega_1 \) and \( \omega_2 \)
   
   d. \( |V_{C0}| \)

![Figure 2.20. Circuit for Exercise 5](image)

6. The series-parallel circuit of Figure 2.21, will behave as a filter if the parallel part is made resonant to the frequency we want to suppress, and the series part is made resonant to the frequency we wish to pass. Accordingly, we can adjust capacitor \( C_2 \) to achieve parallel resonance which will reject the unwanted frequency by limiting the current through the resistive load to its minimum value. Afterwards, we can adjust \( C_1 \) to make the entire circuit series resonant at the desired frequency thus making the total impedance minimum so that maximum current will flow into the load.

For this circuit, we want to set the values of capacitors so that \( v_{LOAD} \) will be maximum at \( f_1 = 10 \, KHz \) and minimum at \( f_2 = 43 \, KHz \). Compute the values of \( C_1 \) and \( C_2 \) that will achieve these values. It is suggested that you use MATLAB to plot \( |v_{LOAD}| \) versus frequency \( f \) in the interval \( 1 \, KHz \leq f \leq 100 \, KHz \) to verify your answers.

![Figure 2.21. Circuit for Exercise 6](image)
Chapter 2  Resonance

2.12 Solutions to Exercises

1. At series resonance $Z_0 = R = 100$ and thus $R = 100 \, \Omega$. We find $L$ from $Q_{0S} = \frac{\omega_0}{L/R}$ where $\omega_0 = 2\pi f_0$. Also,

$$Q_{0S} = \frac{\omega_0}{\omega_2 - \omega_1} = \frac{\omega_0}{BW} = \frac{2\pi \times 10^6}{2\pi \times 20 \times 10^3} = 50$$

Then,

$$L = \frac{R \cdot Q_{0S}}{\omega_0} = \frac{100 \times 50}{2\pi \times 10^6} = 0.796 \, mH$$

and from $\omega_0^2 = 1/LC$

$$C = \frac{1}{\omega_0^2 L} = \frac{1}{(2\pi \times 10^6)^2 \times 7.96 \times 10^4} = 31.8 \, pF$$

Check with MATLAB:

```matlab
f0=10^6; w0=2*pi*f0; Z0=100; BW=2*pi*20000; w1=w0-BW/2; w2=w0+BW/2;...
R=Z0; Qos=w0/BW; L=R*Qos/w0; C=1/(w0^2*L); fprintf('R = %5.2f Ohms 	', R); fprintf('L = %5.2e H 	', L);...
fprintf('C = %5.2e F 	', C); fprintf('
');
R = 100.00 Ohms   L = 7.96e-004 H    C = 3.18e-011 F
```

2.

![Network Diagram](image)

$$Z_{IN} = Z_1 + Z_2 \parallel Z_3$$

where

$$Z_2 \parallel Z_3 = \frac{(3+j4) \cdot (4-j3)}{3+j4+4-j3} = \frac{12-j9+j16+12}{7+j} \cdot \frac{7-j}{7-j}$$

$$= \frac{168+j49-j24+7}{7^2+1^2} = \frac{175+j25}{50} = 3.5+j0.5$$

We let $Z_{IN} = R_{IN} + jX_{IN}$ and $Z_1 = R_1 + jX_1$. For resonance we must have

$$Z_{IN} = R_{IN} + jX_{IN} = R_1 + jX_1 + 3.5 + j0.5 = R_{IN} + 0 = R_1 + jX_1 + 3.5 + j0.5$$
Equating real and imaginary parts we get

\[ R_{IN} = R_1 + 3.5 \]
\[ 0 = jX_1 + j0.5 \]

and while \( R_1 \) can be any real number, we must have \( jX_1 = -j0.5 \) and thus

\[ Z_1 = R_1 - j0.5 \ \Omega \]

3.

a. \( BW = f_2 - f_1 = 1075 - 925 = 150 \ \text{KHz} \). Then,

\[ f_0 = f_1 + BW/2 = 925 + 150/2 = 1000 \ \text{KHz} \]

b. The exact value of \( f_0 \) is the geometric mean of \( f_1 \) and \( f_2 \) and thus

\[ f_0 = \sqrt[3]{f_1 \cdot f_2} = \sqrt[3]{(925 + 1075) \times 10^3} = 997.18 \ \text{KHz} \]

c. \( Q_{OP} = \frac{f_0}{f_2 - f_1} = \frac{1000}{150} = 20/3 \). Also, \( Q_{OP} = \frac{\omega_0 C}{G} \). Then

\[ G = \frac{\omega_0 C}{Q_{OP}} = \frac{2\pi f_0 C}{Q_{OP}} = \frac{2\pi \times 10^6 \times 10^{-6}}{20/3} = \frac{3\pi}{10} = 0.94 \ \Omega^{-1} \]

and

\[ L = \frac{1}{\omega_0 C} = \frac{1}{4\pi^2 f_0^2 C} = \frac{1}{4\pi^2 \times 10^{12} \times 10^{-6}} = 0.025 \ \mu\text{H} \]

4.

a. \( Q_{OP} = \frac{f_0}{BW} = \frac{500}{20} = 25 \). Also, \( Q_{OP} = \frac{\omega_0 C}{G} \) or

\[ C = \frac{Q_{OP} \cdot G}{\omega_0} = \frac{25 \times 10^{-3}}{2\pi \times 5 \times 10^5} = 7.96 \times 10^{-9} \ \text{F} = 7.96 \ \text{nF} \]

\[ L = \frac{1}{\omega_0 C} = \frac{1}{4\pi^2 f_0^2 C} = \frac{1}{4\pi^2 \times 25 \times 10^{10} \times 7.96 \times 10^{-9}} = 12.73 \times 10^{-6} \ \text{H} = 12.73 \ \mu\text{H} \]

\[ I_0 = V_0 Y_0 = V_0 G = 20 \times 10^{-3} \ \text{A} = 20 \ \text{mA} \]

b. \( f_1 = f_0 - BW/2 = 500 - 10 = 490 \ \text{KHz} \) and \( f_2 = f_0 + BW/2 = 500 + 10 = 510 \ \text{KHz} \)
\[ Y_{f_1} = G + j \left( \omega_1 C - \frac{1}{\omega_1 L} \right) \]
\[ = 10^{-3} + j \left( 2\pi \times 490 \times 10^3 \times 7.96 \times 10^{-9} - \frac{1}{2\pi \times 490 \times 10^3 \times 12.73 \times 10^{-6}} \right) \]

Likewise,
\[ Y_{f_2} = G + j \left( \omega_1 C - \frac{1}{\omega_1 L} \right) \]
\[ = 10^{-3} + j \left( 2\pi \times 510 \times 10^3 \times 7.96 \times 10^{-9} - \frac{1}{2\pi \times 510 \times 10^3 \times 12.73 \times 10^{-6}} \right) \]

We will use MATLAB to do the computations.

\[ G=10^{-3}; \quad BC1=2\pi \times 490 \times 10^3 \times 7.96 \times 10^{-9} \quad (-9); \]
\[ BL1=1/(2\pi \times 490 \times 10^3 \times 12.73 \times 10^{-6}); \quad Y1=G+j*(BC1-BL1); \]
\[ BC2=2\pi \times 510 \times 10^3 \times 7.96 \times 10^{-9}; \quad BL2=1/(2\pi \times 510 \times 10^3 \times 12.73 \times 10^{-6}); \]
\[ Y2=G+j*(BC2-BL2); \quad fprintf('magY1 = %.2e mho \
', abs(Y1)); \]
\[ fprintf('magY2 = %.2e mho \
', abs(Y2)); \quad fprintf(' 
'); \]
\[ magY1 = 1.42e-003 \text{ mho} \quad \text{magY2} = 1.41e-003 \text{ mho} \]

We will use the following MATLAB code for the plot

\[ f=100 \times 10^3: 10 \times 10^3: 1000 \times 10^3; \quad w=2\pi \times f; \]
\[ G=10^{-3}; \quad C=7.96 \times 10^{-9}; \quad L=12.73 \times 10^{-6}; \]
\[ BC=w.*C; \quad BL=1./(w.*L); \quad Y=G+j*(BC-BL); \quad \text{plot(f,abs(Y))}; \]
\[ \text{xlabel('Frequency in Hz'); ylabel('Magnitude of Admittance'); grid} \]

The plot is shown below.
a. It is important to remember that the relation \( \omega_0 = 1/\sqrt{LC} \) applies only to series RLC and parallel GLC circuits. For any other circuit we must find the input impedance \( Z_{IN} \), set the imaginary part of \( Z \) equal to zero, and solve for \( \omega_0 \). Thus, for the given circuit

\[
Z_{IN} = R_1 + \frac{1}{j\omega C} \parallel (R_2 + j\omega L) = 1 + \frac{1/j\omega C \cdot (10 + j\omega L)}{10 + j(\omega L - 1/\omega C)}
\]

\[
= \frac{10 + j(\omega L - 1/\omega C) + 10/j\omega C + L/C}{10 + j(\omega L - 1/\omega C)} \cdot \frac{10 - j(\omega L - 1/\omega C)}{10 - j(\omega L - 1/\omega C)}
\]

\[
= \frac{100 + j10(\omega L - 1/\omega C) + 100/(j\omega C) + 10L/C - j10(\omega L - 1/\omega C)}{100 + (\omega L - 1/\omega C)^2}
\]

\[
\quad + \frac{(\omega L - 1/\omega C)^2 - (10/\omega C)(\omega L - 1/\omega C) - jL/C(\omega L - 1/\omega C)}{100 + (\omega L - 1/\omega C)^2}
\]

\[
= \frac{100 + 10L/C + (\omega L - 1/\omega C)^2 - (10/\omega C)(\omega L - 1/\omega C)}{100 + (\omega L - 1/\omega C)^2}
\]

\[
\quad + \frac{100/(j\omega C) - jL/C(\omega L - 1/\omega C)}{100 + (\omega L - 1/\omega C)^2}
\]

For resonance, the imaginary part of \( Z_{IN} \) must be zero, that is,

\[
\frac{100}{j\omega_0 C} - \frac{jL}{C}(\omega_0 L - \frac{1}{\omega_0 C}) = 0
\]

\[
-\frac{j}{C}\left[\frac{100}{\omega_0} + L(\omega_0 L - \frac{1}{\omega_0 C})\right] = 0
\]

\[
\frac{100}{\omega_0} + \omega_0 L^2 - \frac{L}{\omega_0 C} = 0
\]

\[
L^2 C \omega_0^2 + 100C - L = 0
\]
Chapter 2  Resonance

\[ \omega_0^2 = \frac{1}{LC} - \frac{100}{L_2^2} = \frac{1}{10^{-3} \times 10^{-6}} - \frac{100}{10^{-6}} = 10^9 - 10^8 = 9 \times 10^8 \]

and thus

\[ \omega_0 = \sqrt{9 \times 10^8} = 30,000 \text{ r/s} \]

b. \[ BW = \frac{\omega_0}{Q} = \frac{30,000}{50} = 600 \text{ r/s} \]

c. \[ \omega_1 = \omega_0 - \frac{BW}{2} = 30,000 - 300 = 29,700 \text{ r/s} \]
\[ \omega_2 = \omega_0 + \frac{BW}{2} = 30,000 + 300 = 30,300 \text{ r/s} \]

d. At resonance, \( j\omega_0 L = j3 \times 10^4 \times 10^{-3} = j30 \ \Omega \) and \( \frac{1}{j\omega_0 C} = -j10^{-4} \times 10^6 / 3 = -j100/3 \).

The phasor equivalent circuit is shown below.

We let \( z_1 = 1 \ \Omega \), \( z_2 = -j100/3 \ \Omega \), and \( z_3 = 10 + j30 \ \Omega \). Using nodal analysis we get:

\[ \frac{V_{C0} - V_S}{z_1} + \frac{V_{C0}}{z_2} + \frac{V_{C0}}{z_3} = 0 \]
\[ \left( \frac{1}{z_1} + \frac{1}{z_2} + \frac{1}{z_3} \right)V_{C0} = \frac{V_S}{z_1} \]

We will use MATLAB to obtain the value of \( V_{C0} \).

Vs=170; z1=1; z2=-j*100/3; z3=10+j*30; Z=1/z1+1/z2+1/z3; Vc0=Vs/Z;...
fprintf('
'); fprintf('Vc0 = %6.2f', abs(Vc0)); fprintf('
'); fprintf('
')
Vc0 = 168.32

6. First, we will find the appropriate value of \( C_2 \). We recall that at parallel resonance the voltage is maximum and the current is minimum. For this circuit the parallel resonance was found as in (2.37), that is,
\[ \omega_0 = \frac{1}{\sqrt{LC - \frac{R^2}{L^2}}} \]

or

\[ 2\pi \times 43,000 = \sqrt{\frac{1}{2 \times 10^{-3} C_2} - \frac{10^4}{4 \times 10^{-6}}} \]

\[ \frac{10^3}{2C_2} = \frac{10^4}{4 \times 10^{-6}} + (2\pi \times 4.3 \times 10^4)^2 = \frac{10^4 + (2\pi \times 4.3 \times 10^4)^2 \times 4 \times 10^{-6}}{4 \times 10^{-6}} \]

\[ C_2 = 500 \left[ \frac{4 \times 10^{-6}}{10^4 + (2\pi \times 4.3 \times 10^4)^2 \times 4 \times 10^{-6}} \right] = 6.62 \times 10^{-9} F = 6.62 \text{ nF} \]

Next, we must find the value of \( C_1 \) that will make the entire circuit series resonant (minimum impedance, maximum current) at \( f = 10 \text{ KHz} \). In the circuit below we let \( z_1 = -jX_{C1} \), \( z_2 = -jX_{C2} \), \( z_3 = R_1 + jX_L \), and \( z_{LOAD} = 1 \).

\[ \begin{align*}
Z_{IN} &= z_1 + \frac{z_2}{z_3} + z_{LOAD} \\
Z_{IN}(f = 10 \text{ KHz}) &= z_1 + \frac{z_2}{z_3}vert_{f=10 \text{ KHz}} + z_{LOAD} = z_1 + \frac{z_2}{z_3}vert_{f=10 \text{ KHz}} + 1 \quad (1)
\end{align*} \]

where \( \frac{z_2}{z_3}vert_{f=10 \text{ KHz}} \) is found with the MATLAB code below.

```
format short g; f=10000; w=2*pi*f; C2=6.62*10^(-9); XC2=1/(w*C2); L=2*10^(-3);...XL=w*L; R1=100; z2=-j*Xc2; z3=R1+j*XL; Zp=z2*z3/(z2+z3)
Zp =
111.12 + 127.72i
```

and by substitution into (1)
The expression of (2) will be minimum if we let $z_l = -j127.72 \, \Omega$ at $f = 10 \, KHz$. Then, the capacitor $C_l$ value must be such that $1/\omega C = 127.72$ or

$$C_l = \frac{1}{2\pi \times 10^4 \times 127.72} = 1.25 \times 10^{-7} \, F = 0.125 \, \mu F$$

Shown below is the plot of $|V_{LOAD}|$ versus frequency and the MATLAB code that produces this plot.

```
f=1000:100:60000; w=2*pi*f; Vs=170; C1=1.25*10^(-7); C2=6.62*10^(-9);...
L=2.*10.^(3);...
R1=100; Rload=1; z1=j./(w.*C1); z2=j./(w.*C2); z3=R1+j.*w.*L; Zload=Rload;...
Zin=z1+z2.*z3./(z2+z3); Vload=Zload.*Vs./(Zin+Zload); magVload=abs(Vload);...
plot(f,magVload); axis([1000 60000 0 2]);
xlabel('Frequency f'); ylabel('|Vload|'); grid
```

This circuit is considered to be a special type of filter that allows a specific frequency (not a band of frequencies) to pass, and attenuates another specific frequency.
This chapter begins with a discussion of elementary signals that may be applied to electric networks. The unit step, unit ramp, and delta functions are then introduced. The sampling and sifting properties of the delta function are defined and derived. Several examples for expressing a variety of waveforms in terms of these elementary signals are provided.

### 3.1 Signals Described in Math Form

Consider the network of Figure 3.1 where the switch is closed at time $t = 0$.

![Figure 3.1. A switched network with open terminals.](image)

We wish to describe $v_{out}$ in a math form for the time interval $-\infty < t < +\infty$. To do this, it is convenient to divide the time interval into two parts, $-\infty < t < 0$, and $0 < t < \infty$.

For the time interval $-\infty < t < 0$, the switch is open and therefore, the output voltage $v_{out}$ is zero. In other words,

$$v_{out} = 0 \text{ for } -\infty < t < 0 \quad (3.1)$$

For the time interval $0 < t < \infty$, the switch is closed. Then, the input voltage $v_S$ appears at the output, i.e.,

$$v_{out} = v_S \text{ for } 0 < t < \infty \quad (3.2)$$

Combining (3.1) and (3.2) into a single relationship, we get

$$v_{out} = \begin{cases} 0 & -\infty < t < 0 \\ v_S & 0 < t < \infty \end{cases} \quad (3.3)$$

We can express (3.3) by the waveform shown in Figure 3.2.
Chapter 3  Elementary Signals

Figure 3.2. Waveform for \( v_{out} \) as defined in relation (3.3)

The waveform of Figure 3.2 is an example of a discontinuous function. A function is said to be discontinuous if it exhibits points of discontinuity, that is, the function jumps from one value to another without taking on any intermediate values.

3.2 The Unit Step Function \( u_0(t) \)

A well-known discontinuous function is the unit step function \( u_0(t) \) that is defined as

\[
  u_0(t) = \begin{cases} 
  0 & t < 0 \\
  1 & t > 0 
\end{cases} 
\]  \hspace{1cm} (3.4)

It is also represented by the waveform of Figure 3.3.

![Figure 3.3. Waveform for \( u_0(t) \)](image)

In the waveform of Figure 3.3, the unit step function \( u_0(t) \) changes abruptly from 0 to 1 at \( t = 0 \). But if it changes at \( t = t_0 \) instead, it is denoted as \( u_0(t-t_0) \). Its waveform and definition are as shown in Figure 3.4 and relation (3.5).

![Figure 3.4. Waveform for \( u_0(t-t_0) \)](image)

* In some books, the unit step function is denoted as \( u(t) \), that is, without the subscript 0. In this text, however, we will reserve the \( u(t) \) designation for any input.
The Unit Step Function

\[ u_0(t - t_0) = \begin{cases} 
0 & t < t_0 \\
1 & t > t_0 
\end{cases} \]  

(3.5)

If the unit step function changes abruptly from 0 to 1 at \( t = -t_0 \), it is denoted as \( u_0(t + t_0) \). Its waveform and definition are as shown in Figure 3.5 and relation (3.6).

\[ u_0(t + t_0) = \begin{cases} 
0 & t < -t_0 \\
1 & t > -t_0 
\end{cases} \]  

(3.6)

Example 3.1

Consider the network of Figure 3.6, where the switch is closed at time \( t = T \).

Express the output voltage \( v_{out} \) as a function of the unit step function, and sketch the appropriate waveform.

Solution:

For this example, the output voltage \( v_{out} = 0 \) for \( t < T \), and \( v_{out} = v_S \) for \( t > T \). Therefore,

\[ v_{out} = v_S u_0(t - T) \]  

(3.7)

and the waveform is shown in Figure 3.7.
Other forms of the unit step function are shown in Figure 3.8.

Unit step functions can be used to represent other time-varying functions such as the rectangular pulse shown in Figure 3.9.

Thus, the pulse of Figure 3.9(a) is the sum of the unit step functions of Figures 3.9(b) and 3.9(c) is represented as $u_0(t) - u_0(t-1)$. 

$\textbf{Figure 3.7. Waveform for Example 3.1}$

$\textbf{Figure 3.8. Other forms of the unit step function}$

$\textbf{Figure 3.9. A rectangular pulse expressed as the sum of two unit step functions}$
The unit step function offers a convenient method of describing the sudden application of a voltage or current source. For example, a constant voltage source of 24 V applied at \( t = 0 \), can be denoted as \( 24u_0(t) \) V. Likewise, a sinusoidal voltage source \( v(t) = V_m\cos\omega t \) V that is applied to a circuit at \( t = t_0 \), can be described as \( v(t) = (V_m\cos\omega t)u_0(t - t_0) \) V. Also, if the excitation in a circuit is a rectangular, or triangular, or sawtooth, or any other recurring pulse, it can be represented as a sum (difference) of unit step functions.

### Example 3.2

Express the square waveform of Figure 3.10 as a sum of unit step functions. The vertical dotted lines indicate the discontinuities at \( T, 2T, 3T \) and so on.

![Figure 3.10. Square waveform for Example 3.2](image)

#### Solution:

- Line segment ① has height \( A \), starts at \( t = 0 \), and terminates at \( t = T \). Then, as in Example 3.1, this segment is expressed as
  \[
  v_1(t) = A[u_0(t) - u_0(t - T)] \tag{3.8}
  \]

- Line segment ② has height \( -A \), starts at \( t = T \) and terminates at \( t = 2T \). This segment is expressed as
  \[
  v_2(t) = -A[u_0(t - T) - u_0(t - 2T)] \tag{3.9}
  \]

- Line segment ③ has height \( A \), starts at \( t = 2T \) and terminates at \( t = 3T \). This segment is expressed as
  \[
  v_3(t) = A[u_0(t - 2T) - u_0(t - 3T)] \tag{3.10}
  \]

- Line segment ④ has height \( -A \), starts at \( t = 3T \), and terminates at \( t = 4T \). It is expressed as
  \[
  v_4(t) = -A[u_0(t - 3T) - u_0(t - 4T)] \tag{3.11}
  \]

Thus, the square waveform of Figure 3.10 can be expressed as the summation of (3.8) through (3.11), that is,
Chapter 3  Elementary Signals

\[ v(t) = v_1(t) + v_2(t) + v_3(t) + v_4(t) \]
\[ = A[u_0(t) - u_0(t-T)] - A[u_0(t-T) - u_0(t-2T)] + A[u_0(t-2T) - u_0(t-3T)] - A[u_0(t-3T) - u_0(t-4T)] \]  
(3.12)

Combining like terms, we get
\[ v(t) = A[u_0(t) - 2u_0(t-T) + 2u_0(t-2T) - 2u_0(t-3T) + \ldots] \]  
(3.13)

Example 3.3

Express the symmetric rectangular pulse of Figure 3.11 as a sum of unit step functions.

\[ i(t) = \begin{cases} 
A & \text{for } t = 0 \\
0 & \text{for } t < 0 \text{ or } t > T/2 
\end{cases} \]

\[ i(t) = A[u_0(t + T/2) - u_0(t - T/2)] = A[u_0(t + T/2)] - u_0(t - T/2) \]  
(3.14)

Example 3.4

Express the symmetric triangular waveform of Figure 3.12 as a sum of unit step functions.

\[ v(t) = \begin{cases} 
1 & \text{for } -T/2 < t < 0 \\
0 & \text{for } t = 0 \\
1 & \text{for } 0 < t < T/2 
\end{cases} \]

\[ v(t) = A[u_0(t + T/2) - u_0(t - T/2)] - A[u_0(t - T/2) - u_0(t - 3T/2)] + A[u_0(t - 3T/2) - u_0(t - 4T)] \]

Solution:

We first derive the equations for the linear segments 1 and 2 shown in Figure 3.13.
The Unit Step Function

Figure 3.13. Equations for the linear segments of Figure 3.12

For line segment 1,

\[ v_1(t) = \left( \frac{2}{T} t + l \right) \left[ u_0 \left( t + \frac{T}{2} \right) - u_0(t) \right] \]  

(3.15)

and for line segment 2,

\[ v_2(t) = \left( -\frac{2}{T} t + l \right) \left[ u_0(t) - u_0 \left( t - \frac{T}{2} \right) \right] \]  

(3.16)

Combining (3.15) and (3.16), we get

\[ v(t) = v_1(t) + v_2(t) = \left( \frac{2}{T} t + l \right) \left[ u_0 \left( t + \frac{T}{2} \right) - u_0(t) \right] + \left( -\frac{2}{T} t + l \right) \left[ u_0(t) - u_0 \left( t - \frac{T}{2} \right) \right] \]  

(3.17)

Example 3.5

Express the waveform of Figure 3.14 as a sum of unit step functions.

Solution:

As in the previous example, we first find the equations of the linear segments 1 and 2 shown in Figure 3.15.
Following the same procedure as in the previous examples, we get

\[
v(t) = (2t + 1)[u_o(t) - u_o(t - 1)] + 3[u_o(t - 1) - u_o(t - 2)] \\
+ (-t + 3)[u_o(t - 2) - u_o(t - 3)]
\]

Multiplying the values in parentheses by the values in the brackets, we get

\[
v(t) = (2t + 1)u_o(t) - (2t + 1)u_o(t - 1) + 3u_o(t - 1) \\
- 3u_o(t - 2) + (-t + 3)u_o(t - 2) - (-t + 3)u_o(t - 3)
\]

or

\[
v(t) = (2t + 1)u_o(t) + [- (2t + 1) + 3]u_o(t - 1) \\
+ [-3 + (-t + 3)]u_o(t - 2) - (-t + 3)u_o(t - 3)
\]

and combining terms inside the brackets, we get

\[
v(t) = (2t + 1)u_o(t) - 2(t - 1)u_o(t - 1) - tu_o(t - 2) + (t - 3)u_o(t - 3)
\]

Two other functions of interest are the unit ramp function, and the unit impulse or delta function. We will introduce them with the examples that follow.

**Example 3.6**

In the network of Figure 3.16, where \( i_S \) is a constant source, the switch is closed at time \( t = 0 \).
Express the capacitor voltage \( v_C(t) \) as a function of the unit step.

**Solution:**

The current through the capacitor is \( i_C(t) = i_S = \text{constant} \), and the capacitor voltage \( v_C(t) \) is

\[
v_C(t) = \frac{1}{C} \int_{-\infty}^{t} i_C(\tau) d\tau^*
\]

where \( \tau \) is a dummy variable.

Since the switch closes at \( t = 0 \), we can express the current \( i_C(t) \) as

\[
i_C(t) = i_S u_0(t)
\]

and assuming that \( v_C(t) = 0 \) for \( t < 0 \), we can write (3.19) as

\[
v_C(t) = \frac{1}{C} \int_{-\infty}^{t} i_S u_0(\tau) d\tau = \frac{i_S}{C} \int_{-\infty}^{0} u_0(\tau) d\tau + \frac{i_S}{C} \int_{0}^{t} u_0(\tau) d\tau
\]

or

\[
v_C(t) = \frac{i_S}{C} tu_0(t)
\]

Therefore, we see that when a capacitor is charged with a constant current, the voltage across it is a linear function and forms a *ramp* with slope \( i_S / C \) as shown in Figure 3.17.

![Figure 3.17. Voltage across a capacitor when charged with a constant current source.](image_url)

---

*Since the initial condition for the capacitor voltage was not specified, we express this integral with \(-\infty\) at the lower limit of integration so that any non-zero value prior to \( t < 0 \) would be included in the integration.*
3.3 The Unit Ramp Function \( u_1(t) \)

The unit ramp function, denoted as \( u_1(t) \), is defined as

\[
\begin{align*}
    u_1(t) &= \int_{-\infty}^{t} u_0(\tau) \, d\tau \\
    \text{where } \tau \text{ is a dummy variable.}
\end{align*}
\]

We can evaluate the integral of (3.23) by considering the area under the unit step function \( u_0(t) \) from \(-\infty \) to \( t \) as shown in Figure 3.18.

\[
\begin{align*}
    \text{Figure 3.18. Area under the unit step function from } -\infty \text{ to } t
\end{align*}
\]

Therefore, we define \( u_1(t) \) as

\[
\begin{align*}
    u_1(t) &= \begin{cases} 
        0 & t < 0 \\
        t & t \geq 0
    \end{cases} \\
    \text{or } u_1(t) = \int_{-\infty}^{t} u_0(\tau) \, d\tau
\end{align*}
\]  

(3.24)

Since \( u_1(t) \) is the integral of \( u_0(t) \), then \( u_0(t) \) must be the derivative of \( u_1(t) \), i.e.,

\[
\frac{d}{dt} u_1(t) = u_0(t)
\]

(3.25)

Higher order functions of \( t \) can be generated by repeated integration of the unit step function. For example, integrating \( u_0(t) \) twice and multiplying by 2, we define \( u_2(t) \) as

\[
\begin{align*}
    u_2(t) &= \begin{cases} 
        0 & t < 0 \\
        t^2 & t \geq 0
    \end{cases} \\
    \text{or } u_2(t) = 2 \int_{-\infty}^{t} u_1(\tau) \, d\tau
\end{align*}
\]

(3.26)

Similarly,

\[
\begin{align*}
    u_3(t) &= \begin{cases} 
        0 & t < 0 \\
        t^3 & t \geq 0
    \end{cases} \\
    \text{or } u_3(t) = 3 \int_{-\infty}^{t} u_2(\tau) \, d\tau
\end{align*}
\]

(3.27)

and in general,
The Unit Ramp Function

\[
\text{u}_n(t) = \begin{cases} 
0 & \text{if } t < 0 \\
\frac{t^n}{n!} & \text{if } t \geq 0
\end{cases} \quad \text{or} \quad \text{u}_n(t) = 3 \int_{-\infty}^{t} \text{u}_{n-1}(\tau)d\tau \tag{3.28}
\]

Also,

\[
\text{u}_{n-1}(t) = \frac{1}{n!}\frac{d^n}{dt^n} \text{u}_n(t) \tag{3.29}
\]

Example 3.7

In the network of Figure 3.19, the switch is closed at time \( t = 0 \) and \( i_L(t) = 0 \) for \( t < 0 \).

\[R \quad t = 0
\]

\[\downarrow \quad i_s \quad i_L(t) \quad v_L(t)
\]

\[\uparrow \quad L
\]

\[\begin{align*}
\text{Figure 3.19. Network for Example 3.7}
\end{align*}
\]

Express the inductor current \( i_L(t) \) in terms of the unit step function.

Solution:

The voltage across the inductor is

\[
v_L(t) = L\frac{di_L}{dt} \tag{3.30}
\]

and since the switch closes at \( t = 0 \),

\[
i_L(t) = i_s \text{u}_0(t) \tag{3.31}
\]

Therefore, we can write (3.30) as

\[
v_L(t) = L_i \frac{d}{dt} \text{u}_0(t) \tag{3.32}
\]

But, as we know, \( u_0(t) \) is constant (0 or 1) for all time except at \( t = 0 \) where it is discontinuous. Since the derivative of any constant is zero, the derivative of the unit step \( u_0(t) \) has a non-zero value only at \( t = 0 \). The derivative of the unit step function is defined in the next section.
3.4 The Delta Function $\delta(t)$

The unit impulse or delta function, denoted as $\delta(t)$, is the derivative of the unit step $u_0(t)$. It is also defined as

$$\int_{-\infty}^{t} \delta(\tau)d\tau = u_0(t) \quad (3.33)$$

and

$$\delta(t) = 0 \text{ for all } t \neq 0 \quad (3.34)$$

To better understand the delta function $\delta(t)$, let us represent the unit step $u_0(t)$ as shown in Figure 3.20 (a).

![Figure 3.20. Representation of the unit step as a limit.](image)

The function of Figure 3.20 (a) becomes the unit step as $\varepsilon \to 0$. Figure 3.20 (b) is the derivative of Figure 3.20 (a), where we see that as $\varepsilon \to 0$, $1/2\varepsilon$ becomes unbounded, but the area of the rectangle remains $1$. Therefore, in the limit, we can think of $\delta(t)$ as approaching a very large spike or impulse at the origin, with unbounded amplitude, zero width, and area equal to $1$.

Two useful properties of the delta function are the sampling property and the sifting property.

3.5 Sampling Property of the Delta Function $\delta(t)$

The sampling property of the delta function states that

$$f(t)\delta(t-a) = f(a)\delta(t) \quad (3.35)$$

or, when $a = 0$,

$$f(t)\delta(t) = f(0)\delta(t) \quad (3.36)$$
that is, multiplication of any function \( f(t) \) by the delta function \( \delta(t) \) results in sampling the function at the time instants where the delta function is not zero. The study of discrete-time systems is based on this property.

**Proof:**

Since \( \delta(t) = 0 \) for \( t < 0 \) and \( t > 0 \) then,

\[
f(t)\delta(t) = 0 \quad \text{for} \quad t < 0 \quad \text{and} \quad t > 0
\]

(3.37)

We rewrite \( f(t) \) as

\[
f(t) = f(0) + [f(t) - f(0)]
\]

(3.38)

Integrating (3.37) over the interval \(-\infty\) to \( t \) and using (3.38), we get

\[
\int_{-\infty}^{t} f(\tau)\delta(\tau)d\tau = \int_{-\infty}^{t} f(0)\delta(\tau)d\tau + \int_{-\infty}^{t} [f(\tau) - f(0)]\delta(\tau)d\tau
\]

(3.39)

The first integral on the right side of (3.39) contains the constant term \( f(0) \); this can be written outside the integral, that is,

\[
\int_{-\infty}^{t} f(0)\delta(\tau)d\tau = f(0)\int_{-\infty}^{t} \delta(\tau)d\tau
\]

(3.40)

The second integral of the right side of (3.39) is always zero because

\[
\delta(t) = 0 \quad \text{for} \quad t < 0 \quad \text{and} \quad t > 0
\]

and

\[
[f(\tau) - f(0)]|_{\tau=0} = f(0) - f(0) = 0
\]

Therefore, (3.39) reduces to

\[
\int_{-\infty}^{t} f(\tau)\delta(\tau)d\tau = f(0)\int_{-\infty}^{t} \delta(\tau)d\tau
\]

(3.41)

Differentiating both sides of (3.41), and replacing \( \tau \) with \( t \), we get

\[
f(t)\delta(t) = f(0)\delta(t)
\]

(3.42)

**3.6 Sifting Property of the Delta Function \( \delta(t) \)**

The *sifting property of the delta function* states that
that is, if we multiply any function \( f(t) \) by \( \delta(t - \alpha) \), and integrate from \(-\infty\) to \(+\infty\), we will obtain the value of \( f(t) \) evaluated at \( t = \alpha \).

**Proof:**

Let us consider the integral

\[
\int_{a}^{b} f(t)\delta(t - \alpha)dt \quad \text{where} \quad a < \alpha < b
\]

We will use integration by parts to evaluate this integral. We recall from the derivative of products that

\[
d(xy) = xdy + ydx \quad \text{or} \quad xdy = d(xy) - ydx
\]

and integrating both sides we get

\[
\int xdy = xy - \int ydx
\]

Now, we let \( x = f(t) \); then, \( dx = f'(t) \). We also let \( dy = \delta(t - \alpha) \); then, \( y = u_0(t - \alpha) \). By substitution into (3.46), we get

\[
\int_{a}^{b} f(t)\delta(t - \alpha)dt = f(t)u_0(t - \alpha)\bigg|_{a}^{b} - \int_{a}^{b} u_0(t - \alpha)f'(t)dt
\]

We have assumed that \( a < \alpha < b \); therefore, \( u_0(t - \alpha) = 0 \) for \( \alpha < a \), and thus the first term of the right side of (3.47) reduces to \( f(b) \). Also, the integral on the right side is zero for \( \alpha < a \), and therefore, we can replace the lower limit of integration \( a \) by \( \alpha \). We can now rewrite (3.47) as

\[
\int_{a}^{b} f(t)\delta(t - \alpha)dt = f(b) - \int_{a}^{b} f'(t)dt = f(b) - f(b) + f(\alpha)
\]

and letting \( a \to -\infty \) and \( b \to +\infty \) for any \( |\alpha| < \infty \), we get

\[
\int_{-\infty}^{\infty} f(t)\delta(t - \alpha)dt = f(\alpha)
\]

**Sifting Property of \( \delta(t) \)**
3.7 Higher Order Delta Functions

An \( nth \)-order delta function is defined as the \( nth \) derivative of \( u_0(t) \), that is,

\[
\delta^n(t) = \frac{\delta^n}{dt^n}[u_0(t)]
\]  

(3.49)

The function \( \delta'(t) \) is called \( doublet \), \( \delta''(t) \) is called \( triplet \), and so on. By a procedure similar to the derivation of the sampling property of the delta function, we can show that

\[
f(t)\delta'(t-a) = f(a)\delta'(t-a) - f'(a)\delta(t-a)
\]  

(3.50)

Also, the derivation of the sifting property of the delta function can be extended to show that

\[
\int_{-\infty}^{\infty} f(t)\delta^n(t-\alpha)dt = (-1)^n \frac{d^n}{dt^n}[f(t)] \bigg|_{t=\alpha}
\]  

(3.51)

Example 3.8

Evaluate the following expressions:

a. \( 3t^4\delta(t-1) \)

b. \( \int_{-\infty}^{\infty} t\delta(t-2)dt \)

c. \( t^2\delta'(t-3) \)

Solution:

a. The sampling property states that \( f(t)\delta(t-a) = f(a)\delta(t-a) \) For this example, \( f(t) = 3t^4 \) and \( a = 1 \). Then,

\[
3t^4\delta(t-1) = \{3t^4\bigg|_{t=1}\} \delta(t-1) = 3\delta(t-1)
\]

b. The sifting property states that \( \int_{-\infty}^{\infty} f(t)\delta(t-\alpha)dt = f(\alpha) \). For this example, \( f(t) = t \) and \( \alpha = 2 \).

Then,

\[
\int_{-\infty}^{\infty} t\delta(t-2)dt = f(2) = t|_{t=2} = 2
\]

c. The given expression contains the doublet; therefore, we use the relation
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\[ f(t)\delta'(t-a) = f(a)\delta'(t-a) - f'(a)\delta(t-a) \]

Then, for this example,

\[ t^2\delta'(t-3) = t^2\delta(t-3) - \frac{d}{dt}t^2\delta(t-3) = 9\delta'(t-3) - 6\delta(t-3) \]

Example 3.9

a. Express the voltage waveform \( v(t) \) shown in Figure 3.21 as a sum of unit step functions for the time interval \(-1 < t < 7 \) s.

b. Using the result of part (a), compute the derivative of \( v(t) \) and sketch its waveform.

Solution:

a. We first derive the equations for the linear segments of the given waveform. These are shown in Figure 3.22.

Next, we express \( v(t) \) in terms of the unit step function \( u_0(t) \), and we get

\[ v(t) = 2t[u_0(t+1) - u_0(t-1)] + 2[u_0(t-1) - u_0(t-2)] + (-t + 5)[u_0(t-2) - u_0(t-4)] + [u_0(t-4) - u_0(t-5)] + (-t + 6)[u_0(t-5) - u_0(t-7)] \]

(3.52)

Multiplying and collecting like terms in (3.52), we get
Figure 3.22. Equations for the linear segments of Figure 3.21

\[ v(t) = 2tu_0(t+1) - 2tu_0(t-1) - 2u_0(t-1) - 2u_0(t-2) - tu_0(t-2) + 5u_0(t-2) + tu_0(t-4) - 5u_0(t-4) + u_0(t-4) - u_0(t-5) - tu_0(t-5) + 6u_0(t-5) + tu_0(t-7) - 6u_0(t-7) \]

or

\[ v(t) = 2tu_0(t+1) + (-2t+2)u_0(t-1) + (-t+3)u_0(t-2) + (t-4)u_0(t-4) + (-t+5)u_0(t-5) + (t-6)u_0(t-7) \]

b. The derivative of \( v(t) \) is

\[
\frac{dv}{dt} = 2u_0(t+1) + 2t\delta(t+1) - 2u_0(t-1) + (-2t+2)\delta(t-1) - u_0(t-2) + (-t+3)\delta(t-2) + u_0(t-4) + (t-4)\delta(t-4) - u_0(t-5) + (-t+5)\delta(t-5) + u_0(t-7) + (t-6)\delta(t-7)
\]

From the given waveform, we observe that discontinuities occur only at \( t = -1, \ t = 2, \) and \( t = 7. \) Therefore, \( \delta(t-1) = 0, \ \delta(t-4) = 0, \) and \( \delta(t-5) = 0, \) and the terms that contain these delta functions vanish. Also, by application of the sampling property,

\[
2t\delta(t+1) = \left\{ 2t \big|_{t=-1} \right\} \delta(t+1) = -2\delta(t+1)
\]

\[
(-t+3)\delta(t-2) = \left\{ (-t+3) \big|_{t=2} \right\} \delta(t-2) = \delta(t-2)
\]

\[
(t-6)\delta(t-7) = \left\{ (t-6) \big|_{t=7} \right\} \delta(t-7) = \delta(t-7)
\]

and by substitution into (3.53), we get
The plot of \( \frac{dv}{dt} \) is shown in Figure 3.23.

\[
\frac{dv}{dt} = 2u_0(t+1) - 2\delta(t+1) - 2u_0(t-1) - u_0(t-2) + \delta(t-2) + u_0(t-4) - u_0(t-5) + u_0(t-7) + \delta(t-7)
\]  

(3.54)

We observe that a negative spike of magnitude 2 occurs at \( t = -1 \), and two positive spikes of magnitude 1 occur at \( t = 2 \), and \( t = 7 \). These spikes occur because of the discontinuities at these points.

MATLAB\textsuperscript{*} has built-in functions for the unit step, and the delta functions. These are denoted by the names of the mathematicians who used them in their work. The unit step function \( u_0(t) \) is referred to as \texttt{Heaviside(t)}, and the delta function \( \delta(t) \) is referred to as \texttt{Dirac(t)}. Their use is illustrated with the examples below:

```matlab
syms k a t;
% Define symbolic variables
u = k*sym('Heaviside(t-a)');
% Create unit step function at t = a

u = k*Heaviside(t-a)

d = diff(u)
% Compute the derivative of the unit step function

d = k*Dirac(t-a)
```

* An introduction to MATLAB\textsuperscript{®} is given in Appendix A.
int(d) % Integrate the delta function
ans =
Heaviside(t-a)*k

3.8 Summary

• The unit step function \( u_0(t) \) that is defined as

\[
u_0(t) = \begin{cases} 
0 & t < 0 \\
1 & t > 0 
\end{cases}
\]

• The unit step function offers a convenient method of describing the sudden application of a voltage or current source.

• The unit ramp function, denoted as \( u_1(t) \), is defined as

\[
u_1(t) = \int_{-\infty}^{t} u_0(\tau) d\tau
\]

• The unit impulse or delta function, denoted as \( \delta(t) \), is the derivative of the unit step \( u_0(t) \). It is also defined as

\[
\int_{-\infty}^{t} \delta(\tau) d\tau = u_0(t)
\]

and

\[
\delta(t) = 0 \text{ for all } t \neq 0
\]

• The sampling property of the delta function states that

\[
f(t)\delta(t-a) = f(a)\delta(t)
\]

or, when \( a = 0 \),

\[
f(t)\delta(t) = f(0)\delta(t)
\]

• The sifting property of the delta function states that

\[
\int_{-\infty}^{\infty} f(t)\delta(t-\alpha) dt = f(\alpha)
\]

• The sampling property of the doublet function \( \delta'(t) \) states that

\[
f(t)\delta'(t-a) = f(a)\delta'(t-a) - f'(a)\delta(t-a)
\]
3.9 Exercises

1. Evaluate the following functions:
   
   a. \( \sin t \delta(t - \frac{\pi}{6}) \)
   
   b. \( \cos 2t \delta(t - \frac{\pi}{4}) \)
   
   c. \( \cos^2 t \delta(t - \frac{\pi}{2}) \)
   
   d. \( \tan 2t \delta(t - \frac{\pi}{8}) \)
   
   e. \( \int_{-\infty}^{\infty} t^2 e^{-t} \delta(t - 2) \, dt \)
   
   f. \( \sin^2 t \delta(t - \frac{\pi}{2}) \)

2. 
   a. Express the voltage waveform \( v(t) \) shown in Figure 3.24, as a sum of unit step functions for the time interval \( 0 < t < 7 \, s \).
   
   b. Using the result of part (a), compute the derivative of \( v(t) \), and sketch its waveform.

   ![Figure 3.24. Waveform for Exercise 2](image)
3.10 Solutions to Exercises

1. We apply the sampling property of the δ(t) function for all expressions except (e) where we apply
the sifting property. For part (f) we apply the sampling property of the doublet.

We recall that the sampling property states that \( f(t)\delta(t-a) = f(a)\delta(t-a) \). Thus,

\[ a. \quad \sin t \delta \left( t - \frac{\pi}{6} \right) = \sin \left( \frac{\pi}{6} \right) \delta \left( t - \frac{\pi}{6} \right) = \sin \frac{\pi}{6} \delta \left( t - \frac{\pi}{6} \right) = 0.5 \delta \left( t - \frac{\pi}{6} \right) \]

\[ b. \quad \cos 2t \delta \left( t - \frac{\pi}{4} \right) = \cos \left( \frac{\pi}{4} \right) \delta \left( t - \frac{\pi}{4} \right) = 0 \]

\[ c. \quad \cos^2 t \delta \left( t - \frac{\pi}{2} \right) = \frac{1}{2} \left( 1 + \cos 2t \right) \delta \left( t - \frac{\pi}{2} \right) = \frac{1}{2} \left( 1 + \cos \pi \right) \delta \left( t - \frac{\pi}{2} \right) = \frac{1}{2} \left( 1 - 1 \right) \delta \left( t - \frac{\pi}{2} \right) = 0 \]

\[ d. \quad \tan 2t \delta \left( t - \frac{\pi}{8} \right) = \tan \left( \frac{\pi}{8} \right) \delta \left( t - \frac{\pi}{8} \right) = \tan \frac{\pi}{4} \delta \left( t - \frac{\pi}{8} \right) = \delta \left( t - \frac{\pi}{8} \right) \]

We recall that the sampling property states that \( \int_{-\infty}^{\infty} f(t)\delta(t-\alpha)\,dt = f(\alpha) \). Thus,

\[ e. \quad \int_{-\infty}^{\infty} t^2 e^{-t} \delta(t-2)\,dt = t^2 e^{-t} \bigg|_{t=2} = 4e^{-2} = 0.54 \]

We recall that the sampling property for the doublet states that

\[ f(t)\delta(t-a) = f(a)\delta'(t-a) - f'(a)\delta(t-a) \]

Thus,

\[ \sin^2 t \delta' \left( t - \frac{\pi}{2} \right) = \sin^2 \left( \frac{\pi}{2} \right) \delta' \left( t - \frac{\pi}{2} \right) - \frac{d}{dt} \sin^2 t \bigg|_{t=\pi/2} \delta \left( t - \frac{\pi}{2} \right) \]

\[ f. \quad \sin^2 \delta' \left( t - \frac{\pi}{2} \right) = \frac{1}{2} \left( 1 - \cos 2t \right) \bigg|_{t=\pi/2} \delta' \left( t - \frac{\pi}{2} \right) - \sin 2t \bigg|_{t=\pi/2} \delta \left( t - \frac{\pi}{2} \right) \]

\[ \quad = \frac{1}{2} \left( 1 + 1 \right) \delta' \left( t - \frac{\pi}{2} \right) - \sin \pi \delta \left( t - \frac{\pi}{2} \right) = \delta' \left( t - \frac{\pi}{2} \right) \]

2.

\[ v(t) = e^{-2t} [u_0(t) - u_0(t-2)] + (10t - 30)[u_0(t-2) - u_0(t-3)] + (-10t + 50)[u_0(t-3) - u_0(t-5)] + (10t - 70)[u_0(t-5) - u_0(t-7)] \]

or
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\[ v(t) = e^{-2t}u_0(t) - e^{-2t}u_0(t-2) + 10tu_0(t-2) - 30u_0(t-2) - 10tu_0(t-3) + 30u_0(t-3) \]
\[ - 10tu_0(t-3) + 50u_0(t-3) + 10tu_0(t-5) - 50u_0(t-5) + 10tu_0(t-5) \]
\[ - 70u_0(t-5) - 10tu_0(t-7) + 70u_0(t-7) \]
\[ = e^{-2t}u_0(t) + (-e^{-2t} + 10t - 30)u_0(t-2) + (-20t + 80)u_0(t-3) + (20t - 120)u_0(t-5) \]
\[ + (-10t + 70)u_0(t-7) \]

b.
\[ \frac{dv}{dt} = -2e^{-2t}u_0(t) + e^{-2t}\delta(t) + (2e^{-2t} + 10)u_0(t-2) + (-e^{-2t} + 10t - 30)\delta(t-2) \]
\[ - 20u_0(t-3) + (20t + 80)\delta(t-3) + 20u_0(t-5) + (20t - 120)\delta(t-5) \]
\[ = e^{-2t}[u_0(t) - u_0(t-2)] - 10\delta(t-2) + 10[u_0(t-2) - u_0(t-3)] + 20\delta(t-3) \]
\[ - 10[u_0(t-3) - u_0(t-5)] - 20\delta(t-5) + 10[u_0(t-5) - u_0(t-7)] \]

Referring to the given waveform we observe that discontinuities occur only at \( t = 2 \), \( t = 3 \), and \( t = 5 \). Therefore, \( \delta(t) = 0 \) and \( \delta(t-7) = 0 \). Also, by the sampling property of the delta function

\[ (-e^{-2t} + 10t - 30)\delta(t-2) = (-e^{-2t} + 10t - 30)|_{t=2}\delta(t-2) \approx -10\delta(t-2) \]
\[ (-20t + 80)\delta(t-3) = (-20t + 80)|_{t=3}\delta(t-3) = 20\delta(t-3) \]
\[ (20t - 120)\delta(t-5) = (20t - 120)|_{t=5}\delta(t-5) = -20\delta(t-5) \]

and with these simplifications (1) above reduces to

\[ \frac{dv}{dt} = -2e^{-2t}u_0(t) + e^{-2t}u_0(t-2) + 20u_0(t-2) - 10\delta(t-2) \]
\[ - 20u_0(t-3) + 20\delta(t-3) + 20u_0(t-5) - 20\delta(t-5) - 10u_0(t-7) \]
\[ = -2e^{-2t}[u_0(t) - u_0(t-2)] - 10\delta(t-2) + 10[u_0(t-2) - u_0(t-3)] + 20\delta(t-3) \]
\[ - 10[u_0(t-3) - u_0(t-5)] - 20\delta(t-5) + 10[u_0(t-5) - u_0(t-7)] \]

The waveform for \( \frac{dv}{dt} \) is shown below.
Chapter 4

The Laplace Transformation

This chapter begins with an introduction to the Laplace transformation, definitions, and properties of the Laplace transformation. The initial value and final value theorems are also discussed and proved. It concludes with the derivation of the Laplace transform of common functions of time, and the Laplace transforms of common waveforms.

4.1 Definition of the Laplace Transformation

The two-sided or bilateral Laplace Transform pair is defined as

\[
\mathcal{L} \{ f(t) \} = F(s) = \int_{-\infty}^{\infty} f(t)e^{-st} \, dt \quad (4.1)
\]

\[
\mathcal{L}^{-1} \{ F(s) \} = f(t) = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F(s)e^{st} \, ds \quad (4.2)
\]

where \( \mathcal{L} \{ f(t) \} \) denotes the Laplace transform of the time function \( f(t) \), \( \mathcal{L}^{-1} \{ F(s) \} \) denotes the Inverse Laplace transform, and \( s \) is a complex variable whose real part is \( \sigma \), and imaginary part \( \omega \), that is, \( s = \sigma + j\omega \).

In most problems, we are concerned with values of time \( t \) greater than some reference time, say \( t = t_0 = 0 \), and since the initial conditions are generally known, the two-sided Laplace transform pair of (4.1) and (4.2) simplifies to the unilateral or one-sided Laplace transform defined as

\[
\mathcal{L} \{ f(t) \} = F(s) = \int_{t_0}^{\infty} f(t)e^{-st} \, dt = \int_{0}^{\infty} f(t)e^{-st} \, dt \quad (4.3)
\]

\[
\mathcal{L}^{-1} \{ F(s) \} = f(t) = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F(s)e^{st} \, ds \quad (4.4)
\]

The Laplace Transform of (4.3) has meaning only if the integral converges (reaches a limit), that is, if

\[
\left| \int_{0}^{\infty} f(t)e^{-st} \, dt \right| < \infty \quad (4.5)
\]

To determine the conditions that will ensure us that the integral of (4.3) converges, we rewrite (4.5)
as

$$\left| \int_0^\infty f(t)e^{-\sigma t}e^{-j\omega t} dt \right| < \infty \quad (4.6)$$

The term $e^{-j\omega t}$ in the integral of (4.6) has magnitude of unity, i.e., $|e^{-j\omega t}| = 1$, and thus the condition for convergence becomes

$$\left| \int_0^\infty f(t)e^{-\sigma t} dt \right| < \infty \quad (4.7)$$

Fortunately, in most engineering applications the functions $f(t)$ are of exponential order*. Then, we can express (4.7) as,

$$\left| \int_0^\infty f(t)e^{-\sigma t} dt \right| < \left| \int_0^\infty ke^{\sigma_0 t}e^{-\sigma t} dt \right| \quad (4.8)$$

and we see that the integral on the right side of the inequality sign in (4.8), converges if $\sigma > \sigma_0$. Therefore, we conclude that if $f(t)$ is of exponential order, $\mathcal{L}\{f(t)\}$ exists if

$$\text{Re}\{s\} = \sigma > \sigma_0 \quad (4.9)$$

where $\text{Re}\{s\}$ denotes the real part of the complex variable $s$.

Evaluation of the integral of (4.4) involves contour integration in the complex plane, and thus, it will not be attempted in this chapter. We will see, in the next chapter, that many Laplace transforms can be inverted with the use of a few standard pairs, and therefore, there is no need to use (4.4) to obtain the Inverse Laplace transform.

In our subsequent discussion, we will denote transformation from the time domain to the complex frequency domain, and vice versa, as

$$f(t) \Leftrightarrow F(s) \quad (4.10)$$

### 4.2 Properties of the Laplace Transform

#### 1. Linearity Property

The linearity property states that if

$$f_1(t), f_2(t), \ldots, f_n(t)$$

have Laplace transforms

---

* A function $f(t)$ is said to be of exponential order if $|f(t)| < ke^{\sigma_0 t}$ for all $t \geq 0$.  

---
Properties of the Laplace Transform

respectively, and

\[ c_1, c_2, \ldots, c_n \]

are arbitrary constants, then,

\[ c_1 f_1(t) + c_2 f_2(t) + \ldots + c_n f_n(t) \Leftrightarrow c_1 F_1(s) + c_2 F_2(s) + \ldots + c_n F_n(s) \]  \hspace{1cm} (4.11)

**Proof:**

\[ \mathcal{L} \{ c_1 f_1(t) + c_2 f_2(t) + \ldots + c_n f_n(t) \} = \int_{t_0}^{\infty} \left[ c_1 f_1(t) + c_2 f_2(t) + \ldots + c_n f_n(t) \right] dt \]

\[ = c_1 \int_{t_0}^{\infty} f_1(t) e^{-st} dt + c_2 \int_{t_0}^{\infty} f_2(t) e^{-st} dt + \ldots + c_n \int_{t_0}^{\infty} f_n(t) e^{-st} dt \]

\[ = c_1 F_1(s) + c_2 F_2(s) + \ldots + c_n F_n(s) \]

**Note 1:**

It is desirable to multiply \( f(t) \) by \( u_0(t) \) to eliminate any unwanted non-zero values of \( f(t) \) for \( t < 0 \).

**2. Time Shifting Property**

The *time shifting property* states that a right shift in the time domain by \( a \) units, corresponds to multiplication by \( e^{-as} \) in the complex frequency domain. Thus,

\[ f(t-a)u_0(t-a) \Leftrightarrow e^{-as} F(s) \]  \hspace{1cm} (4.12)

**Proof:**

\[ \mathcal{L} \{ f(t-a)u_0(t-a) \} = \int_{0}^{a} 0 e^{-st} dt + \int_{a}^{\infty} f(t-a) e^{-st} dt \]  \hspace{1cm} (4.13)

Now, we let \( t-a = \tau \); then, \( t = \tau + a \) and \( dt = d\tau \). With these substitutions, the second integral on the right side of (4.13) becomes

\[ \int_{0}^{\infty} f(\tau) e^{-s(\tau+a)} d\tau = e^{-as} \int_{0}^{\infty} f(\tau) e^{-s\tau} d\tau = e^{-as} F(s) \]

**3. Frequency Shifting Property**

The *frequency shifting property* states that if we multiply some time domain function \( f(t) \) by an exponential function \( e^{-at} \) where \( a \) is an arbitrary positive constant, this multiplication will produce a shift of the \( s \) variable in the complex frequency domain by \( a \) units. Thus,
Proof:

\[ \mathcal{L} \{ e^{-at} f(t) \} = \int_0^\infty e^{-st} e^{-at} dt = \int_0^\infty f(t)e^{-(s+a)t} dt = F(s + a) \]

Note 2:

A change of scale is represented by multiplication of the time variable \( t \) by a positive scaling factor \( a \). Thus, the function \( f(t) \) after scaling the time axis, becomes \( f(at) \).

4. Scaling Property

Let \( a \) be an arbitrary positive constant; then, the scaling property states that

\[ f(at) \Leftrightarrow \frac{1}{a} F\left( \frac{s}{a} \right) \]  

(4.15)

Proof:

\[ \mathcal{L} \{ f(at) \} = \int_0^\infty f(at) e^{-st} dt \]

and letting \( t = \tau/a \), we get

\[ \mathcal{L} \{ f(at) \} = \int_0^\infty f(\tau)e^{-s(\tau/a)} \left( \frac{\tau}{a} \right) d(\tau) = \frac{1}{a} \int_0^\infty f(\tau)e^{-s\tau/a} d(\tau) = \frac{1}{a} F\left( \frac{s}{a} \right) \]

Note 3:

Generally, the initial value of \( f(t) \) is taken at \( t = 0^- \) to include any discontinuity that may be present at \( t = 0 \). If it is known that no such discontinuity exists at \( t = 0^- \), we simply interpret \( f(0^-) \) as \( f(0) \).

5. Differentiation in Time Domain

The differentiation in time domain property states that differentiation in the time domain corresponds to multiplication by \( s \) in the complex frequency domain, minus the initial value of \( f(t) \) at \( t = 0^- \). Thus,

\[ f'(t) = \frac{d}{dt} f(t) \Leftrightarrow sF(s) - f(0^-) \]  

(4.16)

Proof:

\[ \mathcal{L} \{ f'(t) \} = \int_0^\infty f'(t)e^{-st} dt \]
Properties of the Laplace Transform

Using integration by parts where

\[ \int vdu = uv - \int udv \]  \hspace{1cm} (4.17)

we let \( du = f'(t) \) and \( v = e^{-st} \). Then, \( u = f(t) \), \( dv = -se^{-st} \), and thus

\[ L \{ f'(t) \} = f(t)e^{-st}\bigg|_{0^-}^{\infty} + s \int_0^\infty f(t)e^{-st} dt = \lim_{a \to \infty} \bigg[ f(t)e^{-st}\bigg|_{0^-}^{a} \bigg] + sF(s) \]

\[ = \lim_{a \to \infty} \bigg[ e^{-sa}f(a) - f(0^-) \bigg] + sF(s) = 0 - f'(0^-) + sF(s) \]

The time differentiation property can be extended to show that

\[ \frac{d^2}{dt^2} f(t) \Leftrightarrow s^2F(s) - sf(0^-) - f'(0^-) \]  \hspace{1cm} (4.18)

\[ \frac{d^3}{dt^3} f(t) \Leftrightarrow s^3F(s) - s^2f(0^-) - sf'(0^-) - f''(0^-) \]  \hspace{1cm} (4.19)

and in general

\[ \frac{d^n}{dt^n} f(t) \Leftrightarrow s^nF(s) - s^{n-1}f(0^-) - s^{n-2}f'(0^-) - \ldots - f^{(n-1)}(0^-) \]  \hspace{1cm} (4.20)

To prove (4.18), we let

\[ g(t) = f'(t) = \frac{d}{dt} f(t) \]

and as we found above,

\[ \mathcal{L} \{ g'(t) \} = sL \{ g(t) \} - g(0^-) \]

Then,

\[ \mathcal{L} \{ f''(t) \} = sL \{ f'(t) \} - f'(0^-) = s[sL \{ f(t) \} - f(0^-)] - f'(0^-) \]

\[ = s^2F(s) - sf(0^-) - f'(0^-) \]

Relations (4.19) and (4.20) can be proved by similar procedures.

We must remember that the terms \( f(0^-) \), \( f'(0^-) \), \( f''(0^-) \), and so on, represent the initial conditions. Therefore, when all initial conditions are zero, and we differentiate a time function \( f(t) \) \( n \) times, this corresponds to \( F(s) \) multiplied by \( s \) to the \( nth \) power.
6. Differentiation in Complex Frequency Domain

This property states that differentiation in complex frequency domain and multiplication by minus one, corresponds to multiplication of \( f(t) \) by \( t \) in the time domain. In other words,

\[
\mathcal{L} \{tf(t)\} = Fs = \int_0^\infty ft e^{-st} dt
\]

Proof:

\[
\mathcal{L} \{f(t)\} = F(s) = \int_0^\infty f(t)e^{-st} dt
\]

Differentiating with respect to \( s \), and applying Leibnitz's rule* for differentiation under the integral, we get

\[
\frac{d}{ds} F(s) = \frac{d}{ds} \int_0^\infty f(t)e^{-st} dt = \int_0^\infty \frac{\partial}{\partial s} e^{-st} f(t) dt = \int_0^\infty -te^{-st} f(t) dt
\]

\[
= -\int_0^\infty [tf(t)]e^{-st} dt = -\mathcal{L} \{tf(t)\}
\]

In general,

\[
t^n f(t) \leftrightarrow (-1)^n \frac{d^n}{ds^n} F(s)
\]

The proof for \( n \geq 2 \) follows by taking the second and higher-order derivatives of \( F(s) \) with respect to \( s \).

7. Integration in Time Domain

This property states that integration in time domain corresponds to \( F(s) \) divided by \( s \) plus the initial value of \( f(t) \) at \( t = 0^- \), also divided by \( s \). That is,

\[
\int_{-\infty}^t f(\tau) d\tau \leftrightarrow \frac{F(s)}{s} + \frac{f(0^-)}{s}
\]

* This rule states that if a function of a parameter \( \alpha \) is defined by the equation \( F(\alpha) = \int_a^b f(x, \alpha) dx \) where \( f \) is some known function of integration \( x \) and the parameter \( \alpha \), \( a \) and \( b \) are constants independent of \( x \) and \( \alpha \), and the partial derivative \( \partial f/\partial \alpha \) exists and it is continuous, then \( \frac{dF}{d\alpha} = \int_a^b \frac{\partial f(x, \alpha)}{\partial (\alpha)} dx \).
Properties of the Laplace Transform

Proof:
We express the integral of (4.23) as two integrals, that is,

\[
\int_{-\infty}^{t} f(\tau)d\tau = \int_{-\infty}^{0} f(\tau)d\tau + \int_{0}^{t} f(\tau)d\tau \tag{4.24}
\]

The first integral on the right side of (4.24), represents a constant value since neither the upper, nor the lower limits of integration are functions of time, and this constant is an initial condition denoted as \(f(0^-)\). We will find the Laplace transform of this constant, the transform of the second integral on the right side of (4.24), and will prove (4.23) by the linearity property. Thus,

\[
\mathcal{L}\{f(0^-)\} = \int_{0}^{\infty} f(0^-)e^{-st}dt = f(0^-) \int_{0}^{\infty} e^{-st}dt = f(0^-) \left. \frac{e^{-st}}{-s} \right|_{0}^{\infty} = f(0^-) \times 0 - \left( \frac{f(0^-)}{s} \right) = \frac{f(0^-)}{s} \tag{4.25}
\]

This is the value of the first integral in (4.24). Next, we will show that

\[
\int_{0}^{t} f(\tau)d\tau \Leftrightarrow \frac{F(s)}{s}
\]

We let

\[
g(t) = \int_{0}^{t} f(\tau)d\tau
\]

then,

\[
g'(t) = f(t)
\]

and

\[
g(0) = \int_{0}^{0} f(\tau)d\tau = 0
\]

Now,

\[
\mathcal{L}\{g'(t)\} = G(s) = s\mathcal{L}\{g(t)\} - g(0^-) = G(s) - 0
\]

\[
s\mathcal{L}\{g(t)\} = G(s)
\]

\[
\mathcal{L}\{g(t)\} = \frac{G(s)}{s}
\]

\[
\mathcal{L}\left\{ \int_{0}^{t} f(\tau)d\tau \right\} = \frac{F(s)}{s} \tag{4.26}
\]

and the proof of (4.23) follows from (4.25) and (4.26).
8. Integration in Complex Frequency Domain

This property states that integration in complex frequency domain with respect to $s$ corresponds to division of a time function $f(t)$ by the variable $t$, provided that the limit $\lim_{t \to 0} \frac{f(t)}{t}$ exists. Thus,

$$\frac{f(t)}{t} \Leftrightarrow \int_s F(s) ds$$

(4.27)

**Proof:**

$$F(s) = \int_0^\infty f(t)e^{-st} dt$$

Integrating both sides from $s$ to $\infty$, we get

$$\int_s^\infty F(s) ds = \int_s^\infty \left[ \int_0^\infty f(t)e^{-st} dt \right] ds$$

Next, we interchange the order of integration, i.e.,

$$\int_s^\infty F(s) ds = \int_0^\infty \left[ \int_s^\infty e^{-st} ds \right] f(t) dt$$

and performing the inner integration on the right side integral with respect to $s$, we get

$$\int_s^\infty F(s) ds = \int_0^\infty \left[ -\frac{1}{t} e^{-st} \right]_s^\infty f(t) dt = \int_0^\infty \frac{f(t)}{t} e^{-st} dt = \mathcal{L} \left\{ \frac{f(t)}{t} \right\}$$

9. Time Periodicity

The time periodicity property states that a periodic function of time with period $T$ corresponds to the integral $\int_0^T f(t)e^{-st} dt$ divided by $(1 - e^{-sT})$ in the complex frequency domain. Thus, if we let $f(t)$ be a periodic function with period $T$, that is, $f(t) = f(t + nT)$, for $n = 1, 2, 3, \ldots$ we get the transform pair

$$f(t + nT) \Leftrightarrow \int_0^T f(t)e^{-st} dt = \frac{1}{1 - e^{-sT}}$$

(4.28)
Properties of the Laplace Transform

Proof:

The Laplace transform of a periodic function can be expressed as

\[ \mathcal{L} \{ f(t) \} = \int_{0}^{\infty} f(t)e^{-st} dt = \int_{0}^{T} f(t)e^{-st} dt + \int_{T}^{2T} f(t)e^{-st} dt + \int_{2T}^{3T} f(t)e^{-st} dt + \ldots \]

In the first integral of the right side, we let \( t = \tau \), in the second \( t = \tau + T \), in the third \( t = \tau + 2T \), and so on. The areas under each period of \( f(t) \) are equal, and thus the upper and lower limits of integration are the same for each integral. Then,

\[ \mathcal{L} \{ f(t) \} = \int_{0}^{T} f(\tau)e^{-s\tau} d\tau + \int_{0}^{T} f(\tau + T)e^{-s(\tau + T)} d\tau + \int_{0}^{T} f(\tau + 2T)e^{-s(\tau + 2T)} d\tau + \ldots \quad (4.29) \]

Since the function is periodic, i.e., \( f(\tau) = f(\tau + T) = f(\tau + 2T) = \ldots = f(\tau + nT) \), we can write (4.29) as

\[ \mathcal{L} \{ f(\tau) \} = (1 + e^{-sT} + e^{-2sT} + \ldots) \int_{0}^{T} f(\tau)e^{-s\tau} d\tau \quad (4.30) \]

By application of the binomial theorem, that is,

\[ 1 + a + a^2 + a^3 + \ldots = \frac{1}{1-a} \quad (4.31) \]

we find that expression (4.30) reduces to

\[ \mathcal{L} \{ f(\tau) \} = \frac{\int_{0}^{T} f(\tau)e^{-s\tau} d\tau}{1 - e^{-sT}} \]

10. Initial Value Theorem

The initial value theorem states that the initial value \( f(0^-) \) of the time function \( f(t) \) can be found from its Laplace transform multiplied by \( s \) and letting \( s \to \infty \). That is,

\[ \lim_{t \to 0} f(t) = \lim_{s \to \infty} sF(s) = f(0^-) \quad (4.32) \]

Proof:

From the time domain differentiation property,

\[ \frac{d}{dt} f(t) \Leftrightarrow sF(s) - f(0^-) \]

or
Taking the limit of both sides by letting \( s \to \infty \), we get

\[
\lim_{s \to \infty} [sF(s) - f(0^-)] = \lim_{s \to \infty} \left[ \lim_{T \to \infty} \int_0^T \frac{d}{dt} f(t) e^{-st} \, dt \right]
\]

Interchanging the limiting process, we get

\[
\lim_{s \to \infty} [sF(s) - f(0^-)] = \lim_{T \to \infty} \int_0^T \frac{d}{dt} f(t) \left[ \lim_{s \to \infty} e^{-st} \right] \, dt
\]

and since

\[
\lim_{s \to \infty} e^{-st} = 0
\]

the above expression reduces to

\[
\lim_{s \to \infty} [sF(s) - f(0^-)] = 0
\]

or

\[
\lim_{s \to \infty} sF(s) = f(0^-)
\]

11. Final Value Theorem

The final value theorem states that the final value \( f(\infty) \) of the time function \( f(t) \) can be found from its Laplace transform multiplied by \( s \), then, letting \( s \to 0 \). That is,

\[
\lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = f(\infty)
\]

(4.33)

Proof:

From the time domain differentiation property,

\[
\frac{d}{dt} f(t) \leftrightarrow sF(s) - f(0^-)
\]

or

\[
\mathcal{L} \left\{ \frac{d}{dt} f(t) \right\} = sF(s) - f(0^-) = \int_0^\infty \frac{d}{dt} f(t) e^{-st} \, dt
\]

Taking the limit of both sides by letting \( s \to 0 \), we get
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\[
\lim_{s \to 0} [sF(s) - f(0^-)] = \lim_{s \to 0} \left[ \lim_{T \to \infty, \varepsilon \to 0} \int_{-\varepsilon}^{T} \frac{d}{dt}f(t)e^{-st} dt \right]
\]

and by interchanging the limiting process, we get

\[
\lim_{s \to 0} [sF(s) - f(0^-)] = \lim_{T \to \infty, \varepsilon \to 0} \int_{-\varepsilon}^{T} \frac{d}{dt}f(t) \left[ \lim_{s \to 0} e^{-st} \right] dt
\]

Also, since

\[
\lim_{s \to 0} e^{-st} = 1
\]

the above expression reduces to

\[
\lim_{s \to 0} [sF(s) - f(0^-)] = \lim_{T \to \infty, \varepsilon \to 0} \int_{-\varepsilon}^{T} \frac{d}{dt}f(t) dt = \lim_{T \to \infty, \varepsilon \to 0} \int_{0}^{T} f(t) dt
\]

\[
= \lim_{T \to \infty, \varepsilon \to 0} [f(T) - f(\varepsilon)] = f(\infty) - f(0^-)
\]

and therefore,

\[
\lim_{s \to 0} sF(s) = f(\infty)
\]

12. Convolution in the Time Domain

**Convolution** *in the time domain* corresponds to multiplication in the complex frequency domain, that is,

\[
f_1(t) \ast f_2(t) \Leftrightarrow F_1(s)F_2(s)
\]

**Proof:**

\[
\mathcal{L} \{f_1(t) \ast f_2(t)\} = \mathcal{L} \left[ \int_{-\infty}^{\infty} f_1(\tau)f_2(t-\tau) d\tau \right] = \int_{0}^{\infty} \left[ \int_{0}^{\infty} f_1(\tau)f_2(t-\tau) d\tau \right] e^{-st} dt
\]

\[
= \int_{0}^{\infty} f_1(\tau) \left[ \int_{0}^{\infty} f_2(t-\tau) e^{-st} dt \right] d\tau
\]

We let \( t - \tau = \lambda \); then, \( t = \lambda + \tau \), and \( dt = d\lambda \). By substitution into (4.35),

---

* Convolution is the process of overlapping two signals. The convolution of two time functions \( f_1(t) \) and \( f_2(t) \) is denoted as \( f_1(t) \ast f_2(t) \), and by definition, \( f_1(t) \ast f_2(t) = \int_{-\infty}^{\infty} f_1(\tau)f_2(t-\tau) d\tau \) where \( \tau \) is a dummy variable. It is discussed in detail in Signals and Systems with MATLAB Applications by this author.
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\[ \mathcal{L} \{ f_1(t) * f_2(t) \} = \int_0^\infty f_1(\tau) \left[ \int_0^\infty f_2(\lambda) e^{-s(\lambda + \tau)} d\lambda \right] d\tau = \int_0^\infty f_1(\tau) e^{-st} \int_0^\infty f_2(\lambda) e^{-s\lambda} d\lambda \]

\[ = F_1(s)F_2(s) \]

13. Convolution in the Complex Frequency Domain

_Condvolution in the complex frequency domain_ divided by \(1/2\pi j\), corresponds to multiplication in the time domain. That is,

\[ f_1(t)f_2(t) \Leftrightarrow \frac{1}{2\pi j} F_1(s)F_2(s) \] (4.36)

**Proof:**

\[ \mathcal{L} \{ f_1(t)f_2(t) \} = \int_0^\infty f_1(t)f_2(t) e^{-st} dt \] (4.37)

and recalling that the Inverse Laplace transform from (4.2) is

\[ f_1(t) = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F_1(\mu)e^{\mu t} d\mu \]

by substitution into (4.37), we get

\[ \mathcal{L} \{ f_1(t)f_2(t) \} = \int_0^\infty \left[ \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F_1(\mu)e^{\mu t} d\mu \right] \int_0^\infty f_2(t) e^{-st} dt \]

\[ = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F_1(\mu) \left[ \int_0^\infty f_2(t) e^{-(s-\mu)t} dt \right] d\mu \]

We observe that the bracketed integral is \( F_2(s-\mu) \); therefore,

\[ L \{ f_1(t)f_2(t) \} = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F_1(\mu)F_2(s-\mu) d\mu = \frac{1}{2\pi j} F_1(s)F_2(s) \]

For easy reference, we have summarized the Laplace transform pairs and theorems in Table 4.1.

4.3 The Laplace Transform of Common Functions of Time

In this section, we will present several examples for finding the Laplace transform of common functions of time.

**Example 4.1**

Find \( \mathcal{L} \{ u_0(t) \} \)
<table>
<thead>
<tr>
<th>Property/Theorem</th>
<th>Time Domain</th>
<th>Complex Frequency Domain</th>
</tr>
</thead>
<tbody>
<tr>
<td>1. Linearity</td>
<td>$c_1 f_1(t) + c_2 f_2(t) + \cdots + c_n f_n(t)$</td>
<td>$c_1 F_1(s) + c_2 F_2(s) + \cdots + c_n F_n(s)$</td>
</tr>
<tr>
<td>2. Time Shifting</td>
<td>$f(t-a)u_0(t-a)$</td>
<td>$e^{-as}F(s)$</td>
</tr>
<tr>
<td>3. Frequency Shifting</td>
<td>$e^{-as}f(t)$</td>
<td>$F(s+a)$</td>
</tr>
<tr>
<td>4. Time Scaling</td>
<td>$f(at)$</td>
<td>$\frac{1}{a}F\left(\frac{s}{a}\right)$</td>
</tr>
<tr>
<td>5. Time Differentiation</td>
<td>$\frac{d}{dt}f(t)$</td>
<td>$sF(s) - f(0^-)$</td>
</tr>
<tr>
<td>6. Frequency Differentiation</td>
<td>$tf(t)$</td>
<td>$-\frac{d}{ds}F(s)$</td>
</tr>
<tr>
<td>7. Time Integration</td>
<td>$\int_{-\infty}^{t} f(\tau)d\tau$</td>
<td>$\frac{F(s) + f(0^-)}{s}$</td>
</tr>
<tr>
<td>8. Frequency Integration</td>
<td>$\frac{f(t)}{t}$</td>
<td>$\int_{s}^{\infty} F(s)ds$</td>
</tr>
<tr>
<td>9. Time Periodicity</td>
<td>$f(t + nT)$</td>
<td>$\int_{0}^{T} f(t)e^{-st}dt \over 1 - e^{-sT}$</td>
</tr>
<tr>
<td>10. Initial Value Theorem</td>
<td>$\lim_{t \to 0} f(t)$</td>
<td>$\lim_{s \to \infty} sF(s) = f(0^-)$</td>
</tr>
<tr>
<td>11. Final Value Theorem</td>
<td>$\lim_{t \to \infty} f(t)$</td>
<td>$\lim_{s \to 0} sF(s) = f(\infty)$</td>
</tr>
<tr>
<td>12. Time Convolution</td>
<td>$f_1(t) * f_2(t)$</td>
<td>$F_1(s)F_2(s)$</td>
</tr>
<tr>
<td>13. Frequency Convolution</td>
<td>$f_1(t)f_2(t)$</td>
<td>$\frac{1}{2\pi j} F_1(s)*F_2(s)$</td>
</tr>
</tbody>
</table>
Solution:
We start with the definition of the Laplace transform, that is,
\[ \mathcal{L} \{ f(t) \} = F(s) = \int_0^\infty f(t)e^{-st} \, dt \]
For this example,
\[ \mathcal{L} \{ u_0(t) \} = \int_0^\infty e^{-st} \, dt = \left. \frac{-e^{-st}}{s} \right|_0^\infty = 0 - \left( \frac{1}{s} \right) = \frac{1}{s} \]
Thus, we have obtained the transform pair
\[ u_0(t) \Leftrightarrow \frac{1}{s} \] (4.38)
for \( \text{Re}\{s\} = \sigma > 0 \).*

Example 4.2
Find \( \mathcal{L} \{ u_1(t) \} \)
Solution:
We apply the definition
\[ \mathcal{L} \{ f(t) \} = F(s) = \int_0^\infty f(t)e^{-st} \, dt \]
and for this example,
\[ \mathcal{L} \{ u_1(t) \} = \mathcal{L} \{ t \} = \int_0^\infty te^{-st} \, dt \]
We will perform integration by parts recalling that
\[ \int uv \, dv = uv - \int v \, du \] (4.39)
We let
\[ u = t \quad \text{and} \quad dv = e^{-st} \]
then,
\[ du = 1 \quad \text{and} \quad v = \frac{-e^{-st}}{s} \]
By substitution into (4.39),

---

* This condition was established in (4.9).
\[ \mathcal{L}\{t\} = -\frac{te^{-st}}{s} \bigg|_0^\infty - \int_0^\infty \frac{e^{-st}}{s} \, dt = \left[ -\frac{te^{-st}}{s} - \frac{e^{-st}}{s^2} \right]_0^\infty \] (4.40)

Since the upper limit of integration in (4.40) produces an indeterminate form, we apply L'Hôpital’s rule *, that is,

\[ \lim_{t \to \infty} te^{-st} = \lim_{t \to \infty} t e^{st} = \lim_{t \to \infty} \frac{d}{dt}(t) = \lim_{t \to \infty} \frac{1}{se^{st}} = 0 \]

Evaluating the second term of (4.40), we get \( \mathcal{L}\{t\} = \frac{1}{2} \frac{1}{s} \)

Thus, we have obtained the transform pair

\[ t \Leftrightarrow \frac{1}{s^2} \] (4.41)

for \( \sigma > 0 \).

**Example 4.3**

Find \( \mathcal{L}\{t^n u_0(t)\} \)

**Solution:**

To find the Laplace transform of this function, we must first review the gamma or generalized factorial function \( \Gamma(n) \) defined as

\[ \Gamma(n) = \int_0^\infty x^{n-1} e^{-x} \, dx \] (4.42)

---

* Often, the ratio of two functions, such as \( \frac{f(x)}{g(x)} \), for some value of \( x \), say \( a \), results in an indeterminate form. To work around this problem, we consider the limit \( \lim_{x \to a} \frac{f(x)}{g(x)} \), and we wish to find this limit, if it exists. To find this limit, we use L'Hôpital’s rule which states that if \( f(a) = g(a) = 0 \), and if the limit \( \frac{d}{dx}f(x)/\frac{d}{dx}g(x) \) as \( x \) approaches \( a \) exists, then, \( \lim_{x \to a} \frac{f(x)}{g(x)} = \lim_{x \to a} \left( \frac{d}{dx}f(x)/\frac{d}{dx}g(x) \right) \)
The integral of (4.42) is an improper integral but converges (approaches a limit) for all \( n > 0 \).

We will now derive the basic properties of the gamma function, and its relation to the well known factorial function

\[
n! = n(n-1)(n-2) \cdots 3 \cdot 2 \cdot 1
\]

The integral of (4.42) can be evaluated by performing integration by parts. Thus, in (4.39) we let

\[
u = e^{-x} \quad \text{and} \quad dv = x^{n-1} \, dx
\]

Then,

\[
du = -e^{-x} \, dx \quad \text{and} \quad v = \frac{x^n}{n}
\]

and (4.42) is written as

\[
\Gamma(n) = \frac{x^n e^{-x}}{n} \bigg|_{x=0}^{x=x_0} + \frac{1}{n} \int_0^x e^{-x} \, dx
\]

With the condition that \( n > 0 \), the first term on the right side of (4.43) vanishes at the lower limit \( x = 0 \). It also vanishes at the upper limit as \( x \to \infty \). This can be proved with L’Hôpital’s rule by differentiating both numerator and denominator \( m \) times, where \( m \geq n \). Then,

\[
\lim_{x \to \infty} \frac{n e^{-x}}{n} = \lim_{x \to \infty} \frac{x^n}{n e^x} = \lim_{x \to \infty} \frac{d^m}{dx^m} \frac{x^n}{n e^x} = \lim_{x \to \infty} \frac{d^{m-1}}{dx^{m-1}} \frac{n x^{-n-1}}{e^x} = \ldots
\]

\[
= \lim_{x \to \infty} \frac{n(n-1)(n-2)\cdots(n-m+1)x^{-m}}{n e^x} = \lim_{x \to \infty} \frac{(n-1)(n-2)\cdots(n-m+1)}{x^{m-n} e^x} = \frac{x^n}{n e^x} = 0
\]

Therefore, (4.43) reduces to

\[
\Gamma(n) = \frac{1}{n} \int_0^\infty x^{n-1} e^{-x} \, dx
\]

and with (4.42), we have

* Improper integrals are two types and these are:

a. \( \int_a^b f(x) \, dx \) where the limits of integration \( a \) or \( b \) or both are infinite

b. \( \int_a^b f(x) \, dx \) where \( f(x) \) becomes infinite at a value \( x \) between the lower and upper limits of integration inclusive.
The Laplace Transform of Common Functions of Time

\[ \Gamma(n) = \int_{0}^{\infty} x^{n-1} e^{-x} \, dx = \frac{1}{n} \int_{0}^{\infty} x^{n} e^{-x} \, dx \]  
(4.44)

By comparing the integrals in (4.44), we observe that

\[ \Gamma(n) = \frac{\Gamma(n+1)}{n} \]  
(4.45)

or

\[ n\Gamma(n) = \Gamma(n+1) \]  
(4.46)

It is convenient to use (4.45) for \( n < 0 \), and (4.46) for \( n > 0 \). From (4.45), we see that \( \Gamma(n) \) becomes infinite as \( n \to 0 \).

For \( n = 1 \), (4.42) yields

\[ \Gamma(1) = \int_{0}^{\infty} e^{-x} \, dx = -e^{-x}\bigg|_{0}^{\infty} = 1 \]  
(4.47)

and thus we have the important relation,

\[ \Gamma(1) = 1 \]  
(4.48)

From the recurring relation of (4.46), we obtain

\[ \begin{align*}
\Gamma(2) &= 1 \cdot \Gamma(1) = 1 \\
\Gamma(3) &= 2 \cdot \Gamma(2) = 2 \cdot 1 = 2! \\
\Gamma(4) &= 3 \cdot \Gamma(3) = 3 \cdot 2 = 3!
\end{align*} \]  
(4.49)

and in general

\[ \Gamma(n + 1) = n! \]  
(4.50)

for \( n = 1, 2, 3, \ldots \)

The formula of (4.50) is a noteworthy relation; it establishes the relationship between the \( \Gamma(n) \) function and the factorial \( n! \).

We now return to the problem of finding the Laplace transform pair for \( t^n u_0 t \), that is,

\[ \mathcal{L} \{ t^n u_0 t \} = \int_{0}^{\infty} t^n e^{-st} \, dt \]  
(4.51)

To make this integral resemble the integral of the gamma function, we let \( st = y \), or \( t = y/s \), and
thus \( dt = dy/s \). Now, we rewrite (4.51) as
\[
L \{ t^n u_0(t) \} = \int_0^\infty \left( \frac{y}{s} \right)^n e^{-s} dy = \frac{1}{s^{n+1}} \int_0^\infty y^n e^{-s} dy = \frac{\Gamma(n+1)}{s^{n+1}} = \frac{n!}{s^{n+1}}
\]

Therefore, we have obtained the transform pair
\[
\begin{align*}
t^n u_0(t) & \Leftrightarrow \frac{n!}{s^{n+1}} \\
\end{align*}
\]

(4.52)

for positive integers of \( n \) and \( \sigma > 0 \).

**Example 4.4**

Find \( \mathcal{L} \{ \delta(t) \} \)

**Solution:**
\[
\mathcal{L} \{ \delta(t) \} = \int_0^\infty \delta(t) e^{-st} dt
\]

and using the sifting property of the delta function, we get
\[
\mathcal{L} \{ \delta(t) \} = \int_0^\infty \delta(t) e^{-st} dt = e^{-s(0)} = 1
\]

Thus, we have the transform pair
\[
\delta(t) \Leftrightarrow 1
\]

(4.53)

for all \( \sigma \).

**Example 4.5**

Find \( \mathcal{L} \{ \delta(t-a) \} \)

**Solution:**
\[
\mathcal{L} \{ \delta(t-a) \} = \int_0^\infty \delta(t-a) e^{-st} dt
\]

and again, using the sifting property of the delta function, we get
\[
\mathcal{L} \{ \delta(t-a) \} = \int_0^\infty \delta(t-a) e^{-st} dt = e^{-as}
\]
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Thus, we have the transform pair

\[ \delta(t - a) \Leftrightarrow e^{-as} \quad (4.54) \]

for \( \sigma > 0 \).

**Example 4.6**

Find \( \mathcal{L} \{e^{-at}u_0(t)\} \)

**Solution:**

\[
\mathcal{L} \{e^{-at}u_0(t)\} = \int_0^\infty e^{-at} e^{-st} dt = \int_0^\infty e^{-(s+a)t} dt \\
= \left( \frac{1}{s+a} \right) e^{-(s+a)t} \bigg|_0^\infty = \frac{1}{s+a}
\]

Thus, we have the transform pair

\[ e^{-at}u_0(t) \Leftrightarrow \frac{1}{s+a} \quad (4.55) \]

for \( \sigma > -a \).

**Example 4.7**

Find \( \mathcal{L} \left\{ t^n e^{-at}u_0(t) \right\} \)

**Solution:**

For this example, we will use the transform pair of (4.52), i.e.,

\[ t^n u_0(t) \Leftrightarrow \frac{n!}{s^{n+1}} \quad (4.56) \]

and the frequency shifting property of (4.14), that is,

\[ e^{-at}f(t) \Leftrightarrow F(s + a) \quad (4.57) \]

Then, replacing \( s \) with \( s + a \) in (4.56), we get the transform pair
Chapter 4 The Laplace Transformation

\[ t^n e^{-at} u_0(t) \Leftrightarrow \frac{n!}{(s+a)^{n+1}} \]  \hspace{1cm} (4.58)

where \( n \) is a positive integer, and \( \sigma > -a \). Thus, for \( n = 1 \), we get the transform pair

\[ te^{-at} u_0(t) \Leftrightarrow \frac{1}{(s+a)^2} \]  \hspace{1cm} (4.59)

for \( \sigma > -a \).

For \( n = 2 \), we get the transform

\[ t^2 e^{-at} u_0(t) \Leftrightarrow \frac{2!}{(s+a)^3} \]  \hspace{1cm} (4.60)

and in general,

\[ t^n e^{-at} u_0(t) \Leftrightarrow \frac{n!}{(s+a)^{n+1}} \]  \hspace{1cm} (4.61)

for \( \sigma > -a \).

**Example 4.8**

Find \( \mathcal{L} \{ \sin \omega t \, u_0(t) \} \)

**Solution:**

\[ \mathcal{L} \{ \sin \omega t \, u_0(t) \} = \int_0^\infty (\sin \omega t)e^{-st} \, dt = \lim_{a \to \infty} \int_0^a (\sin \omega t)e^{-st} \, dt \]

and from tables of integrals*

\[ \int e^{ax} \sin bx \, dx = \frac{e^{ax}(a\sin bx - b\cos bx)}{a^2 + b^2} \]

Then,

\[ \mathcal{L} \{ \sin \omega t \, u_0(t) \} = \frac{\omega}{s^2 + \omega^2} \]

---

* This can also be derived from \( \sin \omega t = \frac{1}{j2}(e^{j\omega t} - e^{-j\omega t}) \), and the use of (4.55) where \( e^{-at} u_0(t) \Leftrightarrow \frac{1}{s+a} \). By the linearity property, the sum of these terms corresponds to the sum of their Laplace transforms. Therefore,

\[ L \{ \sin \omega t u_0(t) \} = \frac{j}{2} \left( \frac{1}{s-j\omega} - \frac{1}{s+j\omega} \right) = \frac{\omega}{s^2 + \omega^2} \]
The Laplace Transform of Common Functions of Time

\[ \mathcal{L} \{ \sin \omega t u_0(t) \} = \lim_{a \to \infty} e^{-st} \left( -s \sin \omega t - \omega \cos \omega t \right) \left| _{0}^{a} \right. \]

\[ = \lim_{a \to \infty} \left[ \frac{e^{-as} (s \sin \omega a - \omega \cos \omega a)}{s^2 + \omega^2} + \frac{\omega}{s^2 + \omega^2} \right] = \frac{\omega}{s^2 + \omega^2} \]

Thus, we have obtained the transform pair

\[ \sin \omega t u_0(t) \Leftrightarrow \frac{\omega}{s^2 + \omega^2} \quad (4.62) \]

for \( \sigma > 0 \).

**Example 4.9**

Find \( \mathcal{L} \{ \cos \omega t u_0(t) \} \)

**Solution:**

\[ L \{ \cos \omega t u_0(t) \} = \int_{0}^{\infty} (\cos \omega t) e^{-st} \, dt = \lim_{a \to \infty} \int_{0}^{a} (\cos \omega t) e^{-st} \, dt \]

and from tables of integrals*

\[ \int e^{ax} \cos bxdx = \frac{e^{ax}(a \cos bx + b \sin bx)}{a^2 + b^2} \]

Then,

\[ \mathcal{L} \{ \cos \omega t u_0(t) \} = \lim_{a \to \infty} e^{-st} \left( -s \cos \omega t + \omega \sin \omega t \right) \left| _{0}^{a} \right. \]

\[ = \lim_{a \to \infty} \left[ \frac{e^{-as} (-s \cos \omega a + \omega \sin \omega a)}{s^2 + \omega^2} + \frac{s}{s^2 + \omega^2} \right] = \frac{s}{s^2 + \omega^2} \]

Thus, we have the transform pair

* We can use the relation \( \cos \omega t = \frac{1}{2}(e^{j\omega t} + e^{-j\omega t}) \) and the linearity property, as in the derivation of the transform of \( \sin \omega t \) on the footnote of the previous page. We can also use the transform pair \( \frac{d}{dt} f(t) \Leftrightarrow s F(s) - f(0^-) \); this is the time differentiation property of (4.16). Applying this transform pair for this derivation, we get

\[ L \{ \cos \omega t u_0(t) \} = L \left[ \frac{1}{\omega} \frac{d}{dt} \sin \omega t u_0(t) \right] = \frac{1}{\omega} L \left[ \frac{d}{dt} \sin \omega t u_0(t) \right] = \frac{1}{\omega} \frac{s}{s^2 + \omega^2} = \frac{s}{s^2 + \omega^2} \]
for \( \sigma > 0 \).

**Example 4.10**

Find \( \mathcal{L} \{ e^{-\alpha t} \sin \omega t \ u(t) \} \)

**Solution:**

Since

\[
\sin \omega t u(t) \Leftrightarrow \frac{\omega}{s^2 + \omega^2}
\]

using the frequency shifting property of (4.14), that is,

\[
e^{-\alpha t} f(t) \Leftrightarrow F(s + \alpha)
\]

we replace \( s \) with \( s + \alpha \), and we get

\[
e^{-\alpha t} \sin \omega t u(t) \Leftrightarrow \frac{\omega}{(s + \alpha)^2 + \omega^2}
\]

for \( \sigma > 0 \) and \( \alpha > 0 \).

**Example 4.11**

Find \( \mathcal{L} \{ e^{-\alpha t} \cos \omega t \ u(t) \} \)

**Solution:**

Since

\[
\cos \omega t u(t) \Leftrightarrow \frac{s}{s^2 + \omega^2}
\]

using the frequency shifting property of (4.14), we replace \( s \) with \( s + \alpha \), and we get

\[
e^{-\alpha t} \cos \omega t u(t) \Leftrightarrow \frac{s + \alpha}{(s + \alpha)^2 + \omega^2}
\]

for \( \sigma > 0 \) and \( \alpha > 0 \).

For easy reference, we have summarized the above derivations in Table 4.2.
### Table 4.2 Laplace Transform Pairs for Common Functions

<table>
<thead>
<tr>
<th>$f(t)$</th>
<th>$F(s)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$u_0(t)$</td>
<td>$1/s$</td>
</tr>
<tr>
<td>$tu_0(t)$</td>
<td>$1/s^2$</td>
</tr>
<tr>
<td>$t^n u_0(t)$</td>
<td>$n! \frac{s}{s^n+1}$</td>
</tr>
<tr>
<td>$\delta(t)$</td>
<td>$1$</td>
</tr>
<tr>
<td>$\delta(t-a)$</td>
<td>$e^{-as}$</td>
</tr>
<tr>
<td>$e^{-at} u_0(t)$</td>
<td>$\frac{1}{s+a}$</td>
</tr>
<tr>
<td>$t^n e^{-at} u_0(t)$</td>
<td>$\frac{n!}{(s+a)^{n+1}}$</td>
</tr>
<tr>
<td>$\sin \omega t u_0(t)$</td>
<td>$\frac{\omega}{s^2 + \omega^2}$</td>
</tr>
<tr>
<td>$\cos \omega t u_0(t)$</td>
<td>$\frac{s}{s^2 + \omega^2}$</td>
</tr>
<tr>
<td>$e^{-at} \sin \omega t u_0(t)$</td>
<td>$\frac{\omega}{(s+a)^2 + \omega^2}$</td>
</tr>
<tr>
<td>$e^{-at} \cos \omega t u_0(t)$</td>
<td>$\frac{s+a}{(s+a)^2 + \omega^2}$</td>
</tr>
</tbody>
</table>

#### 4.4 The Laplace Transform of Common Waveforms

In this section, we will present some examples for deriving the Laplace transform of several waveforms using the transform pairs of Tables 4.1 and 4.2.

**Example 4.12**

Find the Laplace transform of the waveform $f_P(t)$ of Figure 4.1. The subscript $P$ stands for pulse.
Solution:

We first express the given waveform as a sum of unit step functions. Then,

$$f_p(t) = A[u_0(t) - u_0(t - a)]$$  \hspace{1cm} (4.67)

Next, from Table 4.1,

$$f(t - a)u_0(t - a) \Leftrightarrow e^{-as}F(s)$$

and from Table 4.2,

$$u_0(t) \Leftrightarrow \frac{1}{s}$$

For this example,

$$Au_0(t) \Leftrightarrow \frac{A}{s}$$

and

$$Au_0(t - a) \Leftrightarrow \frac{e^{-as}A}{s}$$

Then, by the linearity property, the Laplace transform of the pulse of Figure 4.1 is

$$A[u_0(t) - u_0(t - a)] \Leftrightarrow \frac{A}{s}e^{-as} - \frac{A}{s}e^{-as} = \frac{A}{s}(1 - e^{-as})$$

Example 4.13

Find the Laplace transform for the waveform $f_L(t)$ of Figure 4.4. The subscript $L$ stands for line.
The Laplace Transform of Common Waveforms

Solution:

We must first derive the equation of the linear segment. This is shown in Figure 4.3. Then, we express the given waveform in terms of the unit step function.

\[ f_L(t) = (t - 1)u_0(t - 1) \]

From Table 4.1,

\[ f(t - a)u_0(t - a) \leftrightarrow e^{-as}F(s) \]

and from Table 4.2,

\[ tu_0(t) \leftrightarrow \frac{1}{s^2} \]

Therefore, the Laplace transform of the linear segment of Figure 4.2 is

\[ (t - 1)u_0(t - 1) \leftrightarrow e^{-s}\frac{1}{s^2} \quad (4.68) \]

Example 4.14

Find the Laplace transform for the triangular waveform \( f_T(t) \) of Figure 4.4.

Solution:

We must first derive the equations of the linear segments. These are shown in Figure 4.5. Then, we express the given waveform in terms of the unit step function.
For this example,

\[ f_T(t) = t[u_0(t) - u_0(t - 1)] + (-t + 2)[u_0(t - 1) - u_0(t - 2)] \]

\[ = tu_0(t) - tu_0(t - 1) - tu_0(t - 1) + 2u_0(t - 1) + tu_0(t - 2) - 2u_0(t - 2) \]

and collecting like terms,

\[ f_T(t) = tu_0(t) - 2(t - 1)u_0(t - 1) + (t - 2)u_0(t - 2) \]

From Table 4.1,

\[ f(t - a)u_0(t - a) \Leftrightarrow e^{-as}F(s) \]

and from Table 4.2,

\[ tu_0(t) \Leftrightarrow \frac{1}{s^2} \]

Then,

\[ tu_0(t) - 2(t - 1)u_0(t - 1) + (t - 2)u_0(t - 2) \Leftrightarrow \frac{1}{s^2} - 2e^{-s}\frac{1}{s^2} + e^{-2s}\frac{1}{s^2} \]

or

\[ tu_0(t) - 2(t - 1)u_0(t - 1) + (t - 2)u_0(t - 2) \Leftrightarrow \frac{1}{s^2}(1 - 2e^{-s} + e^{-2s}) \]

Therefore, the Laplace transform of the triangular waveform of Figure 4.3 is

\[ f_T(t) \Leftrightarrow \frac{1}{s^2}(1 - e^{-s})^2 \]  \hspace{1cm} (4.69)

Example 4.15

Find the Laplace transform for the rectangular periodic waveform \( f_R(t) \) of Figure 4.6.
The Laplace Transform of Common Waveforms

**Figure 4.6. Waveform for Example 4.15**

**Solution:**

This is a periodic waveform with period $T = 2a$, and thus we can apply the time periodicity property

$$
\mathcal{L} \{ f(t) \} = \frac{\int_0^T f(\tau)e^{-s\tau}d\tau}{1 - e^{-sT}}
$$

where the denominator represents the periodicity of $f(t)$. For this example,

$$
\mathcal{L} \{ f_R(t) \} = \frac{1}{1 - e^{-2as}} \int_0^{2a} f_R(t)e^{-st}dt = \frac{1}{1 - e^{-2as}} \left[ \int_0^a A e^{-st}dt + \int_0^{2a} (-A)e^{-st}dt \right]
$$

or

$$
\mathcal{L} \{ f_R(t) \} = \frac{A}{s(1 - e^{-2as})} \left( -e^{-as} + 1 + e^{-2as} - e^{-as} \right)
$$

where

$$
\begin{align*}
\mathcal{L} \{ f_R(t) \} &= \frac{A}{s(1 - e^{-2as})} \left( 1 - 2e^{-as} + e^{-2as} \right) = \frac{A(1 - e^{-as})^2}{s(1 + e^{-as})(1 - e^{-as})} \\
&= \frac{A(1 - e^{-as})}{s(1 + e^{-as})} \frac{\left( e^{as/2} e^{-as/2} - e^{-as/2} e^{as/2} \right)}{s \left( e^{as/2} e^{-as/2} + e^{-as/2} e^{as/2} \right)} \\
&= \frac{A e^{-as/2} \left( e^{as/2} - e^{-as/2} \right)}{s e^{as/2} \left( e^{as/2} + e^{-as/2} \right)} = \frac{A \sinh(as/2)}{s \cosh(as/2)}
\end{align*}
$$

or

$$
\boxed{f_R(t) \Leftrightarrow \frac{A}{s} \tanh \left( \frac{as}{2} \right)}
$$
Example 4.16

Find the Laplace transform for the half-rectified sine wave $f_{HW}(t)$ of Figure 4.7.

Solution:

This is a periodic waveform with period $T = 2\pi$. We will apply the time periodicity property

$$\mathcal{L}\{f(t)\} = \frac{\int_0^T f(\tau)e^{-s\tau}d\tau}{1 - e^{-sT}}$$

where the denominator represents the periodicity of $f(t)$. For this example,

$$\mathcal{L}\{f_{HW}(t)\} = \frac{1}{1 - e^{-2\pi s}} \int_0^{2\pi} f(t)e^{-st}dt = \frac{1}{1 - e^{-2\pi s}} \int_0^{\pi} \sin te^{-st}dt$$

$$= \frac{1}{1 - e^{-2\pi s}} \left[ -\frac{e^{-st}(s\sin t - \cos t)}{s^2 + 1} \right]_0^\pi = \frac{1}{(s^2 + 1)(1 - e^{-2\pi s})} \left( 1 + e^{-\pi s} \right)$$

or

$$\mathcal{L}\{f_{HW}(t)\} = \frac{1}{(s^2 + 1)(1 + e^{-\pi s})(1 - e^{-\pi s})}$$

(4.71)

$$f_{HW}(t) \Leftrightarrow \frac{1}{(s^2 + 1)(1 - e^{-\pi s})}$$
4.5 Summary

- The two-sided or bilateral Laplace Transform pair is defined as

\[ \mathcal{L}\{f(t)\} = F(s) = \int_{-\infty}^{\infty} f(t)e^{-st}dt \]

\[ \mathcal{L}^{-1}\{F(s)\} = f(t) = \frac{1}{2\pi j} \int_{\sigma-j\omega}^{\sigma+j\omega} F(s)e^{st}ds \]

where \( \mathcal{L}\{f(t)\} \) denotes the Laplace transform of the time function \( f(t) \), \( \mathcal{L}^{-1}\{F(s)\} \) denotes the Inverse Laplace transform, and \( s \) is a complex variable whose real part is \( \sigma \), and imaginary part \( \omega \), that is, \( s = \sigma + j\omega \).

- The unilateral or one-sided Laplace transform defined as

\[ \mathcal{L}\{f(t)\} = F(s) = \int_{0}^{\infty} f(t)e^{-st}dt \]

- We denote transformation from the time domain to the complex frequency domain, and vice versa, as

\[ f(t) \Longleftrightarrow F(s) \]

- The linearity property states that

\[ c_1 f_1(t) + c_2 f_2(t) + \ldots + c_n f_n(t) \Longleftrightarrow c_1 F_1(s) + c_2 F_2(s) + \ldots + c_n F_n(s) \]

- The time shifting property states that

\[ f(t-a)u_0(t-a) \Longleftrightarrow e^{-as}F(s) \]

- The frequency shifting property states that

\[ e^{-at}f(t) \Longleftrightarrow F(s + a) \]

- The scaling property states that

\[ f(at) \Longleftrightarrow \frac{1}{a} F\left(\frac{s}{a}\right) \]

- The differentiation in time domain property states that

\[ f'(t) = \frac{d}{dt} f(t) \Longleftrightarrow sF(s) - f(0^-) \]
\[
\frac{d^2}{dt^2}f(t) \Leftrightarrow s^2F(s) - sf(0^-) - f'(0^-)
\]
\[
\frac{d^3}{dt^3}f(t) \Leftrightarrow s^3F(s) - s^2f(0^-) - sf'(0^-) - f''(0^-)
\]
and in general
\[
\frac{d^n}{dt^n}f(t) \Leftrightarrow s^nF(s) - s^{n-1}f(0^-) - s^{n-2}f'(0^-) - \ldots - f^{n-1}(0^-)
\]
where the terms \(f(0^-), f'(0^-), f''(0^-)\), and so on, represent the initial conditions.

- The differentiation in complex frequency domain property states that
  \[
tf(t) \Leftrightarrow -\frac{d}{ds}F(s)
  \]
  and in general,
  \[
t^n f(t) \Leftrightarrow (-1)^n \frac{d^n}{ds^n}F(s)
  \]
- The integration in time domain property states that
  \[
  \int_{-\infty}^{t} f(\tau)d\tau \Leftrightarrow \frac{F(s)}{s} + \frac{f(0^-)}{s}
  \]
- The integration in complex frequency domain property states that
  \[
  \frac{f(t)}{t} \Leftrightarrow \int_{s}^{\infty} F(s)ds
  \]
  provided that the limit \(\lim_{t \to 0} \frac{f(t)}{t}\) exists.
- The time periodicity property states that
  \[
  f(t + nT) \Leftrightarrow \int_{0}^{T} f(t)e^{-st}dt
  \]
  \[
  1 - e^{-sT}
  \]
- The initial value theorem states that
  \[
  \lim_{t \to 0} f(t) = \lim_{s \to \infty} sF(s) = f(0^-)
  \]
The final value theorem states that
\[
\lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = f(\infty)
\]

Convolution in the time domain corresponds to multiplication in the complex frequency domain, that is,
\[
f_1(t) \ast f_2(t) \Leftrightarrow F_1(s)F_2(s)
\]

Convolution in the complex frequency domain divided by \(1/2\pi j\), corresponds to multiplication in the time domain. That is,
\[
f_1(t)f_2(t) \Leftrightarrow \frac{1}{2\pi j} F_1(s)F_2(s)
\]

The Laplace transforms of some common functions of time are shown below.

- \(u_0(t) \Leftrightarrow 1/s\)
- \(t \Leftrightarrow 1/s^2\)
- \(t^n u_0(t) \Leftrightarrow \frac{n!}{s^{n+1}}\)
- \(\delta(t) \Leftrightarrow 1\)
- \(\delta(t-a) \Leftrightarrow e^{-as}\)
- \(e^{-at}u_0(t) \Leftrightarrow \frac{1}{s+a}\)
- \(te^{-at}u_0(t) \Leftrightarrow \frac{1}{(s+a)^2}\)
- \(t^2 e^{-at}u_0(t) \Leftrightarrow \frac{2!}{(s+a)^3}\)
- \(t^n e^{-at}u_0(t) \Leftrightarrow \frac{n!}{(s+a)^{n+1}}\)
- \(\sin \omega t u_0 t \Leftrightarrow \frac{\omega}{s^2 + \omega^2}\)
- \(\cos \omega t u_0 t \Leftrightarrow \frac{s}{s^2 + \omega^2}\)
The Laplace transforms of some common waveforms are shown below.

\[ e^{-at}\sin\omega t \leftrightarrow \frac{\omega}{(s+a)^2 + \omega^2} \]

\[ e^{-at}\cos\omega t \leftrightarrow \frac{s+a}{(s+a)^2 + \omega^2} \]

- The Laplace transforms of some common waveforms are shown below.

\[ A[u_0(t) - u_0(t-a)] \leftrightarrow \frac{A}{s}e^{-asA} = \frac{A}{s}(1 - e^{-as}) \]

\[ (t-1)u_0(t-1) \leftrightarrow e^{-s\frac{1}{2}} \]

\[ f_T(t) \leftrightarrow \frac{1}{s^2}(1 - e^{-s})^2 \]

\[ f_R(t) \leftrightarrow \frac{A}{s}\tanh\left(\frac{as}{2}\right) \]
\[ f_{HW}(t) \leftrightarrow \frac{1}{(s^2 + 1)(1 - e^{-\pi s})} \]
Chapter 4  The Laplace Transformation

4.6 Exercises

1. Find the Laplace transform of the following time domain functions:
   a. \(12\)
   b. \(6u_0(t)\)
   c. \(24u_0(t - 12)\)
   d. \(5tu_0(t)\)
   e. \(4t^5u_0(t)\)

2. Find the Laplace transform of the following time domain functions:
   a. \(j8\)
   b. \(j5 \angle -90^\circ\)
   c. \(5e^{-5t}u_0(t)\)
   d. \(8t^7e^{-5t}u_0(t)\)
   e. \(15\delta(t - 4)\)

3. Find the Laplace transform of the following time domain functions:
   a. \((t^3 + 3t^2 + 4t + 3)u_0(t)\)
   b. \(3(2t - 3)\delta(t - 3)\)
   c. \((3\sin 5t)u_0(t)\)
   d. \((5\cos 3t)u_0(t)\)
   e. \((2\tan 4t)u_0(t)\) Be careful with this! Comment and skip derivation.

4. Find the Laplace transform of the following time domain functions:
   a. \(3t(sin 5t)u_0(t)\)
   b. \(2t^2(cos 3t)u_0(t)\)
   c. \(2e^{-5t}\sin 5t\)
5. Find the Laplace transform of the following time domain functions:
   a. $5tu_0(t - 3)$
   b. $(2t^2 - 5t + 4)u_0(t - 3)$
   c. $(t - 3)e^{-2t}u_0(t - 2)$
   d. $(2t - 4)e^{2(t-2)}u_0(t - 3)$
   e. $4te^{-3t}(\cos 2t)u_0(t)$

6. Find the Laplace transform of the following time domain functions:
   a. $\frac{d}{dt}(\sin 3t)$
   b. $\frac{d}{dt}(3e^{-4t})$
   c. $\frac{d}{dt}(t^2 \cos 2t)$
   d. $\frac{d}{dt}(e^{-2t}\sin 2t)$
   e. $\frac{d}{dt}(t^2e^{-2t})$

7. Find the Laplace transform of the following time domain functions:
   a. $\frac{\sin t}{t}$
   b. $\int_{0}^{t} \frac{\sin \tau}{\tau} d\tau$
   c. $\frac{\sin at}{t}$
   d. $\int_{t}^{\infty} \frac{\cos \tau}{\tau} d\tau$
8. Find the Laplace transform for the sawtooth waveform $f_{ST}(t)$ of Figure 4.8.

![Figure 4.8. Waveform for Exercise 8.]

9. Find the Laplace transform for the full rectification waveform $f_{FR}(t)$ of Figure 4.9.

![Figure 4.9. Waveform for Exercise 9]
4.7 Solutions to Exercises

1. From the definition of the Laplace transform or from Table 4.2 we get:
   a. $12/s$  
   b. $6/s$  
   c. $e^{-12s} \cdot \frac{24}{s}$  
   d. $5/s^2$  
   e. $4 \cdot \frac{5!}{s^6}$

2. From the definition of the Laplace transform or from Table 4.2 we get:
   a. $j8/s$  
   b. $5/s$  
   c. $\frac{5}{s+5}$  
   d. $8 \cdot \frac{7!}{(s+5)^8}$  
   e. $15e^{-4s}$

3. 
   a. From Table 4.2 and the linearity property $\frac{3!}{s^4} + \frac{3 \times 2!}{s^3} + \frac{4}{s^2} + \frac{3}{s}$
   b. $3(2t-3)\delta(t-3) = 3(2t-3)|_{t=3} = 3\delta(t-3) - 9\delta(t-3)$ and $9\delta(t-3) \Leftrightarrow 9e^{-3s}$
   c. $3 \cdot \frac{5}{s^2 + 5^2}$  
   d. $5 \cdot \frac{s}{s^2 + 3^2}$  
   e. $2\tan 4t = 2 \cdot \frac{\sin 4t}{\cos 4t} \Leftrightarrow 2 \cdot \frac{4/(s^2 + 2^2)}{s/(s^2 + 2^2)} = \frac{8}{s}$. This answer looks suspicious because $8/s \Leftrightarrow 8u_0(t)$ and the Laplace transform is unilateral, that is, there is one-to-one correspondence between the time domain and the complex frequency domain. The fallacy with this procedure is that if we assume that $f_1(t) \Leftrightarrow F_1(s)$ and $f_2(t) \Leftrightarrow F_2(s)$, we cannot conclude that $\frac{f_1(t)}{f_2(t)} \Leftrightarrow \frac{F_1(s)}{F_2(s)}$. For this exercise $f_1(t) \cdot f_2(t) = \sin 4t \cdot \frac{1}{\cos 4t}$ and as we’ve learned multiplication in the time domain corresponds to convolution in the complex frequency domain. Accordingly, we must use the Laplace transform definition $\int_0^\infty (2\tan 4t)e^{-st}dt$ and this requires integration by parts. We skip this analytical derivation. The interested reader may try to find the answer with the MATLAB code `syms s t; 2*laplace(sin(4*t)/cos(4*t))`

4. From (4.22)

\[ t^n f(t) \Leftrightarrow (-1)^n \frac{d^n}{ds^n} F(s) \]

Then,

a. 

\[ 3(-1)^3 \frac{d}{ds} \left( \frac{5}{s^2 + 5^2} \right) = -3 \left[ \frac{-5 \cdot (2s)}{(s^2 + 25)^2} \right] = \frac{30s}{(s^2 + 25)^2} \]
b. 

\[ 2(-1)^2 \frac{d^2}{ds^2} \left( \frac{s}{s^2 + 3^2} \right) = 2 \frac{d}{ds} \left[ \frac{s^2 + 3^2 - s(2s)}{(s^2 + 9)^2} \right] = 2 \frac{d}{ds} \left( \frac{-s^2 + 9}{(s^2 + 9)^2} \right) \]

\[ = 2 \cdot \frac{(s^2 + 9)^2(-2s) - 2(s^2 + 9)(2s)(-s^2 + 9)}{(s^2 + 9)^4} \]

\[ = 2 \cdot \frac{(s^2 + 9)(-2s) - 4s(-s^2 + 9)}{(s^2 + 9)^3} = 2 \cdot \frac{-2s^3 - 18s + 4s^3 - 36s}{(s^2 + 9)^3} \]

\[ = 2 \cdot \frac{2s^3 - 54s}{(s^2 + 9)^3} = 2 \cdot \frac{2s(s^2 - 27)}{(s^2 + 9)^3} = \frac{4s(s^2 - 27)}{(s^2 + 9)^3} \]

c. 

\[ \frac{2 \times 5}{(s + 5)^2 + 5^2} = \frac{10}{(s + 5)^2 + 25} \]

d. 

\[ \frac{8(s + 3)}{(s + 3)^2 + 4^2} = \frac{8(s + 3)}{(s + 3)^2 + 16} \]

e. 

\[ \cos t \big|_{\pi/4} \delta(t - \pi/4) = (\sqrt{2}/2) \delta(t - \pi/4) \text{ and } (\sqrt{2}/2) \delta(t - \pi/4) \Leftrightarrow (\sqrt{2}/2) e^{-(\pi/4)s} \]

5. 

a. 

\[ 5tu_o(t - 3) = [5(t - 3) + 15]u_o(t - 3) \Leftrightarrow e^{-3s} \left( \frac{5}{s} + \frac{15}{s^2} \right) = \frac{5}{s} e^{-3s} \left( \frac{1}{s^2} + 3 \right) \]

b. 

\[ (2t^2 - 5t + 4)u_o(t - 3) = [2(t - 3)^2 + 12t - 18 - 5t + 4]u_o(t - 3) \]

\[ = [2(t - 3)^2 + 7t - 14]u_o(t - 3) \]

\[ = [2(t - 3)^2 + 7(t - 3) + 21 - 14]u_o(t - 3) \]

\[ = [2(t - 3)^2 + 7(t - 3) + 7]u_o(t - 3) \Leftrightarrow e^{-3s} \left( \frac{2\times 2!}{s^3} + \frac{7}{s^2} + \frac{7}{s} \right) \]

c. 

\[ (t - 3)e^{-2t}u_o(t - 2) = [(t - 2) - 1]e^{-2(t - 2)} \cdot e^{-4}u_o(t - 2) \]

\[ \Leftrightarrow e^{-4} \cdot e^{-2t} \left[ \frac{1}{(s + 2)^2} - \frac{1}{(s + 2)} \right] = e^{-4} \cdot e^{-2} \left[ \frac{-(s + 1)}{(s + 2)^2} \right] \]
d. 
\[(2t - 4)e^{2(t-2)}u_0(t-3) = [2(t-3) + 6 - 4]e^{-2(t-3)} \cdot e^{-2}u_0(t-3)\]
\[\iff e^{-2} \cdot e^{-3s} \left[ \frac{2}{(s + 3)^2} + \frac{2}{(s + 3)} \right] = 2e^{-2} \cdot e^{-3s} \left[ \frac{s + 4}{-(s + 3)^2} \right]\]

e. 
\[4te^{-3t}(\cos 2t)u_0(t) \iff 4(-1)^t \frac{d}{ds} \left[ \frac{s + 3}{(s + 3)^2 + 2^2} \right] = -4 \frac{d}{ds} \left[ \frac{s + 3}{s^2 + 6s + 9 + 4} \right]\]
\[\iff -4 \frac{d}{ds} \left[ \frac{s + 3}{s^2 + 6s + 13} \right] = -4 \frac{s^2 + 6s + 13 - (s + 3)(2s + 6)}{(s^2 + 6s + 13)^2}\]
\[\iff -4 \frac{s^2 + 6s + 13 - 2s^2 - 6s - 6 - 18}{(s^2 + 6s + 13)^2} = \frac{4(s^2 + 6s + 5)}{(s^2 + 6s + 13)^2}\]

6.

a. 
\[\sin 3t \iff \frac{3}{s^2 + 3^2} \quad \frac{d}{dt} f(t) \iff sF(s) - f(0^-) \quad f(0^-) = \sin 3t \big|_{t=0} = 0\]
\[\frac{d}{dt} (\sin 3t) \iff s \frac{3}{s^2 + 3^2} - 0 = \frac{3s}{s^2 + 9}\]

b. 
\[3e^{-4t} \iff \frac{3}{s + 4} \quad \frac{d}{dt} f(t) \iff sF(s) - f(0^-) \quad f(0^-) = 3e^{-4t} \big|_{t=0} = 3\]
\[\frac{d}{dt} (3e^{-4t}) \iff s \frac{3}{s + 4} - 3 = \frac{3s}{s + 4} - \frac{3(s + 4)}{s + 4} = -\frac{12}{s + 4}\]

c. 
\[\cos 2t \iff \frac{s}{s^2 + 2^2} \quad t^2 \cos 2t \iff (-1)^2 \frac{d^2}{ds^2} \left[ \frac{s}{s^2 + 4} \right]\]
\[\frac{d}{ds} \left[ \frac{s^2 + 4 - s(2s)}{(s^2 + 4)^2} \right] = \frac{d}{ds} \left[ \frac{-s^2 + 4}{(s^2 + 4)} \right] = \frac{(s^2 + 4)(-2s) - (-s^2 + 4)(s^2 + 4)2(2s)}{(s^2 + 4)^4}\]
\[= \frac{(s^2 + 4)(-2s) - (-s^2 + 4)(4s)}{(s^2 + 4)^3} = -\frac{2s^3 - 8s + 4s^3 - 16s}{(s^2 + 4)^3} = \frac{2s(s^2 - 12)}{(s^2 + 4)^3}\]

Thus,
\[t^2 \cos 2t \iff \frac{2s(s^2 - 12)}{(s^2 + 4)^3}\]
and

\[ \frac{d}{dt}(t^2 \cos 2t) \Leftrightarrow sF(s) - f(0^-) \]

\[ \Leftrightarrow s \frac{2s(s^2 - 12)}{(s^2 + 4)^3} - 0 = \frac{2s^2(s^2 - 12)}{(s^2 + 4)^3} \]

d. \[ \sin 2t \Leftrightarrow \frac{2}{s^2 + 2^2} \quad e^{-2t} \sin 2t \Leftrightarrow \frac{2}{(s + 2)^2 + 4} \quad \frac{d}{dt} f(t) \Leftrightarrow sF(s) - f(0^-) \]

\[ \frac{d}{dt}(e^{-2t} \sin 2t) \Leftrightarrow s \frac{2}{(s + 2)^2 + 4} - 0 = \frac{2s}{(s + 2)^2 + 4} \]

e. \[ t^2 \Leftrightarrow \frac{2!}{s^3} \quad t^2 e^{-2t} \Leftrightarrow \frac{2!}{(s + 2)^3} \quad \frac{d}{dt} f(t) \Leftrightarrow sF(s) - f(0^-) \]

\[ \frac{d}{dt}(t^2 e^{-2t}) \Leftrightarrow s \frac{2!}{(s + 2)^3} - 0 = \frac{2s}{(s + 2)^3} \]

7. a. \[ \sin t \Leftrightarrow \frac{1}{s^2 + 1} \text{ but to find } \mathcal{L} \left\{ \frac{\sin t}{t} \right\} \text{ we must show that the limit } \lim_{t \to 0} \frac{\sin t}{t} \text{ exists. Since} \]

\[ \lim_{x \to 0} \frac{\sin x}{x} = 1 \text{ this condition is satisfied and thus } \frac{\sin t}{t} \Leftrightarrow \int_{\infty}^{\infty} \frac{1}{s} \frac{1}{s^2 + 1} \, ds. \text{ From tables of integrals} \]

\[ \int_{\infty}^{\infty} \frac{1}{x^2 + a^2} \, dx = \frac{1}{a} \tan^{-1}(x/a) + C. \text{ Then, } \int_{\infty}^{\infty} \frac{1}{s^2 + 1} \, ds = \tan^{-1}(1/s) + C \text{ and the constant of integration } C \text{ is evaluated from the final value theorem. Thus,} \]

\[ \lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = \lim_{s \to 0} s[\tan^{-1}(1/s) + C] = 0 \text{ and } \frac{\sin t}{t} \Leftrightarrow \tan^{-1}(1/s) \]

b. From (a) above \[ \frac{\sin t}{t} \Leftrightarrow \tan^{-1}(1/s) \text{ and since } \int_{\infty}^{t} f(\tau) \, d\tau \Leftrightarrow \frac{F(s)}{s} + \frac{f(0^-)}{s} \text{, it follows that} \]

\[ \int_{0}^{t} \frac{\sin \tau}{\tau} \, d\tau \Leftrightarrow \frac{1}{s} \tan^{-1}(1/s) \]
c. 
From (a) above \( \frac{\sin t}{t} \Leftrightarrow \tan^{-1}(1/s) \) and since \( f(at) \Leftrightarrow \frac{1}{a}F\left(\frac{s}{a}\right) \), it follows that

\[
\frac{\sin at}{at} \Leftrightarrow \frac{1}{a} \tan^{-1}\left(\frac{1/s}{a}\right) \quad \text{or} \quad \frac{\sin t}{t} \Leftrightarrow \tan^{-1}(a/s)
\]

d.
\( \cos t \Leftrightarrow \frac{s}{s^2 + 1}, \quad \frac{\cos t}{t} \Leftrightarrow \int_0^{\infty} \frac{s}{s^2 + 1} ds \), and from tables of integrals

\[
\int_0^{\infty} \frac{x}{x^2 + a^2} dx = \frac{1}{2} \ln(x^2 + a^2) + C 
\]

Then, \( \int_0^{\infty} \frac{s}{s^2 + 1} ds = \frac{1}{2} \ln(s^2 + 1) + C \) and the constant of integration \( C \) is evaluated from the final value theorem. Thus,

\[
\lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = \lim_{s \to 0} s \left[ \frac{1}{2} \ln(s^2 + 1) + C \right] = 0 
\]

and using \( \int_{-\infty}^{t} f(\tau) d\tau \Leftrightarrow \frac{F(s)}{s} + \frac{f(0^-)}{s} \) we get

\[
\int_t^{\infty} \frac{\cos \tau}{\tau} d\tau \Leftrightarrow \frac{1}{2s} \ln(s^2 + 1)
\]

e. 
\( e^{-t} \Leftrightarrow \frac{1}{s + 1}, \quad \frac{e^{-t}}{t} \Leftrightarrow \int_0^{\infty} \frac{1}{s + 1} ds \), and from tables of integrals

\[
\int_0^{\infty} \frac{1}{ax + b} dx = \frac{1}{2} \ln(ax + b). \quad \text{Then,} \quad \int_0^{\infty} \frac{1}{s + 1} ds = \ln(s + 1) + C \] and the constant of integration \( C \) is evaluated from the final value theorem. Thus,

\[
\lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = \lim_{s \to 0} s [\ln(s + 1) + C] = 0 \quad \text{and using} \quad \int_{-\infty}^{t} f(\tau) d\tau \Leftrightarrow \frac{F(s)}{s} + \frac{f(0^-)}{s} \quad \text{we get}
\]

\[
\int_t^{\infty} \frac{e^{-\tau}}{\tau} d\tau \Leftrightarrow \frac{1}{s} \ln(s + 1)
\]

8. 
This is a periodic waveform with period \( T = a \) and its Laplace transform is

\[
f_{ST}(t) \quad \frac{A}{a} \quad \frac{A}{a} \quad \frac{A}{a} 
\]

\[
t \quad a \quad 2a \quad 3a
\]
\[ F(s) = \frac{l}{1 - e^{-as}} \int_{0}^{\varphi} A e^{-st} dt = \frac{A}{a(1 - e^{-as})} \int_{0}^{\varphi} e^{-st} dt \quad (1) \]

and from (4.40) of Example 4.2 and limits of integration \( \varphi \) to \( a \) we get

\[ \mathcal{L} \{ t \} \bigg|_{0}^{a} = \int_{0}^{\varphi} e^{-st} dt = \left[ -\frac{te^{-st}}{s} - \frac{e^{-st}}{s^2} \right]_{0}^{\varphi} = \left[ \frac{te^{-st}}{s} + \frac{e^{-st}}{s^2} \right]_{a}^{\varphi} \]

\[ = \left[ \frac{l}{s^2} - \frac{ae^{-as}}{s} - \frac{e^{-as}}{s^2} \right] = \frac{l}{s^2} \left[ 1 - (1 + as)e^{-as} \right] \]

Adding and subtracting \( as \) we get

\[ \mathcal{L} \{ t \} \bigg|_{0}^{a} = \frac{l}{s^2} [(1 + as) - (1 + as)e^{-as} - as] = \frac{l}{s^2} [(1 + as)(1 - e^{-as}) - as] \]

By substitution into (1) we get

\[ F(s) = \frac{A}{a(1 - e^{-as})} \cdot \frac{l}{s^2} [(1 + as)(1 - e^{-as}) - as] = \frac{A}{as^2(1 - e^{-as})} \cdot [(1 + as)(1 - e^{-as}) - as] \]

\[ = \frac{A(1 + as)}{as^2} - \frac{Aa}{as(1 - e^{-as})} = \frac{A}{as} \left[ \frac{(1 + as)}{s} - \frac{a}{(1 - e^{-as})} \right] \]

9.

This is a periodic waveform with period \( T = a = \pi \) and its Laplace transform is

\[ F(s) = \frac{l}{1 - e^{-\pi s}} \int_{0}^{\pi} f(t) e^{-st} dt = \frac{1}{(1 - e^{-\pi s})} \int_{0}^{\pi} \sin t e^{-st} dt \]

From tables of integrals

\[ \int sin b x e^{ax} dx = \frac{e^{ax}(asinh x - bcos b x)}{a^2 + b^2} \]

Then,

\[ F(s) = \frac{l}{1 - e^{-\pi s}} \cdot \frac{e^{-st}(ssint - cost)}{s^2 + 1} \bigg|_{0}^{\pi} = \frac{l}{1 - e^{-\pi s}} \cdot \frac{1 + e^{-\pi s}}{s^2 + 1} \]

\[ = \frac{l}{s^2 + 1} \cdot \frac{1 + e^{-\pi s}}{1 - e^{-\pi s}} = \frac{l}{s^2 + 1} \coth \left( \frac{\pi s}{2} \right) \]
Chapter 5

The Inverse Laplace Transformation

This chapter is a continuation to the Laplace transformation topic of the previous chapter and presents several methods of finding the Inverse Laplace Transformation. The partial fraction expansion method is explained thoroughly and it is illustrated with several examples.

5.1 The Inverse Laplace Transform Integral

The Inverse Laplace Transform Integral was stated in the previous chapter; it is repeated here for convenience.

\[
\mathcal{L}^{-1}\{F(s)\} = f(t) = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F(s)e^{st} ds \tag{5.1}
\]

This integral is difficult to evaluate because it requires contour integration using complex variables theory. Fortunately, for most engineering problems we can refer to Tables of Properties, and Common Laplace transform pairs to lookup the Inverse Laplace transform.

5.2 Partial Fraction Expansion

Quite often the Laplace transform expressions are not in recognizable form, but in most cases appear in a rational form of \( s \), that is,

\[
F(s) = \frac{N(s)}{D(s)} \tag{5.2}
\]

where \( N(s) \) and \( D(s) \) are polynomials, and thus (5.2) can be expressed as

\[
F(s) = \frac{N(s)}{D(s)} = \frac{b_m s^m + b_{m-1}s^{m-1} + b_{m-2}s^{m-2} + \ldots + b_1 s + b_0}{a_n s^n + a_{n-1}s^{n-1} + a_{n-2}s^{n-2} + \ldots + a_1 s + a_0} \tag{5.3}
\]

The coefficients \( a_k \) and \( b_k \) are real numbers for \( k = 1, 2, \ldots, n \), and if the highest power \( m \) of \( N(s) \) is less than the highest power \( n \) of \( D(s) \), i.e., \( m < n \), \( F(s) \) is said to be expressed as a proper rational function. If \( m \geq n \), \( F(s) \) is an improper rational function.

In a proper rational function, the roots of \( N(s) \) in (5.3) are found by setting \( N(s) = 0 \); these are called the zeros of \( F(s) \). The roots of \( D(s) \), found by setting \( D(s) = 0 \), are called the poles of \( F(s) \). We assume that \( F(s) \) in (5.3) is a proper rational function. Then, it is customary and very convenient...
to make the coefficient of \( s^n \) unity; thus, we rewrite \( F(s) \) as

\[
F(s) = \frac{N(s)}{D(s)} = \frac{1}{a_n} \frac{(b_m s^m + b_{m-1} s^{m-1} + b_{m-2} s^{m-2} + \ldots + b_1 s + b_0)}{s^n + \frac{a_{n-1} s^{n-1}}{a_n} + \frac{a_{n-2} s^{n-2}}{a_n} + \ldots + \frac{a_1}{a_n} + \frac{a_0}{a_n}}
\] (5.4)

The zeros and poles of (5.4) can be real and distinct, or repeated, or complex conjugates, or combinations of real and complex conjugates. However, we are mostly interested in the nature of the poles, so we will consider each case separately.

**Case I: Distinct Poles**

If all the poles \( p_1, p_2, p_3, \ldots, p_n \) of \( F(s) \) are distinct (different from each another), we can factor the denominator of \( F(s) \) in the form

\[
F(s) = \frac{N(s)}{(s-p_1) \cdot (s-p_2) \cdot (s-p_3) \cdot \ldots \cdot (s-p_n)}
\] (5.5)

where \( p_k \) is distinct from all other poles. Next, using the *partial fraction expansion method*, *we can express (5.5) as

\[
F(s) = \frac{r_1}{(s-p_1)} + \frac{r_2}{(s-p_2)} + \frac{r_3}{(s-p_3)} + \ldots + \frac{r_n}{(s-p_n)}
\] (5.6)

where \( r_1, r_2, r_3, \ldots, r_n \) are the residues, and \( p_1, p_2, p_3, \ldots, p_n \) are the poles of \( F(s) \).

To evaluate the residue \( r_k \), we multiply both sides of (5.6) by \((s-p_k)\); then, we let \( s \to p_k \), that is,

\[
r_k = \lim_{s \to p_k} (s-p_k)F(s) = (s-p_k)F(s)\bigg|_{s=p_k}
\] (5.7)

**Example 5.1**

Use the partial fraction expansion method to simplify \( F_1(s) \) of (5.8) below, and find the time domain function \( f_1(t) \) corresponding to \( F_1(s) \).

\[
F_1(s) = \frac{3s + 2}{s^2 + 3s + 2}
\] (5.8)

* The partial fraction expansion method applies only to proper rational functions. It is used extensively in integration, and in finding the inverses of the Laplace transform, the Fourier transform, and the z-transform. This method allows us to decompose a rational polynomial into smaller rational polynomials with simpler denominators from which we can easily recognize their integrals and inverse transformations. This method is also being taught in intermediate algebra and introductory calculus courses.
Solution:

Using (5.6), we get

\[ F_1(s) = \frac{3s + 2}{s^2 + 3s + 2} = \frac{3s + 2}{(s + 1)(s + 2)} = \frac{r_1}{s + 1} + \frac{r_2}{s + 2} \quad (5.9) \]

The residues are

\[ r_1 = \lim_{s \to -1} (s + 1)F(s) = \frac{3s + 2}{(s + 2)} \bigg|_{s = -1} = -1 \quad (5.10) \]

and

\[ r_2 = \lim_{s \to -2} (s + 2)F(s) = \frac{3s + 2}{(s + 1)} \bigg|_{s = -2} = 4 \quad (5.11) \]

Therefore, we express (5.9) as

\[ F_1(s) = \frac{3s + 2}{s^2 + 3s + 2} = \frac{-1}{s + 1} + \frac{4}{s + 2} \quad (5.12) \]

and from Table 4.2 of Chapter 4

\[ e^{-at}u_0(t) \Leftrightarrow \frac{1}{s + a} \quad (5.13) \]

Then,

\[ F_1(s) = \frac{-1}{s + 1} + \frac{4}{s + 2} \Leftrightarrow (-e^{-t} + 4e^{-2t})u_0(t) = f_1(t) \quad (5.14) \]

The residues and poles of a rational function of polynomials such as (5.8), can be found easily using the MATLAB \texttt{residue(a,b)} function. For this example, we use the code

\[ \text{Ns} = [3, 2]; \text{Ds} = [1, 3, 2]; [r, p, k] = \text{residue(Ns, Ds)} \]

and MATLAB returns the values

\[ r = 4 \\
-1 \]

\[ p = -2 \\
-1 \]

\[ k = [ ] \]
For this MATLAB code, we defined \( N_s \) and \( D_s \) as two vectors that contain the numerator and denominator coefficients of \( F(s) \). When this code is executed, MATLAB displays the \( r \) and \( p \) vectors that represent the residues and poles respectively. The first value of the vector \( r \) is associated with the first value of the vector \( p \), the second value of \( r \) is associated with the second value of \( p \), and so on.

The vector \( k \) is referred to as the direct term and it is always empty (has no value) whenever \( F(s) \) is a proper rational function, that is, when the highest degree of the denominator is larger than that of the numerator. For this example, we observe that the highest power of the denominator is \( s^2 \), whereas the highest power of the numerator is \( s \) and therefore the direct term is empty.

We can also use the MATLAB \texttt{ilaplace(f)} function to obtain the time domain function directly from \( F(s) \). This is done with the code that follows.

```matlab
syms s t; Fs=(3*s+2)/(s^2+3*s+2); ft=ilaplace(Fs); pretty(ft)
```

When this code is executed, MATLAB displays the expression

\[
4 \exp(-2\ t) - \exp(-t)
\]

Example 5.2

Use the partial fraction expansion method to simplify \( F_2(s) \) of (5.15) below, and find the time domain function \( f_2(t) \) corresponding to \( F_2(s) \).

\[
F_2(s) = \frac{3s^2 + 2s + 5}{s^3 + 12s^2 + 44s + 48} \quad \text{(5.15)}
\]

Solution:

First, we use the MATLAB \texttt{factor(s)} symbolic function to express the denominator polynomial of \( F_2(s) \) in factored form. For this example,

```matlab
syms s; factor(s^3 + 12*s^2 + 44*s + 48)
```

\[
\text{ans} = \frac{(s+2)(s+4)(s+6)}{s^3 + 12s^2 + 44s + 48}
\]

Then,

\[
F_2(s) = \frac{3s^2 + 2s + 5}{(s+2)(s+4)(s+6)} = \frac{r_1}{s+2} + \frac{r_2}{s+4} + \frac{r_3}{s+6} \quad \text{(5.16)}
\]

The residues are
Then, by substitution into (5.16) we get

$$F_2(s) = \frac{3s^2 + 2s + 5}{s^3 + 12s^2 + 44s + 48} = \frac{9/8}{s + 2} + \frac{-37/4}{s + 4} + \frac{89/8}{s + 6}$$  \hspace{1cm} (5.20)$$

From Table 2.2 of Chapter 2

$$e^{-at} u_0(t) \Leftrightarrow \frac{1}{s + a}$$  \hspace{1cm} (5.21)$$

Then,

$$F_2(s) = \frac{9/8}{s + 2} + \frac{-37/4}{s + 4} + \frac{89/8}{s + 6} \Leftrightarrow \left(\frac{9}{8}e^{-2t} - \frac{37}{4}e^{-4t} + \frac{89}{8}e^{-6t}\right)u_0(t) = f_2(t)$$  \hspace{1cm} (5.22)$$

Check with MATLAB:

```matlab
syms s t; Fs = (3*s^2 + 4*s + 5) / (s^3 + 12*s^2 + 44*s + 48); ft = ilaplace(Fs)
ft =
-37/4*exp(-4*t) + 9/8*exp(-2*t) + 89/8*exp(-6*t)
```

**Case II: Complex Poles**

Quite often, the poles of $F(s)$ are complex*, and since complex poles occur in complex conjugate pairs, the number of complex poles is even. Thus, if $p_k$ is a complex root of $D(s)$, then, its complex conjugate pole, denoted as $p_k^*$, is also a root of $D(s)$. The partial fraction expansion method can also be used in this case, but it may be necessary to manipulate the terms of the expansion in order to express them in a recognizable form. The procedure is illustrated with the following example.

**Example 5.3**

Use the partial fraction expansion method to simplify $F_3(s)$ of (5.23) below, and find the time

*$A$ review of complex numbers is presented in Appendix B of Circuit Analysis I with MATLAB Applications.
domain function \( f_3(t) \) corresponding to \( F_3(s) \).

\[
F_3(s) = \frac{s + 3}{s^3 + 5s^2 + 12s + 8} \quad (5.23)
\]

Solution:

Let us first express the denominator in factored form to identify the poles of \( F_3(s) \) using the MATLAB `factor(s)` function. Then,

\[
\text{syms } s; \quad \text{factor}(s^3 + 5s^2 + 12s + 8)
\]

\[
\text{ans} = (s+1)*(s^2+4*s+8)
\]

The `factor(s)` function did not factor the quadratic term. We will use the `roots(p)` function.

\[
p=[1 \ 4 \ 8]; \quad \text{roots}_p=\text{roots}(p)
\]

\[
\text{roots}_p =
\begin{align*}
-2.0000 + 2.0000i \\
-2.0000 - 2.0000i
\end{align*}
\]

Then,

\[
F_3(s) = \frac{s + 3}{s^3 + 5s^2 + 12s + 8} = \frac{s + 3}{(s + 1)(s + 2 + j2)(s + 2 - j2)}
\]

or

\[
F_3(s) = \frac{s + 3}{s^3 + 5s^2 + 12s + 8} = \frac{r_1}{(s + 1)} + \frac{r_2}{(s + 2 + j2)} + \frac{r_2^*}{(s + 2 - j2)} \quad (5.24)
\]

The residues are

\[
r_1 = \frac{s + 3}{s^2 + 4s + 8} \bigg|_{s = -1} = \frac{2}{5} \quad (5.25)
\]

\[
r_2 = \frac{s + 3}{(s + 1)(s + 2 - j2)} \bigg|_{s = -2 - j2} = \frac{1 - j2}{(-1 - j2)(-j4)} = \frac{1 - j2}{-8 + j4} \quad (5.26)
\]

\[
= \frac{(1 - j2)(-8 - j4)}{(-8 + j4)(-8 - j4)} = \frac{-16 + j12}{80} = -\frac{1}{5} + \frac{j3}{20}
\]

\[
r_2^* = (-\frac{1}{5} + \frac{j3}{20})^* = -\frac{1}{5} - \frac{j3}{20} \quad (5.27)
\]

By substitution into (5.24),

\[
\]
Partial Fraction Expansion

\[ F_3(s) = \frac{2/5}{s+1} + \frac{-1/5 + j3/20}{(s+2+j2)} + \frac{-1/5 - j3/20}{(s+2-j2)} \]  
(5.28)

The last two terms on the right side of (5.28), do not resemble any Laplace transform pair that we derived in Chapter 2. Therefore, we will express them in a different form. We combine them into a single term*, and now (5.28) is written as

\[ F_3(s) = \frac{2/5}{s+1} - \frac{1}{5} \cdot \frac{(2s+1)}{(s^2 + 4s + 8)} \]  
(5.29)

For convenience, we denote the first term on the right side of (5.29) as \( F_{31}(s) \), and the second as \( F_{32}(s) \). Then,

\[ F_{31}(s) = \frac{2/5}{s+1} \Leftrightarrow \frac{2}{5} e^{-t} = f_{31}(t) \]  
(5.30)

Next, for \( F_{32}(s) \)

\[ F_{32}(s) = -\frac{1}{5} \cdot \frac{(2s+1)}{(s^2 + 4s + 8)} \]  
(5.31)

and recalling that

\[ e^{-at} \sin \omega t_u t \Leftrightarrow \frac{\omega}{(s+a)^2 + \omega^2} \]  
(5.32)

\[ e^{-at} \cos \omega t_u t \Leftrightarrow \frac{s+a}{(s+a)^2 + \omega^2} \]

we express \( F_{32}(s) \) as

\[ F_{32}(s) = -\frac{2}{5} \left( \frac{s + 2/2 + 2 - 3/2}{(s+2)^2 + 2^2} \right) = -\frac{2}{5} \left( \frac{s + 2}{(s+2)^2 + 2^2} + \frac{-3/2}{(s+2)^2 + 2^2} \right) \]

\[ = -\frac{2}{5} \left( \frac{s + 2}{(s+2)^2 + 2^2} \right) + \frac{6/10}{2} \left( \frac{2}{(s+2)^2 + 2^2} \right) \]

\[ = -\frac{2}{5} \left( \frac{s + 2}{(s+2)^2 + 2^2} \right) + \frac{3}{10} \left( \frac{2}{(s+2)^2 + 2^2} \right) \]  
(5.33)

* Here, we used MATLAB with simple((-1/5 +3j/20)/(s+2+2i))+((-1/5 -3j/20)/(s+2-2i)). The simple function, after several simplification tools that were displayed on the screen, returned \((-2*s-1)/(5*s^2+20*s+40)\)
Addition of (5.30) with (5.33) yields

\[
F_3(s) = F_{31}(s) + F_{32}(s) = \frac{2/5}{s+1} - \frac{2}{5} \left( \frac{s+2}{(s+2)^2 + 2^2} \right) + \frac{3}{10} \left( \frac{2}{(s+2)^2 + 2^2} \right)
\]

\[
\implies \frac{2}{5} e^{-t} - \frac{2}{5} e^{-2t} \cos 2t + \frac{3}{10} e^{-2t} \sin 2t = f_3(t)
\]

Check with MATLAB:

```matlab
syms a s t w; % Define several symbolic variables
Fs=(s + 3)/(s^3 + 5*s^2 + 12*s + 8); ft=ilaplace(Fs)
ft =
2/5*exp(-t) - 2/5*exp(-2*t) *cos(2*t) + 3/10*exp(-2*t) *sin(2*t)
```

**Case III: Multiple (Repeated) Poles**

In this case, \( F(s) \) has simple poles, but one of the poles, say \( p_1 \), has a multiplicity \( m \). For this condition, we express it as

\[
F(s) = \frac{N(s)}{(s-p_1)^m(s-p_2)\ldots(s-p_{n-1})(s-p_n)}
\]  

(5.34)

Denoting the \( m \) residues corresponding to multiple pole \( p_1 \) as \( r_{11}, r_{12}, r_{13}, \ldots, r_{1m} \), the partial fraction expansion of (5.34) is written as

\[
F(s) = \frac{r_{11}}{(s-p_1)^m} + \frac{r_{12}}{(s-p_1)^{m-1}} + \frac{r_{13}}{(s-p_1)^{m-2}} + \ldots + \frac{r_{1m}}{(s-p_1)}
\]

(5.35)

For the simple poles \( p_1, p_2, \ldots, p_n \), we proceed as before, that is, we find the residues as

\[
r_k = \lim_{s \to p_k} (s-p_k)F(s) = (s-p_k)F(s) \bigg|_{s=p_k}
\]

(5.36)

The residues \( r_{11}, r_{12}, r_{13}, \ldots, r_{1m} \) corresponding to the repeated poles, are found by multiplication of both sides of (5.35) by \((s-p)^m\). Then,

\[
(s-p_1)^m F(s) = r_{11} + (s-p_1)r_{12} + (s-p_1)^2 r_{13} + \ldots + (s-p_1)^{m-1} r_{1m}
\]

+ \( (s-p_1)^m \left( \frac{r_2}{(s-p_2)} + \frac{r_3}{(s-p_3)} + \ldots + \frac{r_n}{(s-p_n)} \right) \)

(5.37)
Next, taking the limit as \( s \to p_1 \) on both sides of (5.37), we get

\[
\lim_{s \to p_1} (s - p_1)^m F(s) = r_{11} + \lim_{s \to p_1} \left[(s - p_1)r_{12} + (s - p_1)^2 r_{13} + \ldots + (s - p_1)^{m-1} r_{1m}\right]
\]

\[
+ \lim_{s \to p_1} \left[(s - p_1)^m \left(\frac{r_2}{s - p_2} + \frac{r_3}{s - p_3} + \ldots + \frac{r_n}{s - p_n}\right)\right]
\]

or

\[
r_{11} = \lim_{s \to p_1} (s - p_1)^m F(s) \tag{5.38}
\]

and thus (5.38) yields the residue of the first repeated pole.

The residue \( r_{12} \) for the second repeated pole \( p_1 \), is found by differentiating (5.37) with respect to \( s \) and again, we let \( s \to p_1 \), that is,

\[
r_{12} = \lim_{s \to p_1} \frac{d}{ds}[(s - p_1)^m F(s)] \tag{5.39}
\]

In general, the residue \( r_{1k} \) can be found from

\[
(s - p_1)^m F(s) = r_{11} + r_{12}(s - p_1) + r_{13}(s - p_1)^2 + \ldots \tag{5.40}
\]

whose \((m - 1)th\) derivative of both sides is

\[
(k - 1)! r_{1k} = \lim_{s \to p_1} \frac{1}{(k - 1)!} \frac{d^{k-1}}{ds^{k-1}} [(s - p_1)^m F(s)] \tag{5.41}
\]

or

\[
r_{1k} = \lim_{s \to p_1} \frac{1}{(k - 1)!} \frac{d^{k-1}}{ds^{k-1}} [(s - p_1)^m F(s)] \tag{5.42}
\]

**Example 5.4**

Use the partial fraction expansion method to simplify \( F_4(s) \) of (5.43) below, and find the time domain function \( f_4(t) \) corresponding to \( F_4(s) \).

\[
F_4(s) = \frac{s + 3}{(s + 2)(s + 1)^2} \tag{5.43}
\]

**Solution:**

We observe that there is a pole of multiplicity 2 at \( s = -1 \), and thus in partial fraction expansion form, \( F_4(s) \) is written as
The residues are

\[ r_1 = \frac{s + 3}{(s + 1)^2} \bigg|_{s = -2} = 1 \]

\[ r_{21} = \frac{s + 3}{s + 2} \bigg|_{s = -1} = 2 \]

\[ r_{22} = \frac{d}{ds} \left( \frac{s + 3}{s + 2} \right) \bigg|_{s = -1} = \frac{(s + 2) - (s + 3)}{(s + 2)^2} \bigg|_{s = -1} = -1 \]

The value of the residue \( r_{22} \) can also be found without differentiation as follows:

Substitution of the already known values of \( r_1 \) and \( r_{21} \) into (5.44), and letting \( s = 0^* \), we get

\[ \frac{s + 3}{(s + 1)^2(s + 2)} \bigg|_{s = 0} = \frac{1}{(s + 2)} \bigg|_{s = 0} + \frac{2}{(s + 1)^2} \bigg|_{s = 0} + \frac{r_{22}}{(s + 1)} \bigg|_{s = 0} \]

or

\[ \frac{3}{2} = \frac{1}{2} + 2 + r_{22} \]

from which \( r_{22} = -1 \) as before. Finally,

\[ F_4(s) = \frac{s + 3}{(s + 2)(s + 1)^2} = \frac{1}{(s + 2)} + \frac{2}{(s + 1)^2} + \frac{-1}{(s + 1)} \Leftrightarrow e^{-2t} + 2te^{-t} - e^{-t} = f_4(t) \]  

(5.45)

Check with MATLAB:

```matlab
syms s t; Fs=(s+3)/((s+2)^2*(s+1)^2); ft=ilaplace(Fs)
ft = exp(-2*t) + 2*t*exp(-t) - exp(-t)
```

We can use the following code to check the partial fraction expansion.

```matlab
syms s
Ns = [1 3]; % Coefficients of the numerator N(s) of F(s)
expand((s+1)^2); % Expands (s+1)^2 to s^2 + 2*s + 1;
d1 = [1 2 1]; % Coefficients of (s + 1)^2 = s^2 + 2*s + 1 term in D(s)
d2 = [0 1 2]; % Coefficients of (s + 2) term in D(s)
```

* This is permissible since (5.44) is an identity.

5-10

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Partial Fraction Expansion

Partial Fraction Expansion

\[ D_s = \text{conv}(d1,d2); \quad \text{% Multiplies polynomials } d1 \text{ and } d2 \text{ to express the} \]
\[ \text{denominator } D(s) \text{ of } F(s) \text{ as a polynomial} \]
\[ [r,p,k] = \text{residue}(N_s,D_s) \]
\[ r = \]
\[ 1.0000 \]
\[ -1.0000 \]
\[ 2.0000 \]
\[ p = \]
\[ -2.0000 \]
\[ -1.0000 \]
\[ -1.0000 \]
\[ k = \]
\[ [] \]

**Example 5.5**

Use the partial fraction expansion method to simplify \( F_5(s) \) of (5.46) below, and find the time domain function \( f_5(t) \) corresponding to the given \( F_5(s) \).

\[ F_5(s) = \frac{s^2 + 3s + 1}{(s + 1)^3 (s + 2)^2} \quad (5.46) \]

**Solution:**

We observe that there is a pole of multiplicity 3 at \( s = -1 \), and a pole of multiplicity 2 at \( s = -2 \). Then, in partial fraction expansion form, \( F_5(s) \) is written as

\[ F_5(s) = \frac{r_{11}}{(s + 1)^3} + \frac{r_{12}}{(s + 1)^2} + \frac{r_{13}}{(s + 1)} + \frac{r_{21}}{(s + 2)^2} + \frac{r_{22}}{(s + 2)} \quad (5.47) \]

The residues are

\[ r_{11} = \left. \frac{s^2 + 3s + 1}{(s + 2)^2} \right|_{s = -1} = 1 \]

\[ r_{12} = \left. \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right|_{s = -1} = 3 \]

\[ r_{13} = \left. \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right) \right|_{s = -1} = 3 \]

\[ r_{21} = \left. \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right) \right) \right|_{s = -1} = 3 \]

\[ r_{22} = \left. \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right) \right) \right) \right|_{s = -1} = 3 \]
\[
\begin{align*}
  r_{13} &= \frac{1}{2!} \left. \frac{d^2}{ds^2} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right|_{s = -1} = \frac{1}{2!} \left. \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 2)^2} \right) \right|_{s = -1} \\
  &= \frac{1}{2} \left. \frac{d}{ds} \left( \frac{s + 4}{(s + 2)^3} \right) \right|_{s = -1} = \frac{1}{2} \left. \frac{(s + 2)^3 - 3(s + 2)^2(s + 4)}{(s + 2)^6} \right|_{s = -1} \\
  &= \frac{1}{2} \left. \frac{(s + 2 - 3s - 12}{(s + 2)^4} \right|_{s = -1} = \frac{-s - 5}{(s + 2)^4} \bigg|_{s = -1} = -4
\end{align*}
\]

Next, for the pole at \( s = -2 \),
\[
r_{21} = \left. \frac{s^2 + 3s + 1}{(s + 1)^3} \right|_{s = -2} = 1
\]
and
\[
r_{22} = \left. \frac{d}{ds} \left( \frac{s^2 + 3s + 1}{(s + 1)^3} \right) \right|_{s = -2} = \left. \frac{(s + 1)^3(2s + 3) - 3(s + 1)^2(s^2 + 3s + 1)}{(s + 1)^6} \right|_{s = -2}
\]
\[
= \left. \frac{(s + 1)(2s + 3) - 3(s^2 + 3s + 1)}{(s + 1)^4} \right|_{s = -2} = \frac{-s^2 - 4s}{(s + 1)^4} \bigg|_{s = -2} = 4
\]

By substitution of the residues into (5.47), we get
\[
F(s) = \frac{-1}{(s + 1)^3} + \frac{3}{(s + 1)^2} + \frac{-4}{(s + 1)} + \frac{1}{(s + 2)^2} + \frac{4}{(s + 2)} \quad (5.48)
\]

We will check the values of these residues with the MATLAB code below.

\texttt{syms s; \quad \% The function \texttt{collect(s)} below multiplies \( (s+1)^3 \) by \( (s+2)^2 \)}
\texttt{\% and we use it to express the denominator \( D(s) \) as a polynomial so that we can}
\texttt{\% we can use the coefficients of the resulting polynomial with the \texttt{residue} function}
\texttt{Ds=collect(((s+1)^3)*(s+2)^2))}
\texttt{Ds =}
\texttt{s^5+7*s^4+19*s^3+25*s^2+16*s+4}
\texttt{Ns=[1 3 1]; Ds=[1 7 19 25 16 4]; [r,p,k]=residue(Ns,Ds)}
\texttt{r =}
\texttt{4.0000}
\texttt{1.0000}
\texttt{-4.0000}
\texttt{3.0000}
\texttt{-1.0000}
\( p = \\
-2.0000 \\
-2.0000 \\
-1.0000 \\
-1.0000 \\
-1.0000 \\
\)

\( k = \) 

\[
\begin{array}{c}
\text{From Table 2.2 of Chapter 2} \\
\end{array}
\]

\[
e^{-at} \iff \frac{1}{s+a} \quad te^{-at} \iff \frac{1}{(s+a)^2} \quad t^{n-1}e^{-at} \iff \frac{(n-1)!}{(s+a)^n}
\]

and with these, we derive \( f_3(t) \) from (5.48) as

\[
f_3(t) = \frac{1}{2} t^2 e^{-t} + 3 t e^{-t} - 4 e^{-t} + t e^{-2t} + 4 e^{-2t}
\]  

(5.49)

We can verify (5.49) with MATLAB as follows:

```matlab
syms s t; Fs=-1/((s+1)^3) + 3/((s+1)^2) - 4/(s+1) + 1/((s+2)^2) + 4/(s+2);
ft=ilaplace(Fs);
ft = -1/2*t^2*exp(-t)+3*t*exp(-t)-4*exp(-t) + t*exp(-2*t)+4*exp(-2*t)
```

### 5.3 Case for \( m \geq n \)

Our discussion thus far, was based on the condition that \( F(s) \) is a proper rational function. However, if \( F(s) \) is an improper rational function, that is, if \( m \geq n \), we must first divide the numerator \( N(s) \) by the denominator \( D(s) \) to obtain an expression of the form

\[
F(s) = k_0 + k_1s + k_2s^2 + \ldots + k_{m-n}s^{m-n} + \frac{N(s)}{D(s)}
\]  

(5.50)

where \( N(s)/D(s) \) is a proper rational function.

**Example 5.6**

Derive the Inverse Laplace transform \( f_6(t) \) of

\[
F_6(s) = \frac{s^2 + 2s + 2}{s + 1}
\]  

(5.51)
Solution:

For this example, $F_6(s)$ is an improper rational function. Therefore, we must express it in the form of (5.50) before we use the partial fraction expansion method.

By long division, we get

$$F_6(s) = \frac{s^2 + 2s + 2}{s + 1} = \frac{1}{s + 1} + 1 + s$$

Now, we recognize that

$$\frac{1}{s + 1} \leftrightarrow e^{-t}$$

and

$$1 \leftrightarrow \delta(t)$$

but

$$s \leftrightarrow ?$$

To answer that question, we recall that

$$u_0'(t) = \delta(t)$$

and

$$u_0''(t) = \delta'(t)$$

where $\delta'(t)$ is the doublet of the delta function. Also, by the time differentiation property

$$u_0''(t) = \delta'(t) \leftrightarrow s^2F(s) - sf(0) - f'(0) = s^2F(s) = s^2 \cdot \frac{1}{s} = s$$

Therefore, we have the new transform pair

$$s \leftrightarrow \delta'(t) \quad (5.52)$$

and thus,

$$F_6(s) = \frac{s^2 + 2s + 2}{s + 1} = \frac{1}{s + 1} + 1 + s \leftrightarrow e^{-t} + \delta(t) + \delta'(t) = f_6(t) \quad (5.53)$$

In general,

$$\frac{d^n}{dt^n} \delta(t) \leftrightarrow s^n \quad (5.54)$$

We verify (5.53) with MATLAB as follows:

Ns = [1 2 2]; Ds = [1 1]; [r, p, k] = residue(Ns, Ds)

r =

1
Alternate Method of Partial Fraction Expansion

\[ p = -1 \]
\[ k = \begin{bmatrix} 1 & 1 \end{bmatrix} \]

Here, the direct terms \( k = \begin{bmatrix} 1 & 1 \end{bmatrix} \) are the coefficients of \( \delta(t) \) and \( \delta'(t) \) respectively.

### 5.4 Alternate Method of Partial Fraction Expansion

Partial fraction expansion can also be performed with the *method of clearing the fractions*, that is, making the denominators of both sides the same, then equating the numerators. As before, we assume that \( F(s) \) is a proper rational function. If not, we first perform a long division, and then work with the quotient and the remainder as we did in Example 5.6. We also assume that the denominator \( D(s) \) can be expressed as a product of real linear and quadratic factors. If these assumptions prevail, we let \( (s - a) \) be a linear factor of \( D(s) \), and we assume that \( (s - a)^m \) is the highest power of \( (s - a) \) that divides \( D(s) \). Then, we can express \( F(s) \) as

\[
F(s) = \frac{N(s)}{D(s)} = \frac{r_1}{s - a} + \frac{r_2}{(s - a)^2} + \cdots + \frac{r_m}{(s - a)^m} \tag{5.55}
\]

Let \( s^2 + \alpha s + \beta \) be a quadratic factor of \( D(s) \), and suppose that \( (s^2 + \alpha s + \beta)^n \) is the highest power of this factor that divides \( D(s) \). Now, we perform the following steps:

1. To this factor, we assign the sum of \( n \) partial fractions, that is,

\[
\frac{r_1s + k_1}{s^2 + \alpha s + \beta} + \frac{r_2s + k_2}{(s^2 + \alpha s + \beta)^2} + \cdots + \frac{r_n s + k_n}{(s^2 + \alpha s + \beta)^n}
\]

2. We repeat step 1 for each of the distinct linear and quadratic factors of \( D(s) \)

3. We set the given \( F(s) \) equal to the sum of these partial fractions

4. We clear the resulting expression of fractions and arrange the terms in decreasing powers of \( s \)

5. We equate the coefficients of corresponding powers of \( s \)

6. We solve the resulting equations for the residues

**Example 5.7**

Express \( F_7(s) \) of (5.56) below as a sum of partial fractions using the method of clearing the fractions.

\[
F_7(s) = \frac{-2s + 4}{(s^2 + 1)(s - 1)^2} \tag{5.56}
\]
Solution:

Using Steps 1 through 3 above, we get

\[ F_7(s) = \frac{-2s + 4}{(s^2 + 1)(s - 1)^2} = \frac{r_1s + A}{(s^2 + 1)} + \frac{r_{21}}{(s - 1)^2} + \frac{r_{22}}{(s - 1)} \quad (5.57) \]

With Step 4,

\[ -2s + 4 = (r_1s + A)(s - 1)^2 + r_{21}(s^2 + 1) + r_{22}(s - 1)(s^2 + 1) \quad (5.58) \]

and with Step 5,

\[ -2s + 4 = (r_1 + r_{22})s^3 + (-2r_1 + A + r_{22} - r_{21})s^2 + (r_1 - 2A + r_{22})s + (A - r_{22} + r_{21}) \quad (5.59) \]

Relation (5.59) will be an identity is \( s \) if each power of \( s \) is the same on both sides of this relation. Therefore, we equate like powers of \( s \) and we get

\[
\begin{align*}
0 &= r_1 + r_{22} \\
0 &= -2r_1 + A - r_{22} + r_{21} \\
-2 &= r_1 - 2A + r_{22} \\
4 &= A - r_{22} + r_{21} 
\end{align*} \quad (5.60)
\]

Subtracting the second equation of (5.60) from the fourth, we get

\[ 4 = 2r_1 \]

or

\[ r_1 = 2 \quad (5.61) \]

By substitution of (5.61) into the first equation of (5.60), we get

\[ 0 = 2 + r_{22} \]

or

\[ r_{22} = -2 \quad (5.62) \]

Next, substitution of (5.61) and (5.62) into the third equation of (5.60) yields

\[ -2 = 2 - 2A - 2 \]

or

\[ A = 1 \quad (5.63) \]

Finally by substitution of (5.61), (5.62), and (5.63) into the fourth equation of (5.60), we get
Alternate Method of Partial Fraction Expansion

\[ 4 = 1 + 2 + r_{21} \]
or
\[ r_{21} = 1 \] (5.64)

Substitution of these values into (5.57) yields
\[ F_7(s) = \frac{-2s + 4}{(s^2 + 1)(s - 1)^2} = \frac{2s + 1}{(s^2 + 1)} + \frac{1}{(s - 1)^2} - \frac{2}{(s - 1)} \] (5.65)

**Example 5.8**

Use partial fraction expansion to simplify \( F_8(s) \) of (5.66) below, and find the time domain function \( f_8(t) \) corresponding to \( F_8(s) \).

\[ F_8(s) = \frac{s + 3}{s^3 + 5s^2 + 12s + 8} \] (5.66)

**Solution:**

This is the same transform as in Example 5.3, where we found that the denominator \( D(s) \) can be expressed in factored form of a linear term and a quadratic. Thus, we write \( F_8(s) \) as

\[ F_8(s) = \frac{s + 3}{(s + 1)(s^2 + 4s + 8)} \] (5.67)

and using the method of clearing the fractions, we rewrite (5.67) as

\[ F_8(s) = \frac{s + 3}{(s + 1)(s^2 + 4s + 8)} = \frac{r_1}{s + 1} + \frac{r_2s + r_3}{s^2 + 4s + 8} \] (5.68)

As in Example 5.3,

\[ r_1 = \left. \frac{s + 3}{s^2 + 4s + 8} \right|_{s = -1} = \frac{2}{5} \] (5.69)

Next, to compute \( r_2 \) and \( r_3 \), we follow the procedure of this section and we get

\[ (s + 3) = r_1(s^2 + 4s + 8) + (r_2s + r_3)(s + 1) \] (5.70)

Since \( r_1 \) is already known, we only need two equations in \( r_2 \) and \( r_3 \). Equating the coefficient of \( s^2 \) on the left side, which is zero, with the coefficients of \( s^2 \) on the right side of (5.70), we get

\[ 0 = r_1 + r_2 \] (5.71)
and since \( r_1 = 2/5 \), then \( r_2 = -2/5 \).

To obtain the third residue \( r_3 \), we equate the constant terms of (5.70). Then, \( 3 = 8r_1 + r_3 \) or \( 3 = 8 \times 2/5 + r_3 \), or \( r_3 = -1/5 \). Then, by substitution into (5.68), we get

\[
F_8(s) = \frac{2/5}{s + 1} \frac{1}{5} \frac{(2s + 1)}{(s^2 + 4s + 8)}
\]

as before.

The remaining steps are the same as in Example 5.3, and thus \( f_8(t) \) is the same as \( f_3(t) \), that is,

\[
f_8(t) = f_3(t) = \left( \frac{2}{5} e^{-t} - \frac{2}{5} e^{-2t} \cos 2t + \frac{3}{10} e^{-2t} \sin 2t \right) u_0(t)
\]

### 5.5 Summary

- The Inverse Laplace Transform Integral defined as

\[
\mathcal{L}^{-1}\{F(s)\} = f(t) = \frac{1}{2\pi j} \int_{\sigma-j\infty}^{\sigma+j\infty} F(s)e^{st} ds
\]

is difficult to evaluate because it requires contour integration using complex variables theory.

- For most engineering problems we can refer to Tables of Properties, and Common Laplace transform pairs to lookup the Inverse Laplace transform.

- The partial fraction expansion method offers a convenient means of expressing Laplace transforms in a recognizable form from which we can obtain the equivalent time-domain functions.

- If the highest power \( m \) of the numerator \( N(s) \) is less than the highest power \( n \) of the denominator \( D(s) \), i.e., \( m < n \), \( F(s) \) is said to be expressed as a proper rational function. If \( m \geq n \), \( F(s) \) is an improper rational function.

- The Laplace transform \( F(s) \) must be expressed as a proper rational function before applying the partial fraction expansion. If \( F(s) \) is an improper rational function, that is, if \( m \geq n \), we must first divide the numerator \( N(s) \) by the denominator \( D(s) \) to obtain an expression of the form

\[
F(s) = k_0 + k_1 s + k_2 s^2 + \ldots + k_{m-n} s^{m-n} + \frac{N(s)}{D(s)}
\]

- In a proper rational function, the roots of numerator \( N(s) \) are called the zeros of \( F(s) \) and the roots of the denominator \( D(s) \) are called the poles of \( F(s) \).
The partial fraction expansion method can be applied whether the poles of $F(s)$ are distinct, complex conjugates, repeated, or a combination of these.

When $F(s)$ is expressed as

$$ F(s) = \frac{r_1}{s-p_1} + \frac{r_2}{s-p_2} + \frac{r_3}{s-p_3} + \ldots + \frac{r_n}{s-p_n} $$

$r_1, r_2, r_3, \ldots, r_n$ are called the residues and $p_1, p_2, p_3, \ldots, p_n$ are the poles of $F(s)$.

The residues and poles of a rational function of polynomials can be found easily using the MATLAB `residue(a,b)` function. The direct term is always empty (has no value) whenever $F(s)$ is a proper rational function.

We can use the MATLAB `factor(s)` symbolic function to convert the denominator polynomial form of $F_2(s)$ into a factored form.

We can use the MATLAB `collect(s)` and `expand(s)` symbolic functions to convert the denominator factored form of $F_2(s)$ into a polynomial form.

In this chapter we developed the new transform pair

$$ s \Leftrightarrow \delta'(t) $$

and in general,

$$ \frac{d^n}{dt^n} \delta(t) \Leftrightarrow s^n $$

The method of clearing the fractions is an alternate method of partial fraction expansion.
5.6 Exercises

1. Find the Inverse Laplace transform of the following:
   
   a. \( \frac{4}{s + 3} \)
   
   b. \( \frac{4}{(s + 3)^2} \)
   
   c. \( \frac{4}{(s + 3)^4} \)
   
   d. \( \frac{3s + 4}{(s + 3)^5} \)
   
   e. \( \frac{s^2 + 6s + 3}{(s + 3)^5} \)

2. Find the Inverse Laplace transform of the following:
   
   a. \( \frac{3s + 4}{s^2 + 4s + 85} \)
   
   b. \( \frac{4s + 5}{s^2 + 5s + 18.5} \)
   
   c. \( \frac{s^2 + 3s + 2}{s^3 + 5s^2 + 10.5s + 9} \)
   
   d. \( \frac{s^2 - 16}{s^3 + 8s^2 + 24s + 32} \)
   
   e. \( \frac{s + 1}{s^3 + 6s^2 + 11s + 6} \)

3. Find the Inverse Laplace transform of the following:
   
   a. \( \frac{3s + 2}{s^2 + 25} \)
   
   b. \( \frac{5s^2 + 3}{(s^2 + 4)^2} \) (See hint on next page)
Hint:
\[
\begin{align*}
\frac{1}{2\alpha}(\sin \alpha t + \alpha \cos \alpha t) & \leftrightarrow \frac{s^2}{(s^2 + \alpha^2)^2} \\
\frac{1}{2\alpha^3}(\sin \alpha t - \alpha \cos \alpha t) & \leftrightarrow \frac{1}{(s^2 + \alpha^2)^2}
\end{align*}
\]

c. \[ \frac{2s + 3}{s^2 + 4.25s + 1} \]

d. \[ \frac{s^3 + 8s^2 + 24s + 32}{s^2 + 6s + 8} \]

e. \[ e^{-2s} \frac{3}{(2s + 3)^3} \]

4. Use the Initial Value Theorem to find \( f(0) \) given that the Laplace transform of \( f(t) \) is
\[ \frac{2s + 3}{s^2 + 4.25s + 1} \]

Compare your answer with that of Exercise 3(c).

5. It is known that the Laplace transform \( F(s) \) has two distinct poles, one at \( s = 0 \), the other at \( s = -1 \). It also has a single zero at \( s = 1 \), and we know that \( \lim_{t \to \infty} f(t) = 10 \). Find \( F(s) \) and \( f(t) \).
5.7 Solutions to Exercises

1. 
   a. \[ \frac{4}{s+3} \Leftrightarrow 4e^{-3t} \]
   b. \[ \frac{4}{(s+3)^2} \Leftrightarrow 4te^{-3t} \]
   c. \[ \frac{4}{3!} \Leftrightarrow \frac{4}{3!} t^3 e^{-3t} = \frac{2}{3} t^3 e^{-3t} \]
   d. \[ \frac{3s + 4}{(s + 3)^5} = \frac{3(s + 4/3 + 5/3 - 5/3)}{(s + 3)^5} = 3 \cdot \left( \frac{s + 3}{(s + 3)^5} - \frac{5/3}{(s + 3)^5} \right) = 3 \cdot \frac{1}{(s + 3)^4} - \frac{5}{12} \cdot \frac{1}{(s + 3)^5} \]
   \[ \Leftrightarrow \frac{3}{3!} t^3 e^{-3t} - \frac{5}{4!} t^4 e^{-3t} = \frac{1}{2} \left( t^3 e^{-3t} - \frac{5}{12} t^4 e^{-3t} \right) \]
   e. \[ \frac{s^2 + 6s + 3}{(s + 3)^5} = \frac{s^2 + 6s + 9 - 6}{(s + 3)^5} = \frac{(s + 3)^2}{(s + 3)^5} - \frac{6}{(s + 3)^5} = \frac{1}{(s + 3)^3} - \frac{6}{(s + 3)^5} \]
   \[ \Leftrightarrow \frac{1}{2!} t^2 e^{-3t} - \frac{6}{4!} t^4 e^{-3t} = \frac{1}{2} \left( t^2 e^{-3t} - \frac{1}{2} t^4 e^{-3t} \right) \]

2. 
   a. \[ \frac{3s + 4}{s^2 + 4s + 85} = \frac{3(s + 3/3 + 2/3 - 2/3)}{(s + 2)^2 + 9} = 3 \cdot \left( \frac{s + 2}{(s + 2)^2 + 9} - \frac{1}{9} \cdot \frac{2 \times 9}{(s + 2)^2 + 9} \right) \]
   \[ = 3 \cdot \frac{(s + 2)}{(s + 2)^2 + 9} - \frac{2}{9} \cdot \frac{9}{(s + 2)^2 + 9} \Leftrightarrow 3e^{-2t} \cos 9t = \frac{2}{9} e^{-2t} \sin 9t \]
   b. \[ \frac{4s + 5}{s^2 + 5s + 18.5} = \frac{4s + 5}{s^2 + 5s + 6.25 + 12.25} = \frac{4s + 5}{(s + 2.5)^2 + 3.5^2} = 4 \cdot \frac{s + 5/4}{(s + 2.5)^2 + 3.5^2} \]
   \[ = 4 \cdot \frac{s + 10/4 - 10/4 + 5/4}{(s + 2.5)^2 + 3.5^2} = 4 \cdot \frac{s + 2.5}{(s + 2.5)^2 + 3.5^2} - \frac{1}{3.5} \cdot \frac{5 \times 3.5}{(s + 2.5)^2 + 3.5^2} \]
   \[ = 4 \cdot \frac{(s + 2.5)}{(s + 2.5)^2 + 3.5^2} - \frac{10}{7} \cdot \frac{3.5}{(s + 2.5)^2 + 3.5^2} \Leftrightarrow 4e^{-2.5t} \cos 3.5t - \frac{10}{7} e^{-2.5t} \sin 3.5t \]
   c. Using the MATLAB \texttt{factor(s)} function we get:
   
   \texttt{syms s; factor(s^2+3*s+2), factor(s^3+5*s^2+2+10.5*s+9)}
   
   \texttt{ans = (s+2)*(s+1)}
   
   \texttt{ans = 1/2*(s+2)*(2*s^2+6*s+9)}
Then,

\[
\frac{s^2 + 3s + 2}{s^3 + 5s^2 + 10.5s + 9} = \frac{(s + 1)(s + 2)}{(s + 2)(s^2 + 3s + 4.5)} = \frac{(s + 1)}{(s^2 + 3s + 4.5)} = \frac{s + 1}{s^2 + 3s + 2.25 - 2.25 + 4.5} = \frac{s + 1}{s^2 + 3s + 2.25 - 2.25 + 4.5} = \frac{s + 1.5 - 1.5 + 1}{s^2 + 1.5^2 + (1.5)^2} \cdot \frac{1}{0.5 \times 1.5} \cdot \frac{1.5}{s + 2.25 + (1.5)^2} \Rightarrow e^{-1.5t} \cos 1.5t - \frac{1}{3} e^{-1.5t} \sin 1.5t
\]

d.

\[
\frac{s^2 - 16}{s^3 + 8s^2 + 24s + 32} = \frac{(s + 4)(s - 4)}{(s + 4)(s^2 + 4s + 8)} = \frac{(s - 4)}{(s^2 + 4s + 8)} = \frac{s + 2 - 2 - 4}{(s + 2)^2 + 2^2} = \frac{s + 2}{(s + 2)^2 + 2^2} - \frac{1}{2} \cdot \frac{6 \times 2}{(s + 2)^2 + 2^2} = \frac{s + 2}{(s + 2)^2 + 2^2} - \frac{3 \cdot 2}{(s + 2)^2 + 2^2} \Rightarrow e^{-2t} \cos 2t - 3e^{-2t} \sin 2t
\]

e.

\[
\frac{s + 1}{s^3 + 6s^2 + 11s + 6} = \frac{(s + 1)}{(s + 1)(s + 2)(s + 3)} = \frac{1}{(s + 1)(s + 2)(s + 3)} = \frac{1}{(s + 2)(s + 3)} = \left[ \frac{1}{s + 2} - \frac{1}{s + 3} \right] \Rightarrow e^{-2t} - e^{-3t}
\]

3.

a.

\[
\frac{3s + 2}{s^2 + 25} = \frac{3s}{s^2 + 5^2} + \frac{2 \times 5}{s^2 + 5^2} = 3 \cdot \frac{s}{s^2 + 5^2} + \frac{2 \cdot 5}{5} \cdot \frac{5}{s^2 + 5^2} \Rightarrow 3 \cos 5t + \frac{2}{5} \sin 5t
\]

b.

\[
\frac{5s^2 + 3}{(s^2 + 4)^2} = \frac{5s^2}{(s^2 + 2^2)^2} + \frac{3}{(s^2 + 2^2)^2} \Leftrightarrow 5 \cdot \frac{1}{2 \times 2} (\sin 2t + 2t \cos 2t) + 3 \cdot \frac{1}{2 \times 8} (\sin 2t - 2t \cos 2t) \Leftrightarrow \left( \frac{5}{4} + \frac{3}{16} \right) \sin 2t + \left( \frac{5}{4} - \frac{3}{16} \right)2t \cos 2t = \frac{23}{16} \sin 2t + \frac{17}{8} t \cos 2t
\]
c. 
\[
\frac{2s + 3}{s^2 + 4.25s + 1} = \frac{2s + 3}{(s + 4)(s + 1/4)} = \frac{r_1}{s + 4} + \frac{r_2}{s + 1/4}
\]
\[
r_1 = \frac{2s + 3}{s + 1/4} \bigg|_{s = -4} = -\frac{5}{15/4} = \frac{4}{3} \quad r_2 = \frac{2s + 3}{s + 4} \bigg|_{s = -1/4} = \frac{5/2}{15/4} = \frac{2}{3}
\]
\[
\frac{4/3}{s + 4} + \frac{2/3}{s + 1/4} \Leftrightarrow \frac{2}{3}(2e^{-4t} + e^{-1/4})
\]

4. The initial value theorem states that \( \lim_{t \to 0} f(t) = \lim_{s \to \infty} sF(s). \) Then,

\[
f(0) = \lim_{s \to \infty} \frac{s^2 + 3}{s^2 + 4.25s + 1} = \lim_{s \to \infty} \frac{s^2 + 3s}{s^2 + 4.25s + 1} = \lim_{s \to \infty} 2 + \frac{3}{s} = 2
\]
The value \( f(0) = 2 \) is the same as in the time domain expression that we found in Exercise 3(c).

5. We are given that \( F(s) = \frac{A(s - 1)}{s(s + 1)} \) and \( \lim_{t \to \infty} f(t) = \lim_{s \to 0} sF(s) = 10. \) Then,

\[
\lim_{s \to 0} \frac{A(s - 1)}{s(s + 1)} = A \lim_{s \to 0} \frac{(s - 1)}{s(s + 1)} = -A = 10. \] Therefore,

\[
F(s) = \frac{-10(s - 1)}{s(s + 1)} = \frac{r_1}{s} + \frac{r_2}{s + 1} = \frac{10}{s} - \frac{20}{s + 1} \Leftrightarrow (10 - 20e^{-t})u_0(t), \text{ that is, }
\]
\[
f(t) = (10 - 20e^{-t})u_0(t) \text{ and we see that } \lim_{t \to \infty} f(t) = 10.\]
Chapter 6
Circuit Analysis with Laplace Transforms

This chapter presents applications of the Laplace transform. Several examples are given to illustrate how the Laplace transformation is applied to circuit analysis. Complex impedance, complex admittance, and transfer functions are also defined.

6.1 Circuit Transformation from Time to Complex Frequency

In this section we will derive the voltage-current relationships for the three elementary circuit devices, i.e., resistors, inductors, and capacitors in the complex frequency domain.

a. Resistor

The time and complex frequency domains for purely resistive circuits are shown in Figure 6.1.

\[ v_R(t) = Ri_R(t) \quad i_R(t) = \frac{v_R(t)}{R} \]

Figure 6.1. Resistive circuit in time domain and complex frequency domain

b. Inductor

The time and complex frequency domains for purely inductive circuits is shown in Figure 6.2.

\[ v_L(t) = L \frac{di_L}{dt} \quad i_L(t) = \frac{1}{L} \int_{-\infty}^{t} v_L dt \]

\[ V_L(s) = sLI_L(s) - Li_L(0^-) \quad I_L(s) = \frac{V_L(s)}{Ls} + \frac{i_L(0^-)}{s} \]

Figure 6.2. Inductive circuit in time domain and complex frequency domain

c. Capacitor

The time and complex frequency domains for purely capacitive circuits is shown in Figure 6.3.
Chapter 6  Circuit Analysis with Laplace Transforms

**Figure 6.3.** Capacitive circuit in time domain and complex frequency domain

**Note:**

In the complex frequency domain, the terms $sL$ and $1/sC$ are called *complex inductive impedance*, and *complex capacitive impedance* respectively. Likewise, the terms $sC$ and $1/sL$ are called *complex capacitive admittance* and *complex inductive admittance* respectively.

**Example 6.1**

Use the Laplace transform method to find the voltage $v_C(t)$ across the capacitor for the circuit of Figure 6.4, given that $v_C(0^-) = 6$ V.

**Solution:**

We apply KCL at node $A$ as shown in Figure 6.5.

Then,

$$i_R + i_C = 0$$

or
The Laplace transform of (6.1) is

\[ sV_C(s) - v_C(0^-) + V_C(s) = \frac{12}{s} \]

or

\[ (s + 1)V_C(s) = \frac{12}{s} + 6 \]

or

\[ V_C(s) = \frac{6s + 12}{s(s + 1)} \]

By partial fraction expansion,

\[ V_C(s) = \frac{6s + 12}{s(s + 1)} = \frac{r_1}{s} + \frac{r_2}{s + 1} \]

\[ r_1 = \frac{6s + 12}{s + 1} \bigg|_{s = 0} = 12 \]

\[ r_2 = \frac{6s + 12}{s} \bigg|_{s = -1} = -6 \]

Therefore,

\[ V_C(s) = \frac{12}{s} - \frac{6}{s + 1} \quad \Leftrightarrow \quad 12 - 6e^{-t} = (12 - 6e^{-t})u_0(t) = v_C(t) \]

**Example 6.2**

Use the Laplace transform method to find the current \( i_C(t) \) through the capacitor for the circuit of Figure 6.6, given that \( v_C(0^-) = 6 \text{ V} \).

![Figure 6.6. Circuit for Example 6.2](image-url)
Solution:

This is the same circuit as in Example 6.1. We apply KVL for the loop shown in Figure 6.7.

\[
\begin{align*}
\text{Figure 6.7. Application of KVL for the circuit of Example 6.2}
\end{align*}
\]

\[
Ri_C(t) + \frac{1}{C} \int_{-\infty}^{t} i_C(t) \, dt = 12u_0(t)
\]

and with \( R = 1 \) and \( C = 1 \), we get

\[
i_C(t) + \int_{-\infty}^{t} i_C(t) \, dt = 12u_0(t)
\]  

Next, taking the Laplace transform of both sides of (6.2), we get

\[
I_C(s) + \frac{I_C(s)}{s} + \frac{v_C(0^-)}{s} = \frac{12}{s}
\]

\[
\left(1 + \frac{1}{s}\right)I_C(s) = \frac{12}{s} - \frac{6}{s} = \frac{6}{s}
\]

\[
\left(\frac{s + 1}{s}\right)I_C(s) = \frac{6}{s}
\]

or

\[
I_C(s) = \frac{6}{s+1} \quad \Rightarrow \quad i_C(t) = 6e^{-t}u_0(t)
\]

Check: From Example 6.1,

\[
v_C(t) = (12 - 6e^{-t})u_0(t)
\]

Then,

\[
i_C(t) = C \frac{dv_C}{dt} = \frac{dv_C}{dt} \left(12 - 6e^{-t}\right)u_0(t) = 6e^{-t}u_0(t) + 6\delta(t)
\]

The presence of the delta function in (6.3) is a result of the unit step that is applied at \( t = 0 \).
Example 6.3

In the circuit of Figure 6.8, switch $S_1$ closes at $t = 0$, while at the same time, switch $S_2$ opens. Use the Laplace transform method to find $v_{out}(t)$ for $t > 0$.

Solution:

Since the circuit contains a capacitor and an inductor, we must consider two initial conditions. One is given as $v_C(0^-) = 3 \, V$. The other initial condition is obtained by observing that there is an initial current of $2 \, A$ in inductor $L_1$; this is provided by the $2 \, A$ current source just before switch $S_2$ opens. Therefore, our second initial condition is $i_{L_1}(0^-) = 2 \, A$.

For $t > 0$, we transform the circuit of Figure 6.8 into its $s$-domain* equivalent shown in Figure 6.9.

In Figure 6.9 the current in $L_1$ has been replaced by a voltage source of $1 \, V$. This is found from the relation

$$L_1i_{L_1}(0^-) = \frac{1}{2} \times 2 = 1 \, V$$  

(6.4)

* Henceforth, for convenience, we will refer the time domain as $t$-domain and the complex frequency domain as $s$-domain
The polarity of this voltage source is as shown in Figure 6.9 so that it is consistent with the direction of the current $i_{L1}(t)$ in the circuit of Figure 6.8 just before switch $S_2$ opens.

The initial capacitor voltage is replaced by a voltage source equal to $3/s$.

Applying KCL at node $\Omega$, we get

$$\frac{V_{out}(s) - 1 - 3/s}{1/s + 2 + s/2} + \frac{V_{out}(s)}{1} + \frac{V_{out}(s)}{s/2} = 0 \quad (6.5)$$

and after simplification

$$V_{out}(s) = \frac{2s(s + 3)}{s^3 + 8s^2 + 10s + 4} \quad (6.6)$$

We will use MATLAB to factor the denominator $D(s)$ of (6.6) into a linear and a quadratic factor.

\[
p=[1 \ 8 \ 10 \ 4]; \quad r=\text{roots}(p) \quad \% \text{Find the roots of } D(s)
\]

\[
r =
\begin{align*}
-6.5708 \\
-0.7146 + 0.3132i \\
-0.7146 - 0.3132i
\end{align*}
\]

\[
y=\text{expand}((s + 0.7146 - 0.3132i)*(s + 0.7146 + 0.3132i)) \quad \% \text{Find quadratic form}
\]

\[
y =
\begin{align*}
s^2 + \frac{3573}{2500}s + \frac{3043737}{5000000}
\end{align*}
\]

\[
\frac{3573}{2500} \quad \% \text{Find coefficient of } s
\]

\[
\text{ans} =
\begin{align*}
1.4292
\end{align*}
\]

\[
\frac{3043737}{5000000} \quad \% \text{Find constant term}
\]

\[
\text{ans} =
\begin{align*}
0.6087
\end{align*}
\]

Therefore,

$$V_{out}(s) = \frac{2s(s + 3)}{s^3 + 8s^2 + 10s + 4} = \frac{2s(s + 3)}{(s + 6.57)(s^2 + 1.43s + 0.61)} \quad (6.7)$$

Now, we perform partial fraction expansion.

$$V_{out}(s) = \frac{2s(s + 3)}{(s + 6.57)(s^2 + 1.43s + 0.61)} = \frac{r_1}{s + 6.57} + \frac{r_2 s + r_3}{s^2 + 1.43s + 0.61} \quad (6.8)$$
Circuit Transformation from Time to Complex Frequency

\[ r_1 = \frac{2s(s + 3)}{s^2 + 1.43s + 0.61} \bigg|_{s = -6.57} = 1.36 \]  
\[ (6.9) \]

The residues \( r_2 \) and \( r_3 \) are found from the equality

\[ 2s(s + 3) = r_1(s^2 + 1.43s + 0.61) + (r_2 s + r_3)(s + 6.57) \]  
\[ (6.10) \]

Equating constant terms of (6.10), we get

\[ 0 = 0.61r_1 + 6.57r_3 \]

and by substitution of the known value of \( r_1 \) from (6.9), we get

\[ r_3 = -0.12 \]

Similarly, equating coefficients of \( s^2 \), we get

\[ 2 = r_1 + r_2 \]

and using the known value of \( r_1 \), we get

\[ r_2 = 0.64 \]  
\[ (6.11) \]

By substitution into (6.8),

\[ V_{out}(s) = \frac{1.36}{s + 6.57} + \frac{0.64s - 0.12}{s^2 + 1.43s + 0.61} = \frac{1.36}{s + 6.57} + \frac{0.64s + 0.46 - 0.58}{s^2 + 1.43s + 0.51 + 0.1} \]

or

\[ V_{out}(s) = \frac{1.36}{s + 6.57} + \frac{0.64(s + 0.715)}{(s + 0.715)^2 + (0.316)^2} - \frac{0.58}{(s + 0.715)^2 + (0.316)^2} \]

\[ = \frac{1.36}{s + 6.57} + \frac{0.64(s + 0.715)}{(s + 0.715)^2 + (0.316)^2} - \frac{1.84 \times 0.316}{(s + 0.715)^2 + (0.316)^2} \]  
\[ (6.12) \]

Taking the Inverse Laplace of (6.12), we get

\[ * \text{ We perform these steps to express the term} \frac{0.64s - 0.12}{s^2 + 1.43s + 0.61} \text{ in a form that resembles the transform pairs} \]

\[ e^{-at} \cos \omega t \leftrightarrow \frac{s + a}{(s + a)^2 + \omega^2} \text{ and } e^{-at} \sin \omega t \leftrightarrow \frac{\omega}{(s + a)^2 + \omega^2}. \text{ The remaining steps are carried out in} \]

\[ (6.12). \]
6.2 Complex Impedance $Z(s)$

Consider the $s$-domain $RLC$ series circuit of Figure 6.10, where the initial conditions are assumed to be zero.

\[ v_{out}(t) = (1.36e^{-0.657t} + 0.64e^{-0.715t} \cos 0.316t - 1.84e^{-0.715t} \sin 0.316t)u_0(t) \]  

(6.13)

For this circuit, the sum $R + sL + \frac{1}{sC}$ represents the total opposition to current flow. Then,

\[ I(s) = \frac{V_s(s)}{R + sL + \frac{1}{sC}} \]  

(6.14)

and defining the ratio $V_s(s)/I(s)$ as $Z(s)$, we get

\[ Z(s) \equiv \frac{V_s(s)}{I(s)} = R + sL + \frac{1}{sC} \]  

(6.15)

and thus, the $s$-domain current $I(s)$ can be found from

\[ I(s) = \frac{V_s(s)}{Z(s)} \]  

(6.16)

where

\[ Z(s) = R + sL + \frac{1}{sC} \]  

(6.17)

We recall that $s = \sigma + j\omega$. Therefore, $Z(s)$ is a complex quantity, and it is referred to as the complex input impedance of an $s$-domain $RLC$ series circuit. In other words, $Z(s)$ is the ratio of the voltage excitation $V_s(s)$ to the current response $I(s)$ under zero state (zero initial conditions).
Example 6.4

Find $Z(s)$ for the circuit of Figure 6.11. All values are in $\Omega$ (ohms).

\[
\begin{align*}
V_S(s) & \quad \quad \quad I(s) \\
\uparrow & \quad \quad \quad A \\
\downarrow & \quad \quad \quad s \\
\downarrow & \quad \quad \quad s \\
\end{align*}
\]

Figure 6.11. Circuit for Example 6.4

Solution:

First Method:

We will first find $I(s)$, and we will compute $Z(s)$ from (6.15). We assign the voltage $V_A(s)$ at node $A$ as shown in Figure 6.12.

\[
\begin{align*}
\frac{V_A(s) - V_S(s)}{I} + \frac{V_A(s)}{s} + \frac{V_A(s)}{s + 1/s} &= 0 \\
\left(1 + \frac{1}{s} + \frac{1}{s+1/s}\right)V_A(s) &= V_S(s) \\
V_A(s) &= \frac{s^3 + 1}{s^3 + 2s^2 + s + 1} \cdot V_S(s)
\end{align*}
\]

The current $I(s)$ is now found as

\[
I(s) = \frac{V_S(s) - V_A(s)}{I} = \left(1 - \frac{s^3 + 1}{s^3 + 2s^2 + s + 1}\right)V_S(s) = \frac{2s^2 + 1}{s^3 + 2s^2 + s + 1} \cdot V_S(s)
\]

and thus,

\[
Z(s) = \frac{V_S(s)}{I(s)} = \frac{s^3 + 2s^2 + s + 1}{2s^2 + 1} \quad (6.18)
\]
Second Method:

We can also compute $Z(s)$ by successive combinations of series and parallel impedances, as it is done with series and parallel resistances. For this example, we denote the network devices as $Z_1, Z_2, Z_3$ and $Z_4$ shown in Figure 6.13.

![Figure 6.13. Computation of the impedance of Example 6.4 by series – parallel combinations](image)

To find the equivalent impedance $Z(s)$, looking to the right of terminals $a$ and $b$, we start on the right side of the network and we proceed to the left combining impedances as we combine resistances. Then,

$$Z(s) = [(Z_3 + Z_4) \ || Z_2] + Z_1$$

$Z(s) = \frac{s(s + 1/s)}{s + s + 1/s} + 1 = \frac{s^2 + 1}{(2s^2 + 1)/s} + 1 = \frac{s^3 + s}{2s^2 + 1} + 1 = \frac{s^3 + 2s^2 + s + 1}{2s^2 + 1} \quad (6.19)$

We observe that (6.19) is the same as (6.18).

6.3 Complex Admittance $Y(s)$

Consider the $s$-domain $GLC$ parallel circuit of Figure 6.14 where the initial conditions are zero.

![Figure 6.14. Parallel GLC circuit in s-domain](image)

For this circuit,

$$GV(s) + \frac{1}{sL}V(s) + sCV(s) = I(s)$$

$$\left(G + \frac{1}{sL} + sC\right)(V(s)) = I(s)$$

Defining the ratio $I_s(s)/V(s)$ as $Y(s)$, we get
Complex Admittance $Y(s)$

\[
Y(s) = \frac{I(s)}{V(s)} = G + \frac{1}{sL} + sC = \frac{1}{Z(s)} \tag{6.20}
\]

and thus the $s$-domain voltage $V(s)$ can be found from

\[
V(s) = \frac{I_s(s)}{Y(s)} \tag{6.21}
\]

where

\[
Y(s) = G + \frac{1}{sL} + sC \tag{6.22}
\]

We recall that $s = \sigma + j\omega$. Therefore, $Y(s)$ is a complex quantity, and it is referred to as the complex input admittance of an $s$-domain GLC parallel circuit. In other words, $Y(s)$ is the ratio of the current excitation $I_s(s)$ to the voltage response $V(s)$ under zero state (zero initial conditions).

**Example 6.5**

Compute $Z(s)$ and $Y(s)$ for the circuit of Figure 6.15. All values are in $\Omega$ (ohms). Verify your answers with MATLAB.

![Figure 6.15. Circuit for Example 6.5](image)

**Solution:**

It is convenient to represent the given circuit as shown in Figure 6.16.
where

\[ Z_1 = 13s + \frac{8}{s} = \frac{13s^2 + 8}{s} \]

\[ Z_2 = 10 + 5s \]

\[ Z_3 = 20 + \frac{16}{s} = \frac{4(5s + 4)}{s} \]

Then,

\[ Z(s) = Z_1 + \frac{Z_2Z_3}{Z_2 + Z_3} = \frac{13s^2 + 8}{s} + \frac{(10 + 5s)\left(\frac{4(5s + 4)}{s}\right)}{10 + 5s + \frac{4(5s + 4)}{s}} \]

\[ = \frac{13s^2 + 8 + (10 + 5s)\left(\frac{4(5s + 4)}{s}\right)}{5s^2 + 10s + 4(5s + 4)} = \frac{13s^2 + 8 + 20(5s^2 + 14s + 8)}{5s^2 + 30s + 16} \]

\[ = \frac{65s^4 + 490s^3 + 528s^2 + 400s + 128}{s(5s^2 + 30s + 16)} \]

Check with MATLAB:

\[
\text{syms } s; \quad z1 = 13*s + \frac{8}{s}; \quad z2 = 5*s + 10; \quad z3 = 20 + \frac{16}{s}; \quad z = z1 + z2 * z3 / (z2+z3) \\
\]

\[ z = 13*s+8/s+(5*s+10)*(20+16/s)/(5*s+30+16/s) \]

\[ z10 = \text{simplify}(z) \]

\[ z10 = (65*s^4+490*s^3+528*s^2+400*s+128)/s/(5*s^2+30*s+16) \]

\[ \text{pretty}(z10) \]

\[
\frac{4}{65} + \frac{3}{s} + \frac{2}{s^2} + \frac{2}{s^3} + \frac{2}{s^4} \]

\[ = \frac{65s^4 + 490s^3 + 528s^2 + 400s + 128}{s(5s^2 + 30s + 16)} \]

The complex input admittance \( Y(s) \) is found by taking the reciprocal of \( Z(s) \), that is, \[ Y(s) = \frac{1}{Z(s)} = \frac{s(5s^2 + 30s + 16)}{65s^4 + 490s^3 + 528s^2 + 400s + 128} \] (6.23)
6.4 Transfer Functions

In an s-domain circuit, the ratio of the output voltage $V_{out}(s)$ to the input voltage $V_{in}(s)$ under zero state conditions, is of great interest in network analysis. This ratio is referred to as the voltage transfer function and it is denoted as $G_v(s)$, that is,

$$G_v(s) = \frac{V_{out}(s)}{V_{in}(s)} \quad (6.24)$$

Similarly, the ratio of the output current $I_{out}(s)$ to the input current $I_{in}(s)$ under zero state conditions, is called the current transfer function denoted as $G_i(s)$, that is,

$$G_i(s) = \frac{I_{out}(s)}{I_{in}(s)} \quad (6.25)$$

The current transfer function of (6.25) is rarely used; therefore, from now on, the transfer function will have the meaning of the voltage transfer function, i.e.,

$$G(s) = \frac{V_{out}(s)}{V_{in}(s)} \quad (6.26)$$

Example 6.6

Derive an expression for the transfer function $G(s)$ for the circuit of Figure 6.17, where $R_g$ represents the internal resistance of the applied (source) voltage $V_S$, and $R_L$ represents the resistance of the load that consists of $R_L$, $L$, and $C$.

![Figure 6.17. Circuit for Example 6.6](image)

Solution:

No initial conditions are given, and even if they were, we would disregard them since the transfer function was defined as the ratio of the output voltage $V_{out}(s)$ to the input voltage $V_{in}(s)$ under zero state conditions.
zero initial conditions. The $s$-domain circuit is shown in Figure 6.18.

![Figure 6.18. The $s$-domain circuit for Example 6.6](image)

The transfer function $G(s)$ is readily found by application of the voltage division expression of the $s$-domain circuit of Figure 6.18, i.e.,

$$V_{out}(s) = \frac{R_L + sL + 1/sC}{R_g + R_L + sL + 1/sC}V_{in}(s)$$

Then,

$$G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{R_L + sL + 1/sC}{R_g + R_L + sL + 1/sC}$$

(6.27)

**Example 6.7**

Compute the transfer function $G(s)$ for the circuit of Figure 6.19 in terms of the circuit constants $R_1, R_2, R_3, C_1,$ and $C_2$. Then, replace the complex variable $s$ with $j\omega$, and the circuit constants with their numerical values and plot the magnitude $|G(s)| = V_{out}(s)/V_{in}(s)$ versus radian frequency $\omega$.

![Figure 6.19. Circuit for Example 6.7](image)

**Solution:**

The $s$-domain equivalent circuit is shown in Figure 6.20.
Next, we write nodal equations at nodes 1 and 2.

At node 1,

\[
\frac{V_1(s) - V_{in}(s)}{R_1} + \frac{V_1(s)}{1/sC_1} + \frac{V_1(s) - V_{out}(s)}{R_2} + \frac{V_1(s) - V_2(s)}{R_3} = 0
\]  \hspace{1cm} (6.28)

At node 2,

\[
\frac{V_2(s) - V_1(s)}{R_3} = \frac{V_{out}(s)}{1/sC_2}
\]  \hspace{1cm} (6.29)

Since \(V_2(s) = 0\) (virtual ground), we express (6.29) as

\[
V_1(s) = (-sR_3C_2)V_{out}(s)
\]  \hspace{1cm} (6.30)

and by substitution of (6.30) into (6.28), rearranging, and collecting like terms, we get:

\[
\left[\left(\frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3} + sC_1\right)\left(-sR_3C_2\right) - \frac{1}{R_2}\right] V_{out}(s) = \frac{1}{R_1}V_{in}(s)
\]

or

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{-1}{R_1\left[\left(\frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3} + sC_1\right)\left(sR_3C_2\right) + \frac{1}{R_2}\right]}
\]  \hspace{1cm} (6.31)

By substitution of \(s\) with \(j\omega\) and the given values for resistors and capacitors, we get

\[
G(j\omega) = \frac{-1}{2 \times 10^5\left[\left(\frac{1}{20 \times 10^3} + j2.5 \times 10^{-8}\omega\right)\left(j5 \times 10^4 \times 10^{-8}\omega\right) + \frac{1}{4 \times 10^4}\right]}
\]

or
We use MATLAB to plot the magnitude of (6.32) on a semilog scale with the following code:

```matlab
w=1:10:10000; Gs=1./(2.5.*10.^(-6).*w.^2 - 5.*j.*10.^(-3).*w+5);
semilogx(w,abs(Gs)); grid; hold on
xlabel('Radian Frequency w'); ylabel('|Vout/Vin|');
title('Magnitude Vout/Vin vs. Radian Frequency')
```

The plot is shown in Figure 6.21. We observe that the given op amp circuit is a second order low-pass filter whose cutoff frequency (−3 dB) occurs at about 700 r/s.
6.5 Summary

- The Laplace transformation provides a convenient method of analyzing electric circuits since integrodifferential equations in the \( t \)-domain are transformed to algebraic equations in the \( s \)-domain.

- In the \( s \)-domain the terms \( sL \) and \( 1/sC \) are called complex inductive impedance, and complex capacitive impedance respectively. Likewise, the terms and \( sC \) and \( 1/sL \) are called complex capacitive admittance and complex inductive admittance respectively.

- The expression

\[
Z(s) = R + sL + \frac{1}{sC}
\]

is a complex quantity, and it is referred to as the complex input impedance of an \( s \)-domain \( RLC \) series circuit.

- In the \( s \)-domain the current \( I(s) \) can be found from

\[
I(s) = \frac{V_s(s)}{Z(s)}
\]

- The expression

\[
Y(s) = G + \frac{1}{sL} + sC
\]

is a complex quantity, and it is referred to as the complex input admittance of an \( s \)-domain \( GLC \) parallel circuit.

- In the \( s \)-domain the voltage \( V(s) \) can be found from

\[
V(s) = \frac{I_s(s)}{Y(s)}
\]

- In an \( s \)-domain circuit, the ratio of the output voltage \( V_{out}(s) \) to the input voltage \( V_{in}(s) \) under zero state conditions is referred to as the voltage transfer function and it is denoted as \( G(s) \), that is,

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)}
\]
6.6 Exercises

1. In the circuit of Figure 6.22, switch $S$ has been closed for a long time, and opens at $t = 0$. Use the Laplace transform method to compute $i_L(t)$ for $t > 0$.

![Figure 6.22. Circuit for Exercise 1](image)

2. In the circuit of Figure 6.23, switch $S$ has been closed for a long time, and opens at $t = 0$. Use the Laplace transform method to compute $v_c(t)$ for $t > 0$.

![Figure 6.23. Circuit for Exercise 2](image)

3. Use mesh analysis and the Laplace transform method, to compute $i_1(t)$ and $i_2(t)$ for the circuit of Figure 6.24, given that $i_L(0^-) = 0$ and $v_c(0^-) = 0$.

![Figure 6.24. Circuit for Exercise 3](image)
4. For the \( s \)-domain circuit of Figure 6.25,

a. compute the admittance \( Y(s) = \frac{I_1(s)}{V_1(s)} \)

b. compute the \( t \)-domain value of \( i_1(t) \) when \( v_1(t) = u_0(t) \), and all initial conditions are zero.

![Figure 6.25. Circuit for Exercise 4](image)

5. Derive the transfer functions for the networks (a) and (b) of Figure 6.26.

![Figure 6.26. Networks for Exercise 5](image)

6. Derive the transfer functions for the networks (a) and (b) of Figure 6.27.

![Figure 6.27. Networks for Exercise 6](image)

7. Derive the transfer functions for the networks (a) and (b) of Figure 6.28.
8. Derive the transfer function for the networks (a) and (b) of Figure 6.29.

9. Derive the transfer function for the network of Figure 6.30. Using MATLAB, plot $|G(s)|$ versus frequency in Hertz, on a semilog scale.
6.7 Solutions to Exercises

1. At $t = 0^-$ the $t$-domain circuit is as shown below and the 20 $\Omega$ resistor is shorted out by the inductor.

Then,

$$i_L(t)\bigg|_{t=0^-} = \frac{32}{10} = 3.2 \text{ A}$$

and thus the initial condition has been established as $i_L(0^-) = 3.2 \text{ A}$

For all $t > 0$ the $t$-domain and $s$-domain circuits are as shown below.

From the $s$-domain circuit above we get

$$I_L(s) = \frac{3.2 \times 10^{-3}}{20 + 10^{-3}s} = \frac{3.2}{s + 20000} \Leftrightarrow 3.2e^{-20000t} u_0(t) = i_L(t)$$

2. At $t = 0^-$ the $t$-domain circuit is as shown below.

Then,
Therefore, the initial condition is

\[ v_c(0^-) = (20 \, \text{K}\Omega + 10 \, \text{K}\Omega) \cdot i_2(0^-) = (30 \, \text{K}\Omega) \cdot (1 \, \text{mA}) = 30 \, \text{V} \]

For all \( t > 0 \) the \( s \)-domain circuit is as shown below.

3. The \( s \)-domain circuit is shown below where \( z_1 = 2s \), \( z_2 = 1 + 1/s \), and \( z_3 = s + 3 \)
Then,

\[(z_1 + z_2)I_1(s) - z_2I_2(s) = \frac{1}{s}\]
\[-z_2I_1(s) + (z_2 + z_3)I_2(s) = -\frac{2}{s}\]

and in matrix form

\[
\begin{bmatrix}
(z_1 + z_2) & -z_2 \\
-z_2 & (z_2 + z_3)
\end{bmatrix}
\begin{bmatrix}
I_1(s) \\
I_2(s)
\end{bmatrix}
= \begin{bmatrix}
\frac{1}{s} \\
-\frac{2}{s}
\end{bmatrix}
\]

Using MATLAB we get

\[
Z = [z_1 + z_2 - z_2; -z_2 z_2 + z_3]; Vs = [1/s - 2/s]; Is = Z \backslash Vs; fprintf('\n'); disp('Is1 = '); pretty(Is(1)); disp('Is2 = '); pretty(Is(2))
\]

\[
Is1 = \frac{2}{(6s + 3 + 9s^2 + 2s)s} \quad \text{conj}(s)
\]

\[
Is2 = \frac{2}{(6s + 3 + 9s^2 + 2s)s} \quad \text{conj}(s)
\]

Therefore,

\[
I_1(s) = \frac{s^2 + 2s - 1}{2s^3 + 9s^2 + 6s + 3} \quad (1)
\]

\[
I_2(s) = -\frac{4s^2 + s + 1}{2s^3 + 9s^2 + 6s + 3} \quad (2)
\]

We express the denominator of (1) as a product of a linear and quadratic term using MATLAB.

\[
p = [2 \ 9 \ 6 \ 3]; r = \text{roots}(p); fprintf('\n'); disp('root1 = '); disp(r(1)); disp('root2 = '); disp(r(2)); disp('root3 = '); disp(r(3)); disp('root2 + root3 = '); disp(r(2) + r(3)); disp('root2 * root3 = '); disp(r(2)*r(3))
\]

\[
\text{root1} = \quad -3.8170
\]

\[
\text{root2} =
\]
-0.3415 + 0.5257i  
\text{root3} =  
-0.3415 - 0.5257i  
\text{root2} + \text{root3} =  
-0.6830  
\text{root2} * \text{root3} =  
0.3930  

and with these values (1) is written as

\begin{equation}
I_f(s) = \frac{s^2 + 2s - 1}{(s + 3.817) \cdot (s^2 + 0.683s + 0.393)} = \frac{r_1}{(s + 3.817)} + \frac{r_2s + r_3}{(s^2 + 0.683s + 0.393)} \tag{3}
\end{equation}

Multiplying every term by the denominator and equating numerators we get

\[s^2 + 2s - 1 = r_1(s^2 + 0.683s + 0.393) + (r_2s + r_3)(s + 3.817)\]

Equating \(s^2, s,\) and constant terms we get

\[r_1 + r_2 = 1\]
\[0.683r_1 + 3.817r_2 + r_3 = 2\]
\[0.393r_1 + 3.817r_3 = -1\]

We will use MATLAB to find these residues.

\[A = [1 \ 1 \ 0; 0.683 \ 3.817 \ 1; 0.393 \ 0 \ 3.817]; \ B = [1 \ 2 \ -1]; \ r = A \backslash B; \text{fprintf('n\n');...}\]
\[\text{fprintf('r1 = \%5.2f \n',r(1)); \text{fprintf('r2 = \%5.2f \n',r(2)); \text{fprintf('r3 = \%5.2f',r(3))}\n}}\]
\[r1 = 0.48 \quad r2 = 0.52 \quad r3 = -0.31\]

By substitution of these values into (3) we get

\begin{equation}
I_f(s) = \frac{r_1}{(s + 3.817)} + \frac{r_2s + r_3}{(s^2 + 0.683s + 0.393)} = \frac{0.48}{(s + 3.817)} + \frac{0.52s - 0.31}{(s^2 + 0.683s + 0.393)} \tag{4}
\end{equation}

By inspection, the Inverse Laplace of first term on the right side of (4) is

\[\frac{0.48}{(s + 3.82)} \Leftrightarrow 0.48e^{-3.82t} \quad (5)\]

The second term on the right side of (4) requires some manipulation. Therefore, we will use the MATLAB \texttt{ilaplace(s)} function to find the Inverse Laplace as shown below.

\begin{verbatim}
syms s t
IL=ilaplace((0.52*s-0.31)/(s^2+0.68*s+0.39));
\end{verbatim}
Thus,

\[ i_1(t) = 0.48e^{-3.82t} - 0.93e^{-0.34t} \sin 0.53t + 0.52e^{-0.34t} \cos 0.53t \]

Next, we will find \( I_2(s) \). We found earlier that

\[ I_2(s) = -\frac{4s^2 + s + 1}{2s^3 + 9s^2 + 6s + 3} \]

and following the same procedure we have

\[ I_2(s) = \frac{-4s^2 - s - 1}{(s + 3.817) \cdot (s^2 + 0.683s + 0.393)} = \frac{r_1}{(s + 3.817)} + \frac{r_2s + r_3}{(s^2 + 0.683s + 0.393)} \tag{6} \]

Multiplying every term by the denominator and equating numerators we get

\[-4s^2 - s - 1 = r_1(s^2 + 0.683s + 0.393) + (r_2s + r_3)(s + 3.817)\]

Equating \( s^2, s, \) and constant terms we get

\[ r_1 + r_2 = -4 \]
\[ 0.683r_1 + 3.817r_2 + r_3 = -1 \]
\[ 0.393r_1 + 3.817r_3 = -1 \]

We will use MATLAB to find these residues.

\[ A=[1 \ 1 \ 0; 0.683 \ 3.817 \ 1; 0.393 \ 0 \ 3.817]; B=[-4 \ -1 \ -1]'; r=A\B; fprintf('r1 = %5.2f\t\t',r(1)); fprintf('r2 = %5.2f\t\t',r(2)); fprintf('r3 = %5.2f',r(3)) \]

\[ r1 = -4.49 \quad r2 = 0.49 \quad r3 = 0.20 \]

By substitution of these values into (6) we get

\[ I_1(s) = \frac{r_1}{(s + 3.817)} + \frac{r_2s + r_3}{(s^2 + 0.683s + 0.393)} = \frac{-4.49}{(s + 3.817)} + \frac{0.49s + 0.20}{(s^2 + 0.683s + 0.393)} \tag{7} \]

By inspection, the Inverse Laplace of first term on the right side of (7) is
The second term on the right side of (7) requires some manipulation. Therefore, we will use the MATLAB \texttt{ilaplace(s)} function to find the Inverse Laplace as shown below.

```matlab
syms s t
IL=ilaplace((0.49*s-0.20)/(s^2+0.68*s+0.39)); pretty(IL)

\[
\begin{align*}
&167 / 9800 \quad 17 / 50 \quad 1/2 / 1/2 \\
&\quad \exp(-14 t) \quad \sin(7/50 t) / 49 / 14 / 50 \\
&\quad + \quad \exp(-14 t) \cos(7/50 t) / 100 / 50
\end{align*}
\]

Thus,

\[i_2(t) = -4.47e^{-3.82t} + 0.06e^{-0.34t} \sin 0.53t + 0.49e^{-0.34t} \cos 0.53t\]

4.

\[\begin{aligned}
V_C(s) &+ \quad 1/s \\
V_i(s) &+ \quad I_1(s) \\
I_2(s) &+ \quad 2 \\
V_2(s) &+ \quad 2V_C(s)
\end{aligned}\]

a. Mesh 1:

\[(2 + 1/s) \cdot I_1(s) - I_2(s) = V_i(s)\]

or

\[6(2 + 1/s) \cdot I_1(s) - 6I_2(s) = 6V_i(s) \quad (1)\]

Mesh 2:

\[-I_1(s) + 6I_2(s) = -V_2(s) = -(2/s)I_1(s) \quad (2)\]

Addition of (1) and (2) yields

\[(12 + 6/s) \cdot I_1(s) + (2/s - 1) \cdot I_i(s) = 6V_i(s)\]

or

\[(11 + 8/s) \cdot I_i(s) = 6V_i(s)\]
and thus

\[ Y(s) = \frac{I_j(s)}{V_j(s)} = \frac{6}{11 + 8/s} = \frac{6s}{11s + 8} \]

b. With \( V_j(s) = 1/s \) we get

\[ I_j(s) = Y(s) \cdot V_j(s) = \frac{6s}{11s + 8} \cdot \frac{1}{s} = \frac{6}{11s + 8} = \frac{6/11}{s + 8/11} \Rightarrow \frac{6}{11} e^{-(8/11)t} = i_j(t) \]

5.

Circuit (a):

\[ V_{out}(s) = \frac{1/C_s}{R + 1/C_s} \cdot V_{in}(s) \]

and

\[ G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{1/C_s}{R + 1/C_s} = \frac{1/C_s}{(RCs + 1)/(Cs)} = \frac{1}{RCs + 1} = \frac{1}{RC} s + 1/RC \]

Circuit (b):

\[ V_{out}(s) = \frac{R}{Ls + R} \cdot V_{in}(s) \]

and

\[ G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{R}{Ls + R} = \frac{R/L}{s + R/L} \]

Both of these circuits are first-order low-pass filters.
6.
Circuit (a):
\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{R}{1/CS + R} = \frac{RCs}{RCs + 1} = \frac{s}{s + 1/RC}
\]

and
\[
V_{out}(s) = \frac{R}{1/CS + R} \cdot V_{in}(s)
\]

Circuit (b):
\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{Ls}{R + Ls} = \frac{s}{s + R/L}
\]

Both of these circuits are first-order high-pass filters.

7.
Circuit (a):
\[
V_{out}(s) = \frac{R}{Ls + 1/CS + R} \cdot V_{in}(s)
\]
and
\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{R}{Ls + 1/Cs + R} = \frac{R Cs}{L Cs^2 + 1 + R Cs} = \frac{(R/L)s}{s^2 + (R/L)s + 1/LC}
\]

This circuit is a second-order band-pass filter.

Circuit (b):

\[
V_{out}(s) = \frac{Ls + 1/Cs}{R + Ls + 1/Cs} \cdot V_{in}(s)
\]

and
\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{Ls + 1/Cs}{R + Ls + 1/Cs} = \frac{LCs^2 + 1}{LCs^2 + RCs + 1} = \frac{s^2 + 1/LC}{s^2 + (R/L)s + 1/LC}
\]

This circuit is a second-order band-elimination (band-reject) filter.

8.
Circuit (a):

Let \( z_1 = R_1 \) and \( z_2 = \frac{R_2 \times 1/Cs}{R_2 + 1/Cs} \) and since for inverting op-amp \( \frac{V_{out}(s)}{V_{in}(s)} = \frac{z_2}{z_1} \), for this circuit

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{[(R_2 \times 1/Cs)/(R_2 + 1/Cs)]}{R_1} = \frac{-(R_2 \times 1/Cs)}{R_1 \cdot (R_2 + 1/Cs)} = -\frac{R_1 C}{s + 1/R_2 C}
\]

This circuit is a first-order active low-pass filter.
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Circuit (b):

\[
\begin{align*}
R_1 & \quad C \\
\downarrow & \quad \uparrow \\
V_{in}(s) & \quad V_{ou}(s)
\end{align*}
\]

Let \( z_1 = R_1 + 1/Cs \) and \( z_2 = R_2 \) and since for inverting op-amp \( \frac{V_{out}(s)}{V_{in}(s)} = \frac{z_2}{z_1} \), for this circuit

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{-R_2}{R_1 + 1/Cs} = \frac{-(R_2/R_1)s}{s + 1/R_1C}
\]

This circuit is a first-order active high-pass filter.

9.

\[
R_1 = 11.3 \, \text{K}\Omega \\
R_2 = 22.6 \, \text{K}\Omega \\
R_3 = R_4 = 68.1 \, \text{K}\Omega \\
C_1 = C_2 = 0.01 \, \mu\text{F}
\]

At Node \( V_1 \):

\[
\frac{V_1(s)}{R_3} + \frac{V_1(s) - V_{out}(s)}{R_4} = 0
\]

or

\[
\left( \frac{1}{R_3} + \frac{1}{R_4} \right) V_1(s) = \frac{1}{R_4} V_{out}(s) \quad (1)
\]

At Node \( V_3 \):

\[
\frac{V_3(s) - V_2(s)}{R_2} + \frac{V_3(s)}{1/C_1s} = 0
\]
and since \( V_3(s) = V_i(s) \) we express the last relation above as

\[
\frac{V_i(s) - V_2(s)}{R_2} + C_jsV_i(s) = 0
\]

or

\[
\left( \frac{1}{R_2} + C_js \right)V_i(s) = \frac{1}{R_2}V_2(s) \quad (2)
\]

At Node \( V_2 \):

\[
\frac{V_2(s) - V_{in}(s)}{R_1} + \frac{V_2(s) - V_1(s)}{R_2} + \frac{V_2(s) - V_{out}(s)}{1/C_2s} = 0
\]

or

\[
\left( \frac{1}{R_1} + \frac{R_2}{R_2} + C_2s \right)V_2(s) = \frac{V_{in}(s)}{R_1} + \frac{V_1(s)}{R_2} + C_2sV_{out}(s) \quad (3)
\]

From (1)

\[
V_i(s) = \frac{(1/R_4)}{(R_3 + R_4)/(R_3R_4)}V_{out}(s) = \frac{R_3}{(R_3 + R_4)}V_{out}(s) \quad (4)
\]

From (2)

\[
V_2(s) = R_2\left( \frac{1}{R_2} + C_js \right)V_i(s) = (1 + R_2C_1s)V_i(s)
\]

and with (4)

\[
V_2(s) = \frac{R_3(1 + R_2C_1s)}{(R_3 + R_4)}V_{out}(s) \quad (5)
\]

By substitution of (4) and (5) into (3) we get

\[
\left( \frac{1}{R_1} + \frac{1}{R_2} + C_2s \right)\frac{R_3(1 + R_2C_1s)}{(R_3 + R_4)}V_{out}(s) = \frac{V_{in}(s)}{R_1} + \frac{R_3}{R_2(R_3 + R_4)}V_{out}(s) + C_2sV_{out}(s)
\]

or

\[
\left[ \left( \frac{1}{R_1} + \frac{1}{R_2} + C_2s \right)\frac{R_3(1 + R_2C_1s)}{(R_3 + R_4)} - \frac{1}{R_2(R_3 + R_4)} - C_2s \right]V_{out}(s) = \frac{1}{R_1}V_{in}(s)
\]

and thus

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{1}{R_1 \left[ \left( \frac{1}{R_1} + \frac{1}{R_2} + C_2s \right)\frac{R_3(1 + R_2C_1s)}{(R_3 + R_4)} - \frac{1}{R_2(R_3 + R_4)} - C_2s \right]}
\]

By substitution of the given values and after simplification we get
The plot above indicates that this circuit is a second-order low-pass filter.
Chapter 7

Frequency Response and Bode Plots

This chapter discusses frequency response in terms of both amplitude and phase. This topic will enable us to determine which frequencies are dominant and which frequencies are virtually suppressed. The design of electric filters is based on the study of the frequency response. We will also discuss the Bode method of linear system analysis using two separate plots; one for the magnitude of the transfer function, and the other for the phase, both versus frequency. These plots reveal valuable information about the frequency response behavior.

Note: Throughout this text, the common (base 10) logarithm of a number \( x \) will be denoted as \( \log(x) \) while its natural (base e) logarithm will be denoted as \( \ln(x) \). However, we should remember that in MATLAB the \( \log(x) \) function displays the natural logarithm, and the common (base 10) logarithm is defined as \( \log_{10}(x) \).

7.1 Decibels

The ratio of any two values of the same quantity (power, voltage or current) can be expressed in decibels (\( dB \)). For instance, we say that an amplifier has 10 \( dB \) power gain or a transmission line has a power loss of 7 \( dB \) (or gain \(-7 \, dB\)). If the gain (or loss) is 0 \( dB \), the output is equal to the input. We should remember that a negative voltage or current gain \( A_V \) or \( A_I \) indicates that there is a 180° phase difference between the input and the output waveforms. For instance, if an amplifier has a gain of \( -100 \) (dimensionless number), it means that the output is 180° out-of-phase with the input. For this reason we use absolute values of power, voltage and current when these are expressed in \( dB \) terms to avoid misinterpretation of gain or loss.

By definition,

\[
dB = 10 \log \left( \frac{P_{out}}{P_{in}} \right) \quad (7.1)
\]

Therefore,

10 \( dB \) represents a power ratio of 10

10\( n \) \( dB \) represents a power ratio of \( 10^n \)

20 \( dB \) represents a power ratio of 100

30 \( dB \) represents a power ratio of 1,000

60 \( dB \) represents a power ratio of 1,000,000
Also,

1 dB represents a power ratio of approximately 1.25
3 dB represents a power ratio of approximately 2
7 dB represents a power ratio of approximately 5

From these, we can estimate other values. For instance, $4 \, dB = 3 \, dB + 1 \, dB$ which is equivalent to a power ratio of approximately $2 \times 1.25 = 2.5$. Likewise, $27 \, dB = 20 \, dB + 7 \, dB$ and this is equivalent to a power ratio of approximately $100 \times 5 = 500$.

Since $y = \log x^2 = 2 \log x$ and $P = V^2 / R = I^2 R$, if we let $R = 1$ the dB values for the voltage and current ratios become:

$$ dB_v = 10 \log \left| \frac{V_{out}}{V_{in}} \right|^2 = 20 \log \left| \frac{V_{out}}{V_{in}} \right| $$

and

$$ dB_i = 10 \log \left| \frac{I_{out}}{I_{in}} \right|^2 = 20 \log \left| \frac{I_{out}}{I_{in}} \right| $$

Example 7.1

Compute the gain in $dB_W$ for the amplifier shown in Figure 7.1.

Solution:

$$ dB_W = 10 \log \frac{P_{out}}{P_{in}} = 10 \log \frac{10}{1} = 10 \log 10 = 10 \times 1 = 10 \, dB_W $$

Example 7.2

Compute the gain in $dB_V$ for the amplifier shown in Figure 7.2 given that $\log 2 = 0.3$.

$$ dB_V = 10 \log \frac{V_{out}}{V_{in}} = 10 \log \frac{2}{1} = 10 \log 2 = 10 \times 0.3 = 3 \, dB_V $$
Solution:

\[ dB_V = 20 \log \frac{V_{out}}{V_{in}} = 20 \log \frac{2}{1} = 20 \log 0.3 = 20 \times 0.3 = 6 \ dB_V \]

### 7.2 Bandwidth and Frequency Response

Electric and electronic circuits, such as filters and amplifiers, exhibit a band of frequencies over which the output remains nearly constant. Consider, for example, the magnitude of the output voltage \( |V_{out}| \) of an electric or electronic circuit as a function of radian frequency \( \omega \) as shown in Figure 7.3.

![Figure 7.3. Definition of the bandwidth.](image)

As shown in Figure 7.3, the **bandwidth** is \( BW = \omega_2 - \omega_1 \) where \( \omega_1 \) and \( \omega_2 \) are the **lower** and **upper cutoff frequencies** respectively. At these frequencies, \( |V_{out}| = \sqrt{2}/2 = 0.707 \) and these two points are known as the **3 dB down** or **half-power points**. They derive their name from the fact that since power \( p = v^2/R = i^2 R \), for \( R = 1 \) and for \( v = 0.707 |V_{out}| \) or \( i = 0.707 |V_{out}| \) the power is \( 1/2 \), that is, it is “halved”. Alternately, we can define the bandwidth as the frequency band between half-power points.

Most amplifiers are used with a feedback path which returns (feeds) some or all its output to the input as shown in Figure 7.4.

![Figure 7.4. Amplifier with partial output feedback](image)

Figure 7.5 shows an amplifier where the entire output is fed back to the input.
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The symbol $\Sigma$ (Greek capital letter sigma) inside the circle indicates the summing point where the output signal, or portion of it, is combined with the input signal. This summing point may be also indicated with a large plus (+) symbol inside the circle. The positive (+) sign below the summing point implies positive feedback which means that the output, or portion of it, is added to the input. On the other hand, the negative (−) sign implies negative feedback which means that the output, or portion of it, is subtracted from the input. Practically, all amplifiers use used with negative feedback since positive feedback causes circuit instability.

7.3  Octave and Decade

Let us consider two frequencies $u_1$ and $u_2$ defining the frequency interval $u_2 - u_1$, and let

$$u_2 - u_1 = \log_{10} \omega_2 - \log_{10} \omega_1 = \log_{10} \frac{\omega_2}{\omega_1}$$  \hspace{1cm} (7.4)

If these frequencies are such that $\omega_2 = 2\omega_1$, we say that these frequencies are separated by one octave and if $\omega_2 = 10\omega_1$, they are separated by one decade.

Let us now consider a transfer function $G(s)$ whose magnitude is evaluated at $s = j\omega$, that is,

$$|G(s)| = \frac{C}{s^k} \bigg|_{s = j\omega} = |G(\omega)| = \frac{C}{\omega^k}$$  \hspace{1cm} (7.5)

Taking the log of both sides of (7.5) and multiplying by 20, we get

$$20\log_{10}|G(\omega)| = 20\log_{10}C - 20\log_{10}\omega^k = -20k\log_{10}\omega + 20\log_{10}C$$

or

$$|G(\omega)|_{dB} = -20k\log_{10}\omega + 20\log_{10}C$$  \hspace{1cm} (7.6)

Relation (7.6) is an equation of a straight line in a semilog plot with abscissa $log_{10}\omega$ where

$$\text{slope} = -20k \frac{dB}{\text{decade}}$$

and intercept $= C \ dB$ shown in Figure 7.6.
With these concepts in mind, we can now proceed to discuss Bode Plots and Asymptotic Approximations.

### 7.4 Bode Plot Scales and Asymptotic Approximations

Bode plots are magnitude and phase plots where the abscissa (frequency axis) is a logarithmic (base 10) scale, and the radian frequency \( \omega \) is equally spaced between powers of 10 such as \( 10^{-1}, 10^0, 10^1, 10^2 \) and so on.

The ordinate (\( dB \) axis) of the magnitude plot has a scale in \( dB \) units, and the ordinate of the phase plot has a scale in degrees as shown in Figure 7.7.
It is convenient to express the magnitude in \( dB \) so that a transfer function \( G(s) \), composed of products of terms can be computed by the sum of the \( dB \) magnitudes of the individual terms. For example,

\[
\frac{20 \cdot \left( 1 + \frac{j\omega}{100} \right)}{1 + j\omega} = 20 \ dB + \left( 1 + \frac{j\omega}{100} \right) dB + \frac{1}{1 + j\omega} dB
\]

and the Bode plots then can be approximated by straight lines called *asymptotes*.

### 7.5 Construction of Bode Plots when the Zeros and Poles are Real

Let us consider the transfer function

\[
G(s) = \frac{A \cdot (s + z_1) \cdot (s + z_2) \cdot \ldots \cdot (s + z_m)}{s \cdot (s + p_1) \cdot (s + p_2) \cdot \ldots \cdot (s + p_n)} \quad (7.7)
\]

where \( A \) is a real constant, and the zeros \( z_i \) and poles \( p_i \) are real numbers. We will consider complex zeros and poles in the next section. Letting \( s = j\omega \) in (7.7) we get

\[
G(j\omega) = \frac{A \cdot (j\omega + z_1) \cdot (j\omega + z_2) \cdot \ldots \cdot (j\omega + z_m)}{j\omega \cdot (j\omega + p_1) \cdot (j\omega + p_2) \cdot \ldots \cdot (j\omega + p_n)} \quad (7.8)
\]

Next, we multiply and divide each numerator factor \( j\omega + z_i \) by \( z_i \) and each denominator factor \( j\omega + p_i \) by \( p_i \) and we get:

\[
G(j\omega) = \frac{A \cdot z_1 \left( \frac{j\omega}{z_1} + 1 \right) \cdot z_2 \left( \frac{j\omega}{z_2} + 1 \right) \cdot \ldots \cdot z_m \left( \frac{j\omega}{z_m} + 1 \right)}{j\omega \cdot p_1 \left( \frac{j\omega}{p_1} + 1 \right) \cdot p_2 \left( \frac{j\omega}{p_2} + 1 \right) \cdot \ldots \cdot p_n \left( \frac{j\omega}{p_n} + 1 \right)} \quad (7.9)
\]

Letting

\[
K = \frac{A \cdot z_1 \cdot z_2 \cdot \ldots \cdot z_m}{p_1 \cdot p_2 \cdot \ldots \cdot p_n} = A \prod_{i=1}^{m} z_i \prod_{i=1}^{n} p_i \quad (7.10)
\]

we can express (7.9) in \( dB \) magnitude and phase form,
The constant $K$ can be positive or negative. Its magnitude is $|K|$ and its phase angle is $0^\circ$ if $K > 0$, and $-180^\circ$ if $K < 0$. The magnitude and phase plots for the constant $K$ are shown in Figure 7.8.

For a zero of order $n$, that is, $(j\omega)^n$ at the origin, the Bode plots for the magnitude and phase are as shown in Figures 7.9 and 7.10 respectively.

For a pole of order $n$, that is, $1/(j\omega)^n = (j\omega)^{-n}$ at the origin, the Bode plots are as shown in Figures 7.11 and 7.12 respectively.

Next, we consider the term $G(j\omega) = (a+j\omega)^n$.

The magnitude of this term is

$$|G(j\omega)| = \sqrt{(a^2 + \omega^2)^n} = (a^2 + \omega^2)^{n/2}$$

and taking the log of both sides and multiplying by 20 we get

$$20\log|G(j\omega)| = 10n\log(a^2 + \omega^2)$$

It is convenient to normalize (7.14) by letting

$$u = \frac{\omega}{a}$$
Then, (7.14) becomes

\[ 20 \log |G(\omega)| = 10n \log \left( \frac{a^2 \cdot \frac{a^2 + \omega^2}{a^2}}{a^2} \right) = 10n \log a^2 + 10n \log (1 + u^2) \]

\[ = 10n \log (1 + u^2) + 20 \log a \]
For $u \ll 1$ the first term of (7.16) becomes $10n \log 1 = 0 \, \text{dB}$. For $u \gg 1$, this term becomes approximately $10n \log u^2 = 20n \log u$ and this has the same form as $G(j\omega) = (j\omega)^n$ which is shown in Figure 7.9 for $n = 1, n = 2$, and $n = 3$.

The frequency at which two asymptotes intersect each other forming a corner is referred to as the corner frequency. Thus, the two lines defined by the first term of (7.16), one for $u \ll 1$ and the other for $u \gg 1$ intersect at the corner frequency $u = 1$.

The second term of (7.16) represents the ordinate axis intercept defined by this straight line.
The phase response for the term $G(j\omega) = (a + j\omega)^n$ is found as follows:

We let

$$u = \frac{\omega}{a} \tag{7.17}$$

and

$$\phi(u) = \tan^{-1}u \tag{7.18}$$

Then,

$$(a + j\omega)^n = a^n (1 + ju)^n = a^n \left(\sqrt{1 + u^2} \angle \tan^{-1}u\right)^n = a^n (1 + u^2)^{n/2} e^{in\phi(u)} \tag{7.19}$$

Figure 7.13 shows plots of the magnitude of (7.16) for $a = 10$, $n = 1$, $n = 2$, and $n = 3$.

![Figure 7.13. Magnitude for zeros of Order n for $(a+j\omega)^n$](image)

As shown in Figure 7.13, a quick sketch can be obtained by drawing the straight line asymptotes given by $10\log 1 = 0$ and $10n\log u^2$ for $u \ll 1$ and $u \gg 1$ respectively.

The phase angle of (7.19) is $n\phi(u)$. Then, with (7.18) and letting

$$n\phi(u) = \theta(u) = n\tan^{-1}u \tag{7.20}$$
we get

\[ \lim_{u \to 0} \theta(u) = \lim_{u \to 0} n \tan^{-1} u = 0 \]  
(7.21)

and

\[ \lim_{u \to \infty} \theta(u) = \lim_{u \to \infty} n \tan^{-1} u = \frac{n\pi}{2} \]  
(7.22)

At the corner frequency \( u = a \) we get \( u = 1 \) and with (7.20)

\[ \theta(1) = n \tan^{-1} 1 = \frac{n\pi}{4} \]  
(7.23)

Figure 7.14 shows the phase angle plot for (7.19).

The magnitude and phase plots for \( G(j\omega) = 1/(a + j\omega)^n \) are similar to those of \( G(j\omega) = (a + j\omega)^n \) except for a minus sign. In this case (7.16) becomes

\[ -20 \log |G(ju)| = -10 n \log (1 + u^2) - 20 \log a \]  
(7.24)

and (7.20) becomes

\[ \theta(u) = -n \tan^{-1} u \]  
(7.25)

The plots for (7.24) and (7.25) are shown in Figures 7.15 and 7.16 respectively.
7.6 Construction of Bode Plots when the Zeros and Poles are Complex

The final type of terms appearing in the transfer function $G(s)$ are quadratic term of the form $as^2 + bs + c$ whose roots are complex conjugates. In this case, we express the complex conjugate roots as $1/(a+j\omega)^n$.
Construction of Bode Plots when the Zeros and Poles are Complex

\[(s + \alpha - j\beta)(s + \alpha + j\beta) = (s + \alpha)^2 + \beta^2\]
\[= s^2 + 2\alpha s + \alpha^2 + \beta^2\]  \hspace{1cm} (7.26)

and letting
\[\alpha = \zeta \omega_n\]  \hspace{1cm} (7.27)

and
\[\alpha^2 + \beta^2 = \omega_n^2\]  \hspace{1cm} (7.28)

by substitution into (7.26) we get
\[s^2 + 2\alpha s + \alpha^2 + \beta^2 = s^2 + 2\zeta \omega_n s + \omega_n^2\]  \hspace{1cm} (7.29)

Next, we let
\[G(s) = s^2 + 2\zeta \omega_n s + \omega_n^2\]  \hspace{1cm} (7.30)

Then,
\[G(j\omega) = (j\omega)^2 + j2\omega_n \omega + \omega_n^2\]
\[= (\omega_n^2 - \omega^2) + j2\omega_n \omega\]  \hspace{1cm} (7.31)

The magnitude of (7.31) is
\[|G(j\omega)| = \sqrt{(\omega_n^2 - \omega^2)^2 + 4\zeta^2 \omega_n^2 \omega^2}\]  \hspace{1cm} (7.32)

and taking the log of both sides and multiplying by 20 we get
\[20 \log |G(j\omega)| = 10 \log [(\omega_n^2 - \omega^2)^2 + 4\zeta^2 \omega_n^2 \omega^2]\]  \hspace{1cm} (7.33)

As in the previous section, it is convenient to normalize (7.33) by dividing by \(\omega_n^4\) to yield a function of the normalized frequency variable \(u\) such that
\[u = \omega / \omega_n\]  \hspace{1cm} (7.34)

Then, (7.33) is expressed as
\[20 \log |G(ju)| = 10 \log [(\omega_n^2 - \omega^2)^2 + 4\zeta^2 \omega_n^2 \omega^2]\]
or
The first term in (7.35) is a constant which represents the ordinate axis intercept defined by this straight line. For the second term, if \( u^2 \ll 1 \), this term reduces to approximately \( 10 \log 1 = 0 \text{ dB} \) and if \( u^2 \gg 1 \), this term reduces to approximately \( 10 \log u^4 \) and this can be plotted as a straight line increasing at \( 40 \text{ dB/decade} \). Using these two straight lines as asymptotes for the magnitude curve we see that the asymptotes intersect at the corner frequency \( u = 1 \). The exact shape of the curve depends on the value of \( \zeta \) which is called the damping coefficient.

A plot of (7.35) for \( \zeta = 0.2 \), \( \zeta = 0.4 \), and \( \zeta = 0.707 \) is shown in Figure 7.17.

The phase shift associated with \( (\omega_n^2 - \omega^2) + j2\omega_n \omega \) is also simplified by the substitution \( u \equiv \omega / \omega_n \) and thus

\[
\theta(u) = \tan^{-1}\left( \frac{2\zeta u}{1 - u^2} \right) \tag{7.36}
\]

The two asymptotic relations of (7.36) are

\[
\lim_{u \to 0} \theta(u) = \lim_{u \to 0} \tan^{-1}\left( \frac{2\zeta u}{1 - u^2} \right) = 0 \tag{7.37}
\]

and

\[
\lim_{u \to \infty} \theta(u) = \lim_{u \to \infty} \tan^{-1}\left( \frac{2\zeta u}{1 - u^2} \right) = \pi \tag{7.38}
\]

At the corner frequency \( \omega = \omega_n \), \( u = 1 \) and

\[
\theta(1) = \lim_{u \to 1} \tan^{-1}\left( \frac{2\zeta u}{1 - u^2} \right) = \frac{\pi}{2} \tag{7.39}
\]

A plot of the phase for \( \zeta = 0.2 \), \( \zeta = 0.4 \), and \( \zeta = 0.707 \) is shown in Figure 7.18.
Construction of Bode Plots when the Zeros and Poles are Complex

Figure 7.17. Magnitude for zeros of $10\log \omega_n^4 + 10\log [(1-u^2)^2 + 4\zeta^2 u^2]$ which is given by

$$G(j\omega) = \frac{1}{(\omega_n^2 - \omega^2) + j2\omega_n\omega}$$
are similar to those of
\[ G(j\omega) = (\omega_n^2 - \omega^2) + j2\omega_n\omega \]
extcept for a minus sign. In this case, (7.35) becomes
\[ -10 \log \omega_n^4 - 10 \log [(1 - u^2)^2 + 4\zeta^2 u^2] \]
and (7.36) becomes
\[ \theta(u) = -\tan^{-1} \left( \frac{2\zeta u}{1-u^2} \right) \]
A plot of (7.40) for \( \zeta = 0.2 \), \( \zeta = 0.4 \), and \( \zeta = 0.707 \) is shown in Figure 7.19.

\[ \text{Figure 7.19. Magnitude for poles of } 1/(\omega_n^2 - \omega^2) + j2\omega_n\omega \]
\[ 
\begin{array}{c|c}
\text{Frequency } u (r/s) & \text{Magnitude in dB} \\
\hline
0.01 & 20 \\
0.10 & 10 \\
1.00 & 0 \\
10.00 & -10 \\
100.00 & -20 \\
\end{array}
\]

A plot of the phase for \( \zeta = 0.2 \), \( \zeta = 0.4 \), and \( \zeta = 0.707 \) is shown in Figure 7.20.
Construction of Bode Plots when the Zeros and Poles are Complex

Figure 7.20. Phase for poles of \( \frac{1}{10 \log \omega_n^4 + 10 \log [(1 - u^2)^2 + 4 \zeta^2 u^2]} \)

Example 7.3

For the circuit shown in Figure 7.21

a. Compute the transfer function \( G(s) \).

b. Construct a straight line approximation for the magnitude of the Bode plot.

c. From the Bode plot obtain the values of \( 20 \log |G(j\omega)| \) at \( \omega = 30 \text{ r/s} \) and \( \omega = 4000 \text{ r/s} \). Compare these values with the actual values.

d. If \( v_s(t) = 10 \cos(5000t + 60^\circ) \), use the Bode plot to compute the output \( v_{out}(t) \).

Solution:

a. We transform the given circuit to its equivalent in the \( s – domain \) shown in Figure 7.22.
and by the voltage division expression,

\[
V_{out}(s) = \frac{110}{10^4/s + 0.1s + 110} \cdot V_{in}(s)
\]

Therefore, the transfer function is

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{110s}{0.1s^2 + 110s + 10^4} = \frac{1100s}{s^2 + 1100s + 10^5} = \frac{1100s}{(s+100)(s+1000)}
\]  

(7.42)

b. Letting \( s = j\omega \) we get

\[
G(j\omega) = \frac{1100j\omega}{(j\omega + 100)(j\omega + 1000)}
\]

or in standard form

\[
G(j\omega) = \frac{0.011j\omega}{(1+j\omega/100)(1+j\omega/1000)}
\]  

(7.43)

Letting the magnitude of (7.43) be denoted as \( A \), and expressing it in decibels we get

\[
A_{dB} = 20\log |G(j\omega)| = 20\log 0.011 + 20\log |j\omega| - 20\log \left| \frac{1 + j\omega}{10} \right| - 20\log \left| \frac{1 + j\omega}{1000} \right|
\]  

(7.44)

We observe that the first term on the right side of (7.44) is a constant whose value is 20\log 0.011 = -39.17. The second term is a straight line with slope equal to 20 dB/decade. For \( \omega < 100 \, r/s \) the third term is approximately zero and for \( \omega > 100 \) it decreases with slope equal to -20 dB/decade. Likewise, for \( \omega < 1000 \, r/s \) the fourth term is approximately zero and for \( \omega > 1000 \) it also decreases with slope equal to -20 dB/decade.

For Bode plots we use semilog paper. Instructions to construct semilog paper with Microsoft Excel are provided in Appendix D.
In the Bode plot of Figure 7.23 the individual terms are shown with dotted lines and the sum of these with a solid line.

![Figure 7.23. Magnitude plot of (7.44)](image)

**Figure 7.23. Magnitude plot of (7.44)**

c. The plot of Figure 7.23 shows that the magnitude of (7.43) at $\omega = 30 \text{ r/s}$ is approximately $-9 \text{ dB}$ and at $\omega = 4000 \text{ r/s}$ is approximately $-10 \text{ dB}$. The actual values are found as follows:

At $\omega = 30 \text{ r/s}$, (7.43) becomes

$$G(j30) = \frac{0.011 \times j30}{(1 + j0.3)(1 + j0.03)}$$

and using MATLAB we get

```
g30=0.011*30j/((1+0.3j)*(1+0.03j));...
fprintf(' 
'); fprintf('mag = %6.2f 	',abs(g30));...
fprintf('magdB = %6.2f dB',20*log10(abs(g30))); fprintf(' 
'); fprintf(' 
')
mag = 0.32 magdB = -10.01 dB
```

Therefore,

$$|G(j30)| = 0.32$$

and

$$20\log|G(j30)| = 20\log 0.32 \approx -10 \text{ dB}$$
Likewise, at $\omega = 4000 \text{ r/s}$, (7.43) becomes

$$G(j1000) = \frac{0.11(j4000)}{(1 + j40)(1 + j4)}$$

and using MATLAB we get

```matlab
g4000 = 0.011*4000j/((1+40j)*(1+4j));
fprintf(' 
'); fprintf('mag = %6.2f 	',abs(g4000));
fprintf('magdB = %6.2f dB',20*log10(abs(g4000))); fprintf(' 
'); fprintf(' 
')
mag = 0.27     magdB = -11.48 dB
```

Therefore,

$$|G(j4000)| = 0.27$$

and

$$20\log|G(j4000)| = 20\log 0.27 = -11.48 \text{ dB}$$

d. From the Bode plot of Figure 7.23, we see that the value of $A_{dB}$ at $\omega = 5000 \text{ r/s}$ is approximately $-12 \text{ dB}$. Then, since in general $a_{dB} = 20\log b$, and that $y = \log x$ implies $x = 10^y$, we have

$$|A| = 10^{\left(-\frac{12}{20}\right)} = 0.25$$

and therefore

$$V_{out_{max}} = |A||V_s| = 0.25 \times 10 = 2.5 \text{ V}$$

If we wish to obtain a more accurate value, we substitute $\omega = 5000$ into (7.43) and we get

```matlab
g5000 = 0.011*5000j/((1+50j)*(1+5j));
fprintf(' 
'); fprintf('mag = %6.2f 	',abs(g5000));
fprintf('phase = %6.2f deg.',angle(g5000)*180/pi); fprintf(' 
'); fprintf(' 
')
mag = 0.22     phase = -77.54 deg.
```

$$G(j5000) = \frac{0.011(j5000)}{(1 + j50)(1 + j5)} = 0.22 \angle -77.54$$

Then,

$$V_{out_{max}} = |A| \times 10 = 0.22 \times 10 = 2.2 \text{ V}$$

and in the $t$–domain

$$v_{out}(t) = 2.2 \cos(5000t - 77.54^\circ)$$
We can use the MATLAB function **bode(sys)** to draw the Bode plot of a Linear Time Invariant (LTI) System where **sys = tf(num,den)** creates a continuous-time transfer function **sys** with numerator **num** and denominator **den**, and **tf** creates a transfer function. With this function, the frequency range and number of points are chosen automatically. The function **bode(sys,[wmin,wmax])** draws the Bode plot for frequencies between **wmin** and **wmax** (in radians/second) and the function **bode(sys,w)** uses the user-supplied vector **w** of frequencies, in radians/second, at which the Bode response is to be evaluated. To generate logarithmically spaced frequency vectors, we use the command **logspace(first_exponent,last_exponent,number_of_values)**. For example, to generate plots for 100 logarithmically evenly spaced points for the frequency interval $10^{-1} \leq \omega \leq 10^{2}$ r/s, we use the statement **logspace(−1,2,100)**.

The **bode(sys,w)** function displays both magnitude and phase. If we want to display the magnitude only, we can use the **bodemag(sys,w)** function.

MATLAB requires that we express the numerator and denominator of $G(s)$ as polynomials of $s$ in descending powers.

Let us plot the transfer function of Example 7.3 using MATLAB.

From (7.42),

$$G(s) = \frac{1100s}{s^2 + 1100s + 10^5}$$

and the MATLAB code to generate the magnitude and phase plots is

```matlab
num=[0 1100 0]; den=[1 1100 10^5]; w=logspace(0,5,100); bode(num,den,w)
```

However, since for this example we are interested in the magnitude only, we will use the code

```matlab
num=[0 1100 0]; den=[1 1100 10^5]; sys=tf(num,den);...
 w=logspace(0,5,100); bodemag(sys,w); grid
```

and upon execution, MATLAB displays the plot shown in Figure 7.24.

**Example 7.4**

For the circuit of Example 7.3

a. Draw a Bode phase plot.

b. Using the Bode phase plot estimate the frequency where the phase is zero degrees.

c. Compute the actual frequency where the phase is zero degrees.

d. Find $v_{out}(t)$ if $v_{in}(t) = 10\cos(\omega t + 60^\circ)$ and $\omega$ is the value found in part (c).
Figure 7.24. Bode plot for Example 7.3.

Solution:

a. From (7.43) of Example 7.3

\[ G(j\omega) = \frac{0.011j\omega}{(1 + j\omega/100)(1 + j\omega/1000)} \quad (7.45) \]

and in magnitude-phase form

\[ G(j\omega) = \frac{0.011|\omega|}{|(1 + j\omega/100)||1 + j\omega/1000||} \angle(\alpha - \beta - \gamma) \]

where

\[ \angle \alpha = 90^\circ \quad \angle -\beta = -\tan^{-1}(\omega/100) \quad \angle -\gamma = -\tan^{-1}(\omega/1000) \]

For \( \omega = 100 \)

\[ \angle -\beta = -\tan^{-1}1 = -45^\circ \]

For \( \omega = 1000 \)

\[ \angle -\gamma = -\tan^{-1}1 = -45^\circ \]

The straight-line phase angle approximations are shown in Figure 7.25.
Construction of Bode Plots when the Zeros and Poles are Complex

Figure 7.25. Bode plot for Example 7.4.

Figure 7.26 shows the magnitude and phase plots generated with the following MATLAB code:

```matlab
num = [0 1100 0]; den = [1 1100 10^5]; w = logspace(0, 5, 100); bode(num, den, w)
```

b. From the Bode plot of Figure 7.25 we find that the phase is zero degrees at approximately $\omega = 310 \text{ r/s}$

c. From (7.45)

$$G(j\omega) = \frac{0.011j\omega}{(1 + j\omega/100)(1 + j\omega/1000)}$$

and in magnitude-phase form

$$G(j\omega) = \frac{0.011\omega \angle 90^\circ}{|(1 + j\omega/100)| \angle \tan^{-1}(\omega/100)| (1 + j\omega/1000)| \angle \tan^{-1}(\omega/1000)}$$

The phase will be zero when

$$\tan^{-1}(\omega/100) + \tan^{-1}(\omega/1000) = 90^\circ$$
This is a trigonometric equation and we will solve it for $\omega$ with the `solve(equ)` MATLAB function as follows:

```matlab
syms w; x=solve(atan(w/100)+atan(w/1000)-pi/2); combine(x)
ans =
    316.2278
```

Therefore, $\omega = 316.23 \, \text{r/s}$

d. Evaluating (7.45) at $\omega = 316.23 \, \text{r/s}$ we get:

$$G(j316.23) = \frac{0.011(j316.23)}{(1+j316.23/100)(1+j316.23/1000)}$$

(7.46)

and with MATLAB

```matlab
Gj316=0.011*316.23j/((1+316.23j/100)*(1+316.23j/1000)); fprintf('magGj316 = \%5.2f \n', abs(Gj316)); fprintf('phaseGj316 = \%5.2f deg.', angle(Gj316)*180/pi)
magGj316 = 1.00 phaseGj316 = -0.00 deg.
```

We are given that $|V_{in}| = 10 \, \text{V}$ and with $|G(j316.23)| = 1$ we get...
Corrected Amplitude Plots

The phase angle of the input voltage is given as \( \theta_{in} = 60^\circ \) and with \( \theta(j316.23) = 0^\circ \) we find that the phase angle of the output voltage is

\[ \theta_{out} = \theta_{in} + \theta(j316.23) = 60^\circ + 0^\circ = 60^\circ \]

and thus

\[ V_{out} = 10 \angle 60^\circ \]

or

\[ v_{out}(t) = 10 \cos(316.23t + 60^\circ) \]

### 7.7 Corrected Amplitude Plots

The amplitude plots we have considered thus far are approximate. We can make the straight line more accurate by drawing smooth curves connecting the points at one-half the corner frequency \( \omega_n/2 \), the corner frequency \( \omega_n \) and twice the corner frequency \( 2\omega_n \) as shown in Figure 7.27.

At the corner frequency \( \omega_n \), the value of the amplitude \( A \) in \( dB \) is

\[ A_{dB}\big|_{\omega = \omega_n} = \pm 20 \log |1 + j| = \pm 20 \log \sqrt{2} = \pm 3 \, dB \]  

(7.47)

where the plus (+) sign applies to a first order zero, and the minus (−) to a first order pole.

Similarly,

\[ A_{dB}\big|_{\omega = \omega_n/2} = \pm 20 \log |1 + j/2| = \pm 20 \log \frac{\sqrt{5}}{4} = \pm 0.97 \, dB \approx \pm 1 \, dB \]  

(7.48)

and

\[ A_{dB}\big|_{\omega = 2\omega_n} = \pm 20 \log |1 + j2| = \pm 20 \log \sqrt{5} = \pm 6.99 \, dB \approx \pm 7 \, dB \]  

(7.49)

As we can see from Figure 7.27, the straight line approximations, shown by dotted lines, yield 0 \( dB \) at half the corner frequency and at the corner frequency. At twice the corner frequency, the straight line approximations yield \( \pm 6 \, dB \) because \( \omega_n \) and \( 2\omega_n \) are separated by one octave which is equivalent to \( \pm 3 \, dB \) per decade. Therefore, the corrections to be made are \( \pm 1 \, dB \) at half the corner frequency \( \omega_n/2 \), \( \pm 3 \, dB \) at the corner frequency \( \omega_n \), and \( \pm 1 \, dB \) at twice the corner frequency \( 2\omega_n \).

The corrected amplitude plots for a first order zero and first order pole are shown by solid lines in Figure 7.27.
The corrections for straight-line amplitude plots when we have complex poles and zeros require different type of correction because they depend on the damping coefficient $\zeta$. Let us refer to the plot of Figure 7.28.

We observe that as the damping coefficient $\zeta$ becomes smaller and smaller, larger and larger peaks in the amplitude occur in the vicinity of the corner frequency $\omega_n$. We also observe that when $\zeta \geq 0.707$, the amplitude at the corner frequency $\omega_n$ lies below the straight line approximation.

We can obtain a fairly accurate amplitude plot by computing the amplitude at four points near the corner frequency $\omega_n$ as shown in Figure 7.28.

Figure 7.27. Corrections for magnitude Bode plots
The amplitude plot of Figure 7.28 is for complex poles. In analogy with (7.30), i.e.,

\[ G(s) = s^2 + 2\zeta\omega_ns + \omega_n^2 \]

which was derived earlier for complex zeros, the transfer function for complex poles is

\[ G(s) = \frac{C}{s^2 + 2\zeta\omega_ns + \omega_n^2} \] (7.50)

where \( C \) is a constant.

Dividing each term of the denominator of (7.50) by \( \omega_n \) we get

\[ G(s) = \frac{C}{\omega_n^2 (s/\omega_n)^2 + 2\zeta(s/\omega_n) + 1} \]

and letting \( C/\omega_n^2 = K \) and \( s = j\omega \), we get

\[ G(j\omega) = \frac{K}{1 - (\omega/\omega_n)^2 + j2\zeta\omega/\omega_n} \] (7.51)

As before, we let \( \omega/\omega_n = u \). Then (7.51) becomes
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\[ G(j\omega) = \frac{K}{1 - \omega^2 + j2\zeta\omega} \]  
(7.52)

and in polar form,

\[ G(j\omega) = \frac{K}{|1 - \omega^2 + j2\zeta\omega|} \angle \theta \]  
(7.53)

The magnitude of (7.53) in dB is

\[ A_{dB} = 20\log|G(j\omega)| = 20\log K - 20\log|1 - \omega^2 + j2\zeta\omega| \]
\[ = 20\log K - 20\log \sqrt{(1 - \omega^2)^2 + 4\zeta^2 \omega^2} = 20\log K - 10\log[u^4 + 2\omega^2(2\zeta^2 - 1) + 1] \]  
(7.54)

and the phase is

\[ \theta(\omega) = -\tan^{-1} \frac{2\zeta\omega}{1 - \omega^2} \]  
(7.55)

In (7.54) the term \(20\log K\) is constant and thus the amplitude \(A_{dB}\), as a function of frequency, is dependent only the second term on the right side. Also, from this expression, we observe that as \(\omega \to 0\),

\[ -10\log[u^4 + 2\omega^2(2\zeta^2 - 1) + 1] \to 0 \]  
(7.56)

and as \(\omega \to \infty\),

\[ -10\log[u^4 + 2\omega^2(2\zeta^2 - 1) + 1] \to -40\log \omega \]  
(7.57)

We are now ready to compute the values of \(A_{dB}\) at points 1, 2, 3, and 4 of the plot of Figure 7.29. At point 1, the corner frequency \(\omega_n\) corresponds to \(\omega = 1\). Then, from (7.54)

\[ A_{dB}(\omega_n/2) = A_{dB}\left(\frac{1}{2}\right) = -10\log[u^4 + 2\omega^2(2\zeta^2 - 1) + 1]|_{\omega = 1/2} \]
\[ = -10\log\left[\frac{1}{16} + 2 \cdot \frac{1}{4}(2\zeta^2 - 1) + 1\right] = -10\log\left[\frac{1}{16} + \zeta^2 - \frac{1}{2} + 1\right] \]
\[ = -10\log(\zeta^2 + 0.5625) \]
\[ A_{dB}(\omega_n/2)|_{point \ 1} = -10\log(0.4^2 + 0.5625) = 1.41 \text{ dB} \]  
(7.58)
To find the amplitude at point 2, in (7.54) we let $K = 1$ and we form the magnitude in $dB$. Then,

$$A_{dB}|_{point\ 2} = 20\log\frac{1}{|1 - (\omega/\omega_n)^2 + j2\zeta\omega/\omega_n|}$$

We now recall that the logarithmic function is a monotonically increasing function and therefore (7.59) has a maximum when the absolute magnitude of this expression is maximum. Also, the square of the absolute magnitude is maximum when the absolute magnitude is maximum.

The square of the absolute magnitude is

$$\frac{1}{[1 - (\omega/\omega_n)^2]^2 + 4(\zeta\omega/\omega_n)^2}$$

---

**Figure 7.29.** Corrections for magnitude Bode plots with complex poles when $\zeta = 0.4$
or

\[
\frac{1}{1 - 2 \omega^2 / \omega_n^2 + \omega^4 / \omega_n^4 + 4 \zeta^2 \omega^2 / \omega_n^2}
\]  

(7.61)

To find the maximum, we take the derivative with respect to \(\omega\) and we set it equal to zero, that is,

\[
\frac{4\omega / \omega_n^2 - 4\omega^3 / \omega_n^4 - 8\zeta^2 \omega / \omega_n^2}{\left[1 - (\omega / \omega_n)^2\right]^2 + 4((\zeta \omega) / \omega_n)^2} = 0
\]

(7.62)

The expression of (7.62) will be zero when the numerator is set to zero, that is,

\[
(\omega / \omega_n^2)(4 - 4\omega^2 / \omega_n^2 - 8\zeta^2) = 0
\]

(7.63)

Of course, we require that the value of \(\omega\) must be a nonzero value. Then,

\[
4 - 4\omega^2 / \omega_n^2 - 8\zeta^2 = 0
\]

or

\[
(4\omega^2) / \omega_n^2 = 4 - 8\zeta^2
\]

from which

\[
\omega_{max} = \omega = \omega_n \sqrt{1 - 2\zeta^2}
\]

(7.64)

provided that \(1 - 2\zeta^2 > 0\) or \(\zeta < 1 / \sqrt{2}\) or \(\zeta < 0.707\).

The dB value of the amplitude at point 2 is found by substitution of (7.64) into (7.54), that is,

\[
A_{dB}(\omega_{max}) = -10 \log \left[ u^4 + 2u^2 (2\zeta^2 - 1) + 1 \right]_{u = \sqrt{1 - 2\zeta}}
\]

\[
= -10 \log \left[ (1 - 2\zeta^2)^2 + 2(1 - 2\zeta^2)(2\zeta^2 - 1) + 1 \right]
\]

\[
= -10 \log [4\zeta^2 (1 - \zeta^2)]
\]

(7.65)

and for \(\zeta = 0.4\)

\[
A_{dB}(\omega_{max}) = -10 \log (4 \times 0.4^2 (1 - 0.4^2)) = 2.69 \text{ dB}
\]

The dB value of the amplitude at point 3 is found by substitution of \(\omega = \omega_n = u = 1\) into (7.54). Then,
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\[ A_{dB}(\omega_n) = -10 \log \left[ u^4 + 2u^2(2\zeta^2 - 1) + 1 \right] \]

\[ = -10 \log \left[ 1 + 2(2\zeta^2 - 1) + 1 \right] \]

\[ = -10 \log [4\zeta^2] = -20 \log (2\zeta) \]  

(7.66)

and for \( \zeta = 0.4 \)

\[ A_{dB}(\omega_n) = -20 \log (2 \times 0.4) = 1.94 \text{ dB} \]

Finally, at point 4, the dB value of the amplitude crosses the 0 dB axis. Therefore, at this point we are interested not in \( A_{dB}(\omega_0 \text{ dB}) \) but in the location of \( \omega_0 \text{ dB} \) in relation to the corner frequency \( \omega_n \). at this point we must have from (7.57)

\[ 0 \text{ dB} = -10 \log \left[ u^4 + 2u^2(2\zeta^2 - 1) + 1 \right] \]

and since \( \log 1 = 0 \), it follows that

\[ u^4 + 2u^2(2\zeta^2 - 1) + 1 = 1 \]

\[ u^4 + 2u^2(2\zeta^2 - 1) = 0 \]

\[ u^2(u^2 + 2(2\zeta^2 - 1)) = 0 \]

or

\[ u^2 + 2(2\zeta^2 - 1) = 0 \]

Solving for \( u \) and making use of \( u = \omega / \omega_n \) we get

\[ \omega_0 \text{ dB} = \omega_n \sqrt{2(1 - 2\zeta^2)} \]

From (7.67),

\[ \omega_{max} = \omega_n \sqrt{1 - 2\zeta^2} \]

therefore, if we already know the frequency at which the dB amplitude is maximum, we can compute the frequency at point 4 from

\[ \omega_0 \text{ dB} = \sqrt{2}\omega_{max} \]  

(7.67)

Example 7.5

For the circuit of Figure 7.30
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Figure 7.30. Circuit for Example 7.5.

a. Compute the transfer function $G(s)$

b. Find the corner frequency $\omega_n$ from $G(s)$.

c. Compute the damping coefficient $\zeta$.

d. Construct a straight line approximation for the magnitude of the Bode plot.

e. Compute the amplitude in $dB$ at one-half the corner frequency $\omega_n/2$, at the frequency $\omega_{max}$ at which the amplitude reaches its maximum value, at the corner frequency $\omega_n$, and at the frequency $\omega_{dB}$ where the $dB$ amplitude is zero. Then, draw a smooth curve to connect these four points.

Solution:

a. We transform the given circuit to its equivalent in the $s$–domain shown in Figure 7.31 where $R = 1$, $Ls = 0.05s$, and $1/Cs = 125/s$.

Figure 7.31. Circuit for Example 7.5 in $s$–domain

and by the voltage division expression,

$$V_{out}(s) = \frac{25/s}{0.2 + 0.01s + 25/s} \cdot V_{in}(s)$$

Therefore, the transfer function is
Corrected Amplitude Plots

\[ G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{25}{0.01s^2 + 0.2s + 25} = \frac{2500}{s^2 + 20s + 2500} \]  \hspace{1cm} (7.68)

b. From (7.50)

\[ G(s) = \frac{K}{s^2 + 2\zeta\omega_n s + \omega_n^2} \]  \hspace{1cm} (7.69)

and from (7.68) and (7.69) \( \omega_n^2 = 2500 \) or

\[ \omega_n = 50 \text{ rad/s} \]  \hspace{1cm} (7.70)

c. From (7.68) and (7.69) \( 2\zeta\omega_n = 20 \). Then, the damping coefficient \( \zeta \) is

\[ \zeta = \frac{20}{2\omega_n} = \frac{20}{2 \times 50} = 0.2 \]  \hspace{1cm} (7.71)

d. For \( \omega < \omega_n \), the straight line approximation lies along the 0 dB axis, whereas for \( \omega > \omega_n \), the straight line approximation has a slope of \(-40 \text{ dB}\). The corner frequency \( \omega_n \) was found in part (b) to be \( 50 \text{ rad/s} \). The dB amplitude plot is shown in Figure 7.31.

e. From (7.61),

\[ A_{dB}(\omega_n/2) = -10\log(\zeta^2 + 0.5625) \]

where from (7.74) \( \zeta = 0.2 \) and thus \( \zeta^2 = 0.04 \). Then,

\[ A_{dB}(\omega_n/2) = -10\log(0.04 + 0.5625) = -10\log(0.6025) = 2.2 \text{ dB} \]

and this value is indicated as Point 1 on the plot of Figure 7.32.

Next, from (7.64)

\[ \omega_{max} = \omega_n\sqrt{1 - 2\zeta^2} \]

Then,

\[ \omega_{max} = 50\sqrt{1 - 2 \times 0.04} = 50\sqrt{0.92} = 47.96 \text{ rad/s} \]

Therefore, from (7.65)

\[ A_{dB}(\omega_{max}) = -10\log[4\zeta^2(1 - \zeta^2)] = -10\log[(0.16) \times (0.96)] = 8.14 \text{ dB} \]

and this value is indicated as Point 2 on the plot of Figure 7.32.

The dB amplitude at the corner frequency is found from (7.66), that is,
Then,

\[ A_{dB}(\omega_n) = -20 \log(2\zeta) \]

and this value is indicated as Point 3 on the plot of Figure 7.32.

Finally, the frequency at which the amplitude plot crosses the 0 dB axis is found from (7.67), that is,

\[ \omega_0 dB = \sqrt{2}\omega_{max} \]

or

\[ \omega_0 dB = \sqrt{2} \times 47.96 = 67.83 \text{ rad/s} \]

This frequency is indicated as Point 4 on the plot of Figure 7.32.

![Figure 7.32. Amplitude plot for Example 7.5](image)

The amplitude plot of Figure 7.32 reveals that the given circuit behaves as a low pass filter.
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Using the transfer function of (7.68) with MATLAB, we get the Bode magnitude plot shown in Figure 7.33.

\[ \text{num} = [0 \ 0 \ 2500]; \text{den} = [1 \ 20 \ 2500]; \text{sys} = \text{tf(num,den)}; \text{w} = \text{logspace(0,5,100)}; \text{bodemag(sys,w)} \]

Figure 7.33. Bode plot for Example 7.5 using MATLAB
7.8 Summary

- The decibel, denoted as dB, is a unit used to express the ratio between two amounts of power, generally \( P_{out}/P_{in} \). By definition, the number of dB is obtained from \( dB = 10 \log_{10}(P_{out}/P_{in}) \). It can also be used to express voltage and current ratios provided that the voltages and currents have identical impedances. Then, for voltages we use the expression \( dB_v = 20 \log_{10}(V_{out}/V_{in}) \), and for currents we use the expression \( dB_i = 20 \log_{10}(I_{out}/I_{in}) \).

- The bandwidth, denoted as \( BW \), is a term generally used with electronic amplifiers and filters. For low-pass filters the bandwidth is the band of frequencies from zero frequency to the cutoff frequency where the amplitude falls to 0.707 of its maximum value. For high-pass filters the bandwidth is the band of frequencies from 0.707 of maximum amplitude to infinite frequency. For amplifiers, band-pass, and band-elimination filters the bandwidth is the range of frequencies where the maximum amplitude falls to 0.707 of its maximum value on either side of the frequency response curve.

- If two frequencies \( \omega_1 \) and \( \omega_2 \) are such that \( \omega_2 = 2 \omega_1 \), we say that these frequencies are separated by one octave and if \( \omega_2 = 10 \omega_1 \), they are separated by one decade.

- Frequency response is a term used to express the response of an amplifier or filter to input sinusoids of different frequencies. The response of an amplifier or filter to a sinusoid of frequency \( \omega \) is completely described by the magnitude \( |G(j\omega)| \) and phase \( \angle G(j\omega) \) of the transfer function.

- Bode plots are frequency response diagrams of magnitude and phase versus frequency \( \omega \).

- In Bode plots the 3-dB frequencies, denoted as \( \omega_n \), are referred to as the corner frequencies.

- In Bode plots, the transfer function is expressed in linear factors of the form \( j\omega + z_i \) for the zero (numerator) linear factors and \( j\omega + p_i \) for the pole linear factors. When quadratic factors with complex roots occur in addition to the linear factors, these quadratic factors are expressed in the form \( (j\omega)^2 + 2\zeta \omega_n \omega + \omega_n^2 \).

- In magnitude Bode plots with quadratic factors the difference between the asymptotic plot and the actual curves depends on the value of the damping factor \( \zeta \). But regardless of the value of \( \zeta \), the actual curve approaches the asymptotes at both low and high frequencies.

- In Bode plots the corner frequencies \( \omega_n \) are easily identified by expressing the linear terms as \( z_i(j\omega/z_i + 1) \) and \( p_i(j\omega/p_i + 1) \) for the zeros and poles respectively. For quadratic factor the corner frequency \( \omega_n \) appears in the expression \( (j\omega)^2 + 2\zeta \omega_n \omega + \omega_n^2 \) or \( (j\omega/\omega_n)^2 + 2\zeta \omega/\omega_n + 1 \).
In both the magnitude and phase Bode plots the frequency (abscissa) scale is logarithmic. The ordinate in the magnitude plot is expressed in $\text{dB}$ and in the phase plot is expressed in degrees.

In magnitude Bode plots, the asymptotes corresponding to the linear terms of the form $(j\omega/z_i + 1)$ and $(j\omega/p_i + 1)$ have a slope $\pm 20 \text{dB/decade}$ where the positive slope applies to zero (numerator) linear factors, and the negative slope applies to pole (denominator) linear factors.

In magnitude Bode plots, the asymptotes corresponding to the quadratic terms of the form $(j\omega/\omega_n)^2 + j2\zeta\omega/\omega_n + 1$ have a slope $\pm 40 \text{dB/decade}$ where the positive slope applies to zero (numerator) quadratic factors, and the negative slope applies to pole (denominator) quadratic factors.

In phase Bode plots with linear factors, for frequencies less than one tenth the corner frequency we assume that the phase angle is zero. At the corner frequency the phase angle is $\pm 45^\circ$. For frequencies ten times or greater than the corner frequency, the phase angle is approximately $\pm 90^\circ$ where the positive angle applies to zero (numerator) linear factors, and the negative angle applies to pole (denominator) linear factors.

In phase Bode plots with quadratic factors, the phase angle is zero for frequencies less than one tenth the corner frequency. At the corner frequency the phase angle is $\pm 90^\circ$. For frequencies ten times or greater than the corner frequency, the phase angle is approximately $\pm 180^\circ$ where the positive angle applies to zero (numerator) quadratic factors, and the negative angle applies to pole (denominator) quadratic factors.

Bode plots can be easily constructed and verified with the MATLAB function $\text{bode(sys)}$ function. With this function, the frequency range and number of points are chosen automatically. The function $\text{bode(sys),\{wmin,wmax\}}$ draws the Bode plot for frequencies between $\text{wmin}$ and $\text{wmax}$ (in radian/second) and the function $\text{bode(sys,w)}$ uses the user-supplied vector $\text{w}$ of frequencies, in radian/second, at which the Bode response is to be evaluated. To generate logarithmically spaced frequency vectors, we use the command $\text{logspace(first_exponent,last_exponent, number_of_values)}$. 
7.9 Exercises

1. For the transfer function

\[ G(s) = \frac{10^5(s + 5)}{(s + 100)(s + 5000)} \]

a. Draw the magnitude Bode plot and find the approximate maximum value of \(|G(j\omega)|\) in dB.

b. Find the value of \(\omega\) where \(|G(j\omega)| = 1\) for \(\omega > 5\) r/s

c. Check your plot with the plot generated with MATLAB.

2. For the transfer function of Exercise 1

a. Draw a Bode plot for the phase angle and find the approximate phase angle at \(\omega = 30\) r/s, \(\omega = 50\) r/s, \(\omega = 100\) r/s, and \(\omega = 5000\) r/s.

b. Compute the actual values of the phase angle at the frequencies specified in (a).

c. Check your magnitude plot of Exercise 1 and the phase plot of this exercise with the plots generated with MATLAB.

3. For the circuit of Figure 7.34

a. Compute the transfer function.

b. Draw the Bode amplitude plot for \(20\log|G(j\omega)|\)

c. From the plot of part (b) determine the type of filter represented by this circuit and estimate the cutoff frequency.

d. Compute the actual cutoff frequency of this filter.

e. Draw a straight line phase angle plot of \(G(j\omega)\).

f. Determine the value of \(\theta(\omega)\) at the cutoff frequency from the plot of part (c).

g. Compute the actual value of \(\theta(\omega)\) at the cutoff frequency.

---

**Figure 7.34. Circuit for Exercise 3**
7.10 Answers to Exercises

1. a. 

\[
G(j\omega) = \frac{10^5 (j\omega + 5)}{(j\omega + 100)(j\omega + 5000)} = \frac{10^5 \times (1 + j\omega/5)}{100 \times (1 + j\omega/100) \times 5000 \times (1 + j\omega/5000)} = \frac{(1 + j\omega/5)}{(1 + j\omega/100) \cdot (1 + j\omega/5000)}
\]

\[
20\log|G(j\omega)| = 20\log|1 + j\omega/5| - 20\log|1 + j\omega/100| - 20\log|1 + j\omega/5000|
\]

The corner frequencies are at \( \omega = 5 \, \text{r/s} \), \( \omega = 100 \, \text{r/s} \), and \( \omega = 5000 \, \text{r/s} \). The asymptotes are shown as solid lines.

From this plot we observe that \( 20\log|G(j\omega)|_{\text{max}} \approx 26 \, \text{dB} \) for the interval \( 10^2 \leq \omega \leq 5 \times 10^3 \).
b. By inspection, $20\log|G(j\omega)| = 0 \text{ dB}$ at $\omega = 9.85 \times 10^4 \text{ r/s}$

2. From the solution of Exercise 1

$$G(j\omega) = \frac{(1 + j\omega/5)}{(1 + j\omega/100) \cdot (1 + j\omega/5000)}$$

and in magnitude-phase form

$$G(j\omega) = \frac{|(1 + j\omega/5)|}{|(1 + j\omega/100)| \cdot |(1 + j\omega/5000)|} \angle (\alpha - \beta - \gamma)$$

that is, $\theta(\omega) = \alpha - \beta - \gamma$ where $\alpha = \tan^{-1}\omega/5$, $\beta = -\tan^{-1}\omega/100$, and $\gamma = -\tan^{-1}\omega/5000$

The corner frequencies are at $\omega = 5 \text{ r/s}$, $\omega = 100 \text{ r/s}$, and $\omega = 5000 \text{ r/s}$ where at those frequencies $\alpha = 45^\circ$, $\beta = -45^\circ$, and $\gamma = -45^\circ$ respectively. The asymptotes are shown as solid lines.

From the phase plot we observe that $\theta(30 \text{ r/s}) \approx 60^\circ$, $\theta(50 \text{ r/s}) \approx 53^\circ$, $\theta(100 \text{ r/s}) \approx 38^\circ$, and $\theta(5000 \text{ r/s}) \approx -39^\circ$
b. We use MATLAB for the computations.

\[
\theta_{g30} = \frac{(1 + 30j/5) \cdot (1 + 30j/5000)}{(1 + 30j/100) \cdot (1 + 30j/5000)}; \ldots \\
\theta_{g50} = \frac{(1 + 50j/5) \cdot (1 + 50j/5000)}{(1 + 50j/100) \cdot (1 + 50j/5000)}; \ldots \\
\theta_{g100} = \frac{(1 + 100j/5) \cdot (1 + 100j/5000)}{(1 + 100j/100) \cdot (1 + 100j/5000)}; \ldots \\
\theta_{g5000} = \frac{(1 + 5000j/5) \cdot (1 + 5000j/5000)}{(1 + 5000j/100) \cdot (1 + 5000j/5000)}; \ldots \\
\]

printf(' 
');
fprintf('theta30r = %5.2f deg. 	', angle(theta_g30)*180/pi);
fprintf('theta50r = %5.2f deg. ', angle(theta_g50)*180/pi);
fprintf(' 
');
fprintf('theta100r = %5.2f deg. 	', angle(theta_g100)*180/pi);
fprintf('theta5000r = %5.2f deg. ', angle(theta_g5000)*180/pi);
fprintf(' 
')

theta30r = 63.49 deg. theta50r = 57.15 deg.
theta100r = 40.99 deg. theta5000r = -43.91 deg.

Thus, the actual values are

\[
\angle G(j30) = \angle \frac{(1 + j30/5)}{(1 + j30/100) \cdot (1 + j30/5000)} = 63.49^\circ \\
\angle G(j50) = \angle \frac{(1 + j50/5)}{(1 + j50/100) \cdot (1 + j50/5000)} = 57.15^\circ \\
\angle G(j100) = \angle \frac{(1 + j100/5)}{(1 + j100/100) \cdot (1 + j100/5000)} = 40.99^\circ \\
\angle G(j5000) = \angle \frac{(1 + j5000/5)}{(1 + j5000/100) \cdot (1 + j5000/5000)} = -43.91^\circ \\
\]

c. The Bode plot generated with MATLAB is shown below.

```matlab
syms s; expand((s+100)*(s+5000))
ans =
s^2+5100*s+500000
num=[0 10^5 5*10^5]; den=[1 5.1*10^3 5*10^5]; w=logspace(0,10^4); bode(num,den,w)
```

\[
\]
3. a. The equivalent \( s \)-domain circuit is shown below.

\[
V_{in}(s) \quad 0.25s \quad 1 \quad 25/s \quad V_{out}(s)
\]

By the voltage division expression

\[ V_{out}(s) = \frac{1 + 25/s}{0.25s + 1 + 25/s} \cdot V_{in}(s) \]

and

\[ G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{s + 25}{0.25s^2 + s + 25} = \frac{4(s + 25)}{s^2 + 4s + 100} \] (1)

b. From (1) with \( s = j\omega \)

\[ G(j\omega) = \frac{4(j\omega + 25)}{-\omega^2 + 4j\omega + 100} \] (2)

From (7.53)
\[ G(s) = \frac{C}{s^2 + 2\zeta\omega_n s + \omega_n^2} \quad (3) \]

and from (1) and (3) \( \omega_n^2 = 100 \), \( \omega_n = 10 \), and \( 2\zeta\omega_n = 4 \), \( \zeta = 0.2 \)

Following the procedure of page 7-26 we let \( u = \omega/\omega_n = \omega/10 \). The numerator of (2) is a linear factor and thus we express it as \( 100(1 + j\omega/25) \). Then (2) is written as

\[
G(j\omega) = \frac{100(1 + j\omega/25)}{100(-\omega^2/100 + 4j\omega/100 + 100/100)} = \frac{(1 + j\omega/25)}{1 - (\omega/10)^2 + j0.4/10}
\]

or

\[
G(j\omega) = \frac{|1 + j\omega/25|\angle\theta}{|1 - (\omega/10)^2 + j0.4\omega/10|\angle\phi} \quad (4)
\]

The amplitude of \( G(j\omega) \) in dB is

\[
20\log|G(j\omega)| = 20\log|1 + j\omega/25| - 20\log|1 - (\omega/10)^2 + j0.4\omega/10| \quad (5)
\]

The asymptote of the first term on the right side of (5) has a corner frequency of \( 25 \text{ r/s} \) and rises with slope of \( 20 \text{ dB/decade} \). The second term has a corner frequency of \( 10 \text{ r/s} \) and rises with slope of \( -40 \text{ dB/decade} \). The amplitude plot is shown below.

c. The plot above indicates that the circuit is a low-pass filter and the 3 dB cutoff frequency \( \omega_c \) occurs at approximately \( 13 \text{ r/s} \).
d. The actual cutoff frequency occurs where

\[ |G(j\omega_c)| = \frac{|G(j\omega)|_{max}}{\sqrt{2}} = \frac{1}{\sqrt{2}} = 0.70 \]

At this frequency (2) is written as

\[ G(j\omega_c) = \frac{100 + 4j\omega_c}{(100-\omega_c^2) + 4j\omega} \]

and considering its magnitude we get

\[ \frac{\sqrt{100^2 + (4\omega_c)^2}}{\sqrt{(100-\omega_c^2)^2 + (4\omega_c)^2}} = \frac{1}{\sqrt{2}} \]

\[ 2[100^2 + (4\omega_c)^2] = (100-\omega_c^2)^2 + (4\omega_c)^2 \]

\[ 20000 + 32\omega_c^2 = 10000 - 200\omega_c^2 + \omega_c^4 + 16\omega_c^2 \]

\[ \omega_c^4 - 216\omega_c^2 - 10000 = 0 \]

We will use MATLAB to find the four roots of this equation.

```matlab
syms w; solve(w^4 - 216*w^2 - 10000)
```

```matlab
ans =
[ 2*(27+1354^(1/2))^(1/2)]
[ 2*(27-1354^(1/2))^(1/2)]
```

\[ w1 = 2*(27+1354^(1/2))^(1/2) \]

\[ w1 = 15.9746 \]

\[ w2 = -2*(27+1354^(1/2))^(1/2) \]

\[ w2 = -15.9746 \]

\[ w3 = 2*(27-1354^(1/2))^(1/2) \]

\[ w3 = 0.0000 + 6.2599i \]

\[ w4 = -2*(27-1354^(1/2))^(1/2) \]

\[ w4 = -0.0000 - 6.2599i \]
From these four roots we accept only the first, that is, $\omega_c \approx 16 \text{ r/s}$

e. From (4)

$$\theta = \tan^{-1}(\omega/25)$$

and

$$\phi = \frac{0.4\omega/10}{1 - (\omega/10)^2}$$

For a first order zero or pole not at the origin, the straight line phase angle plot approximations are as follows:

I. For frequencies less than one tenth the corner frequency we assume that the phase angle is zero. For this exercise the corner frequency of $\theta(\omega)$ is $\omega_n = 25 \text{ r/s}$ and thus for $1 \leq \omega \leq 2.5 \text{ r/s}$ the phase angle is zero as shown on the Bode plot below.

II. For frequencies ten times or greater than the corner frequency, the phase angle is approximately $\pm 90^\circ$. The numerator phase angle $\theta(\omega)$ is zero at one tenth the corner frequency, it is $45^\circ$ at the corner frequency, and $90^\circ$ for frequencies ten times or greater the corner frequency. For this exercise, in the interval $2.5 \leq \omega \leq 250 \text{ r/s}$ the phase angle is zero at $2.5 \text{ r/s}$ and rises to $90^\circ$ at $250 \text{ r/s}$. 

\[\begin{array}{c}
\text{Phase angle (Degrees)} \\
\hline
\theta(\omega) \\
\varphi(\omega_n) = 10 \text{ r/s} \\
\angle G(j\omega) \\
\phi(\omega) \\
\hline
\end{array}\]
As shown in Figure 7.20, for complex poles the phase angle is zero at zero frequency, 
\[ -90^\circ \] at the corner frequency and approaches \[ -180^\circ \] as the frequency becomes large. The
phase angle asymptotes are shown on the plot of the previous page.

f. From the plot of the previous page we observe that the phase angle at the cutoff frequency is
approximately \[ -63^\circ \]

g. The exact phase angle at the cutoff frequency \( \omega_c = 16 \, \text{rad/s} \) is found from (1) with \( s = j16 \).

\[
G(j16) = \frac{4(j16 + 25)}{(j16)^2 + 4(j16) + 100}
\]

We need not simplify this expression since we can use MATLAB.

\[
g16 = \frac{(64j+100)/((16j)^2+64j+100)}{\text{angle}(g16)*180/\pi}
\]

\[
\text{ans } = -125.0746
\]

This value is about twice as that we observed from the asymptotic plot of the previous page. Errors such as this occur because of the high non-linearity between frequency intervals. Therefore, we should use the straight line asymptotes only to observe the shape of the phase angle. It is best to use MATLAB as shown below.

\[
\text{num}=[0 \ 4 \ 100]; \text{den}=[1 \ 4 \ 100]; \text{w} = \text{logspace}(0,2,1000); \text{bode(num,den,w)}
\]
This chapter begins with the interactions between electric circuits and changing magnetic fields. It defines self and mutual inductances, flux linkages, induced voltages, the dot convention, Lenz’s law, and magnetic coupling. It concludes with a detailed discussion on transformers.

### 8.1 Self-Inductance

About 1830, Joseph Henry, while working at the university which is now known as Princeton, found that electric current flowing in a circuit has a property analogous to mechanical momentum which is a measure of the motion of a body and it is equal to the product of its mass and velocity, i.e., $Mv$. In electric circuits this property is sometimes referred to as the **electrokinetic momentum** and it is equal to the product of $Li$ where $i$ is the current analogous to velocity and the **self-inductance** $L$ is analogous to the mass $M$. About the same time, Michael Faraday visualized this property in a magnetic field in space around a current carrying conductor. This electrokinetic momentum is denoted by the symbol $\lambda$, that is,

$$\lambda = Li \quad (8.1)$$

Newton’s second law states that the force necessary to change the velocity of a body with mass $M$ is equal to the rate of change of the momentum, i.e.,

$$F = \frac{d}{dt}(Mv) = M\frac{dv}{dt} = Ma \quad (8.2)$$

where $a$ is the acceleration. The analogous electrical relation says that the voltage $v$ necessary to produce a change of current in an inductive circuit is equal to the rate of change of electrokinetic momentum, i.e,

$$v = \frac{d}{dt}(Li) = L\frac{di}{dt} \quad (8.3)$$

### 8.2 The Nature of Inductance

Inductance is associated with the magnetic field which is always present when there is an electric current. Thus when current flows in an electric circuit, the conductors (wires) connecting the devices in the circuit are surrounded by a magnetic field. Figure 8.1 shows a simple loop of wire and its magnetic field which is represented by the small loops. The direction of the magnetic field (not shown) can be
determined by the left-hand rule if conventional current flow is assumed, or by the right-hand rule if electron current flow is assumed. The magnetic field loops are circular in form and are called lines of magnetic flux. The unit of magnetic flux is the weber (Wb).

![Figure 8.1. Magnetic field around a current carrying wire](image1)

In a loosely wound coil of wire such as the one shown in Figure 8.2, the current through the wound coil produces a denser magnetic field and many of the magnetic lines link the coil several times.

![Figure 8.2. Magnetic field around a current carrying wound coil](image2)

The magnetic flux is denoted as \( \phi \) and, if there are \( N \) turns and we assume that the flux \( \phi \) passes through each turn, the total flux denoted as \( \lambda \) is called flux linkage. Then,

\[
\lambda = N \phi \tag{8.4}
\]

By definition, a linear inductor one in which the flux linkage is proportional to the current through it, that is,

\[
\lambda = Li \tag{8.5}
\]

where the constant of proportionality \( L \) is called inductance in webers per ampere.

We now recall Faraday’s law of electromagnetic induction which states that

\[
v = \frac{d\lambda}{dt} \tag{8.6}
\]

and from (8.3) and (8.5),

\[
v = L \frac{di}{dt} \tag{8.7}
\]
8.3 Lenz’s Law

Heinrich F. E. Lenz was a German scientist who, without knowledge of the work of Faraday and Henry, duplicated many of their discoveries nearly simultaneously. The law which goes by his name, is a useful rule for predicting the direction of an induced current. Lenz’s law states that:

Whenever there is a change in the amount of magnetic flux linking an electric circuit, an induced voltage of value directly proportional to the time rate of change of flux linkages is set up tending to produce a current in such a direction as to oppose the change in flux.

To understand Lenz’s law, let us consider the transformer shown in Figure 8.3.

![Figure 8.3. Basic transformer construction](image)

Here, we assume that the current in the primary winding has the direction shown and it produces the flux $\phi$ in the direction shown in Figure 8.3 by the arrow below the dotted line. Suppose that this flux is decreasing. Then in the secondary winding there will be a voltage induced whose current will be in a direction to increase the flux. In other words, the current produced by the induced voltage will tend to prevent any decrease in flux. Conversely, if the flux produced by the primary winding in increasing, the induced voltage in the secondary will produce a current in a direction which will oppose an increase in flux.

8.4 Mutually Coupled Coils

Consider the inductor (coil) shown in Figure 8.4. There are many magnetic lines of flux linking the coil $L_1$ with $N_1$ turns but for simplicity, only two are shown in Figure 8.4. The current $i_1$ produces a magnetic flux $\phi_{11}$. Then by (8.4) and (8.5)

$$\lambda_1 = N_1 \phi_{11} = L_1 i_1$$

(8.8)

and by Faraday’s law of (8.6), in terms of the self-inductance $L_1$,

$$v_1 = \frac{d\lambda_1}{dt} = N_1 \frac{d\phi_{11}}{dt} = L_1 \frac{di_1}{dt}$$

(8.9)
Next, suppose another coil $L_2$ with $N_2$ turns is brought near the vicinity of coil $L_1$ and some lines of flux are also linking coil $L_2$ as shown in Figure 8.5.

It is convenient to express the flux $\phi_{11}$ as the sum of two fluxes $\phi_{L1}$ and $\phi_{21}$, that is,

$$\phi_{11} = \phi_{L1} + \phi_{21}$$

(8.10)

where the linkage flux $\phi_{L1}$ is the flux which links coil $L_1$ only and not coil $L_2$, and the mutual flux $\phi_{21}$ is the flux which links both coils $L_1$ and $L_2$. We have assumed that the linkage and mutual fluxes $\phi_{L1}$ and $\phi_{21}$ link all turns of coil $L_1$ and the mutual flux $\phi_{21}$ links all turns of coil $L_2$.

The arrangement above forms an elementary transformer where coil $L_1$ is called the primary winding and coil $L_2$ the secondary winding.
Mutually Coupled Coils

In a linear transformer the mutual flux $\varphi_{21}$ is proportional to the primary winding current $i_1$ and since there is no current in the secondary winding, the flux linkage in the secondary winding is by (8.8),

$$\lambda_2 = N_2 \varphi_{21} = M_{21} i_1$$  \hspace{1cm} (8.11)

where $M_{21}$ is the mutual inductance (in Henries) and thus the open-circuit secondary winding voltage $v_2$ is

$$v_2 = \frac{d\lambda_2}{dt} = N_2 \frac{d\varphi_{21}}{dt} = M_{21} \frac{di_1}{dt}$$  \hspace{1cm} (8.12)

In summary, when there is no current in the secondary winding the voltages are

$$v_1 = L_1 \frac{di_1}{dt} \text{ and } v_2 = M_{21} \frac{di_1}{dt}$$  \hspace{1cm} (8.13)

if $i_1 \neq 0$ and $i_2 = 0$

Next, we will consider the case where there is a voltage in the secondary winding producing current $i_2$ which in turn produces flux $\varphi_{22}$ as shown in Figure 8.6.

![Figure 8.6. Flux in secondary winding](image)

Then in analogy with (8.8) and (8.9)

$$\lambda_2 = N_2 \varphi_{22} = L_2 i_2$$  \hspace{1cm} (8.14)

and by Faraday’s law in terms of the self-inductance $L_2$

$$v_2 = \frac{d\lambda_2}{dt} = N_2 \frac{d\varphi_{22}}{dt} = L_2 \frac{di_2}{dt}$$  \hspace{1cm} (8.15)
If another coil $L_1$ with $N_1$ turns is brought near the vicinity of coil $L_2$, some lines of flux are also linking coil $L_1$ as shown in Figure 8.7.

![Figure 8.7. Lines of flux linking open primary coil](image)

Following the same procedure as above we express the flux $\varphi_{22}$ as the sum of two fluxes $\varphi_{L2}$ and $\varphi_{I2}$, that is,

$$\varphi_{22} = \varphi_{L2} + \varphi_{I2}$$  \hspace{1cm} (8.16)

where the **linkage flux** $\varphi_{L2}$ is the flux which links coil $L_2$ only and not coil $L_1$, and the **mutual flux** $\varphi_{I2}$ is the flux which links both coils $L_2$ and $L_1$. As before, we have assumed that the linkage and mutual fluxes link all turns of coil $L_2$ and the mutual flux links all turns of coil $L_1$.

Since there is no current in the primary winding, the flux linkage in the primary winding is

$$\lambda_1 = N_1 \varphi_{I2} = M_{12}i_2$$  \hspace{1cm} (8.17)

where $M_{12}$ is the **mutual inductance** (in Henries) and thus the open-circuit primary winding voltage $v_1$ is

$$v_1 = \frac{d\lambda_1}{dt} = N_1 \frac{d\varphi_{I2}}{dt} = M_{12} \frac{di_2}{dt}$$  \hspace{1cm} (8.18)

In summary, when there is no current in the primary winding, the voltages are

$$v_2 = L_2 \frac{di_2}{dt} \quad \text{and} \quad v_1 = M_{12} \frac{di_2}{dt}$$  \hspace{1cm} (8.19)

\[\text{if} \quad i_1 = 0 \quad \text{and} \quad i_2 \neq 0\]
We will see later that

\[ M_{12} = M_{21} = M \]  

(8.20)

The last possible arrangement is shown in Figure 8.8 where \( i_1 \neq 0 \) and also \( i_2 \neq 0 \).

![Flux linkages when both primary and secondary currents are present](image)

**Figure 8.8. Flux linkages when both primary and secondary currents are present**

The total flux \( \varphi_1 \) linking coil \( L_1 \) is

\[ \varphi_1 = \varphi_{L1} + \varphi_{I2} = \varphi_{II} + \varphi_{I2} \]  

(8.21)

and the total flux \( \varphi_2 \) linking coil \( L_2 \) is

\[ \varphi_2 = \varphi_{L2} + \varphi_{I2} = \varphi_{II} + \varphi_{22} \]  

(8.22)

and since \( \lambda = N \varphi \), we express (8.21) and (8.22) as

\[ \lambda_1 = N_1 \varphi_{II} + N_1 \varphi_{I2} \]  

(8.23)

and

\[ \lambda_2 = N_2 \varphi_{II} + N_2 \varphi_{I2} \]  

(8.24)

Differentiating (8.23) and (8.24) and using (8.13), (8.14), (8.19) and (8.20) we get:

\[ v_1 = L_1 \frac{di_1}{dt} + M \frac{di_2}{dt} \]  

(8.25)

\[ v_2 = M \frac{di_1}{dt} + L_2 \frac{di_2}{dt} \]
In (8.25) the voltage terms
\[ L_1 \frac{di_1}{dt} \quad \text{and} \quad L_2 \frac{di_2}{dt} \]
are referred to as *self-induced voltages* and the terms
\[ M \frac{di_1}{dt} \quad \text{and} \quad M \frac{di_2}{dt} \]
are referred to as *mutual voltages*.

In our previous studies we used the passive sign convention as a basis to denote the polarity (+) and (−) of voltages and powers. While this convention can be used with the self-induced voltages, it cannot be used with mutual voltages because there are four terminals involved. Instead, the polarity of the mutual voltages is denoted by the *dot convention*. To understand this convention, we first consider the transformer circuit designations shown in Figures 8.9(a) and 8.9(b) where the dots are placed on the upper terminals and the lower terminals respectively.

**Figure 8.9. Arrangements where the mutual voltage has a positive sign**

These designations indicate the condition that a current \( i \) entering the dotted (undotted) terminal of one coil induce a voltage across the other coil with positive polarity at the dotted (undotted) terminal of the other coil. Thus, the mutual voltage term has a positive sign. Following the same rule we see that in the circuits of Figure 8.10 (a) and 8.10(b) the mutual voltage has a negative sign.

**Example 8.1**

For the circuit of Figure 8.11 find \( v_1 \) and \( v_2 \) if

a. \( i_1 = 50 \, mA \) and \( i_2 = 25 \, mA \)

b. \( i_1 = 0 \) and \( i_2 = 20 \sin 377t \, mA \)

c. \( i_1 = 15 \cos 377t \, mA \) and \( i_2 = 40 \sin(377t + 60^\circ) \, mA \)
Solution:

a. Since both currents $i_1$ and $i_2$ are constants, their derivatives are zero, i.e.,

$$\frac{di_1}{dt} = \frac{di_2}{dt} = 0$$

and thus

$$v_1 = v_2 = 0$$

b. The dot convention in the circuit of Figure 8.11 shows that the mutual voltage terms are positive and thus

$$v_1 = L_1 \frac{di_1}{dt} + M \frac{di_2}{dt} = 0.05 \times 0 + 20 \times 10^{-3} \times 20 \times 377 \times \cos 377t$$

$$= 150.8 \cos 377t \text{ mV}$$

$$v_2 = M \frac{di_1}{dt} + L_2 \frac{di_2}{dt} = 20 \times 10^{-3} \times 0 + 0.05 \times 20 \times 377 \times \cos 377t$$

$$= 377 \cos 377t \text{ mV}$$
Example 8.2

For the circuit of Figure 8.12 find the open-circuit voltage $v_2$ for $t > 0$ given that $i_1(0^-) = 0$.

![Figure 8.12. Circuit for Example 8.2](image)

Solution:

For $t > 0$

\[
\frac{di_1}{dt} + Ri_1 = 24
\]

\[
0.05 \frac{di_1}{dt} + 5i_1 = 24
\]

\[
\frac{di_1}{dt} + 100i_1 = 480
\]

Now,

\[
i_1 = i_f + i_n
\]

where $i_f$ is the forced response component of $i_1$ and it is obtained from

\[
i_f = \frac{24}{5} = 4.8 \text{ A}
\]
and $i_n$ is the natural response component of $i_I$ and it is obtained from

$$i_n = Ae^{-Rt/L} = Ae^{-100t}$$

Then,

$$i_I = i_f + i_n = 4.8 + Ae^{-100t}$$

and with the initial condition

$$i_I(0^+) = i_I(0^-) = 0 = 4.8 + Ae^0$$

we get $A = -4.8$

Therefore,

$$i_I = i_f + i_n = 4.8 - 4.8e^{-100t}$$

and in accordance with the dot convention,

$$v_2 = -M \frac{di_I}{dt} = -0.02(480e^{-100t}) = -9.6e^{-100t}$$

### 8.5 Establishing Polarity Markings

In our previous discussion and in Examples 8.1 and 8.2, the polarity markings (dots) were given. There are cases, however, when these are not known. The following method is generally used to establish the polarity marking in accordance with the dot convention.

Consider the transformer and its circuit symbol shown in Figure 8.13.

![Figure 8.13. Establishing polarity markings](image)

We recall that the direction of the flux $\phi$ can be found by the right-hand rule which states that if the fingers of the right hand encircle a winding in the direction of the current, the thumb indicates the direction of the flux. Let us place a dot at the upper end of $L_I$ and assume that the current $i_I$ enters the top end thereby producing a flux in the clockwise direction shown. Next, we want the current in
Chapter 8  Self and Mutual Inductances - Transformers

\( L_2 \) to enter the end which will produce a flux in the same direction, in this case, clockwise. This will be accomplished if the current \( i_2 \) in \( L_2 \) enters the lower end as shown and thus we place a dot at that end.

**Example 8.3**

For the transformer shown in Figure 8.14, find \( v_1 \) and \( v_2 \).

\[
\begin{align*}
    i_1 &= 2 \sin 377t \ A \\
    i_2 &= -5 \cos 377t \ A \\
    v_1 &= L_1 \frac{di_1}{dt} - M \frac{di_2}{dt} = 2262 \cos 377t - 3770 \sin 377t \ V \\
    v_2 &= -M \frac{di_1}{dt} + L_2 \frac{di_2}{dt} = -1508 \cos 377t + 7540 \sin 377t \ V 
\end{align*}
\]
### Example 8.4

For the circuit below, find the voltage ratio $|V_2/V_1|$. *

![Circuit for Example 8.4](image)

**Solution:**

The dots are given to us as shown. Now, we arbitrarily assign currents $I_1$ and $I_2$ as shown in Figure 8.17 and we write mesh equations for each mesh.

![Mesh currents for the circuit of Example 8.4](image)

With this current assignments $I_2$ leaves the dotted terminal of the right mesh and therefore the mutual voltage has a negative sign. Then,

Mesh 1:

$$R_1I_1 + j\omega L_1 I_1 - j\omega MI_2 = V_{in}$$

or

$$(0.5 + j18.85)I_1 - j18.85I_2 = 120\angle0^\circ$$ \tag{8.26}

* Henceforth we will be using bolded capital letters to denote phasor quantities.
Mesh 2:

\[-j\omega M I_1 + j\omega L_2 I_2 + R_{LOAD} I_2 = 0\]

or

\[-j18.85I_1 + (1000 + j37.7)I_2 = 0\]  \hspace{1cm} (8.27)

We will find the ratio \(|V_2/V_1|\) using the MATLAB code below where \(V_1 = j\omega L_1 I_1 = j18.85I_1\) and \(Z = [0.5 + 18.85j, -18.85j; -18.85j, 500 + 37.7j]\); \(V = [120 \, 0]'\); \(I = Z\backslash V;\)

\[
\text{fprintf('V1 = %7.3f V 	', abs(18.85j*I(1)));
fprintf('V2 = %7.3f V 	', abs(500*I(2)));
fprintf('Ratio V2/V1 = %7.3f 	',abs((500*I(2))/(18.85j*I(1))))
\]

\[
V1 = 120.093 \, \text{V} \quad V2 = 119.753 \, \text{V} \quad \text{Ratio V2/V1 = 0.997}
\]

That is,

\[
\frac{|V_2|}{|V_1|} = \frac{119.75}{120.09} = 0.997
\]  \hspace{1cm} (8.28)

and thus the magnitude of \(V_{LOAD} = V_2\) is practically the same as the magnitude of \(V_{in}\). However, we suspect that \(V_{LOAD}\) will be out of phase with \(V_{in}\). We can find the phase of \(V_{LOAD}\) by adding the following statement to the MATLAB code above.

\[
\text{fprintf('Phase V2= %6.2f deg', angle(500*I(2))*180/pi)}
\]

Phase \(V2 = -0.64\, \text{deg}\)

This is a very small phase difference from the phase of \(V_{in}\) and thus we see that both the magnitude and phase of \(V_{LOAD}\) are essentially the same as that of \(V_{in}\).

If we increase the load resistance \(R_{LOAD}\) to 1 K\(\Omega\) we will find that again the magnitude and phase of \(V_{LOAD}\) are essentially the same as that of \(V_{in}\). Therefore, the transformer of this example is an isolation transformer, that is, it isolates the load from the source and the value of \(V_{in}\) appears across the load even though the load changes. An isolation transformer is also referred to as a 1:1 transformer.

If in a transformer the secondary winding voltage is considerably higher than the input voltage, the transformer is referred to as a step-up transformer. Conversely, if the secondary winding voltage is considerably lower than the input voltage, the transformer is referred to as a step-down transformer.

### 8.6 Energy Stored in a Pair of Mutually Coupled Inductors

We know that the energy stored in an inductor is

\[
W(t) = \frac{1}{2}L_i^2(t)\]  \hspace{1cm} (8.29)
Energy Stored in a Pair of Mutually Coupled Inductors

In the transformer circuits shown in Figure 8.18, the stored energy is the sum of the energies supplied to the primary and secondary terminals. From (8.25),

\[
\begin{align*}
\text{Figure 8.18. Transformer circuits for computation of the energy} \\
\end{align*}
\]

\[v_1 = L_1 \frac{di_1}{dt} + M \frac{di_2}{dt}\]

\[v_2 = M \frac{di_1}{dt} + L_2 \frac{di_2}{dt}\]

and after replacing \(M\) with \(M_{12}\) and \(M_{21}\) in the appropriate terms, the instantaneous power delivered to these terminals are:

\[
\begin{align*}
p_1 &= v_1 i_1 = \left( L_1 \frac{di_1}{dt} + M_{12} \frac{di_2}{dt} \right) i_1 \\
p_2 &= v_2 i_2 = \left( M_{21} \frac{di_1}{dt} + L_2 \frac{di_2}{dt} \right) i_2
\end{align*}
\]

Now, let us suppose that at some reference time \(t_0\), both currents \(i_1\) and \(i_2\) are zero, that is,

\[i_1(t_0) = i_2(t_0) = 0\]  

(8.32)

In this case, there is no energy stored, and thus

\[W(t_0) = 0\]  

(8.33)

Next, let us assume that at time \(t_1\), the current \(i_1\) is increased to some finite value, while \(i_2\) is still zero. In other words, we let

\[i_1(t_1) = I_1\]  

(8.34)

and

\[i_2(t_1) = 0\]  

(8.35)
Then, the energy accumulated at this time is

\[ W_1 = \int_{t_0}^{t_1} (p_1 + p_2)\,dt \]  

(8.36)

and since \( i_2(t_1) = 0 \), then \( p_2(t_1) = 0 \) and also \( di_2/dt = 0 \). Therefore, from (8.31) and (8.36) we get

\[ W_1 = \int_{t_0}^{t_1} L_1 i_1 \frac{di_1}{dt}\,dt = L_1 \int_{t_0}^{t_1} i_1\,di_1 = \frac{1}{2} L_1 I_1^2 \]  

(8.37)

Finally, let us at some later time \( t_2 \), maintain \( i_1 \) at its previous value, and increase \( i_2 \) to a finite value, that is, we let

\[ i_1(t_2) = I_1 \]  

(8.38)

and

\[ i_2(t_2) = I_2 \]  

(8.39)

During this time interval, \( di_1/dt = 0 \) and using (8.31) the energy accumulated is

\[ W_2 = \int_{t_1}^{t_2} (p_1 + p_2)\,dt = \int_{t_1}^{t_2} \left( M_{12} I_1 \frac{di_2}{dt} + L_2 i_2 \frac{di_2}{dt} \right)\,dt \]

\[ = \int_{t_1}^{t_2} (M_{12} I_1 + L_2 i_2)\,di_2 = M_{12} I_1 I_2 + \frac{1}{2} L_2 I_2^2 \]  

(8.40)

Therefore, the energy stored in the transformer from \( t_0 \) to \( t_2 \) is from (8.37) and (8.40),

\[ W_{12}^{t_2}_{t_0} = \frac{1}{2} L_1 I_1^2 + M_{12} I_1 I_2 + \frac{1}{2} L_2 I_2^2 \]  

(8.41)

Now, let us reverse the order in which we increase \( i_1 \) and \( i_2 \). That is, in the time interval \( t_0 \leq t \leq t_1 \), we increase \( i_2 \) so that \( i_2(t_1) = I_2 \) while keeping \( i_1 = 0 \). Then, at \( t = t_2 \), we keep \( i_2 = I_2 \) while we increase \( i_1 \) so that \( i_1(t_2) = I_1 \). Using the same steps in equations (8.33) through (8.40), we get

\[ W_{12}^{t_2}_{t_0} = \frac{1}{2} L_1 I_1^2 + M_{21} I_1 I_2 + \frac{1}{2} L_2 I_2^2 \]  

(8.42)

Since relations (8.41) and (8.42) represent the same energy, we must have

\[ M_{12} = M_{21} = M \]  

(8.43)
and thus we can express (8.41) and (8.42) as

\[ W_{t_0}^{I_2} = \frac{1}{2}L_1I_1^2 + MI_1I_2 + \frac{1}{2}L_2I_2^2 \]  \hspace{1cm} (8.44)

Relation (8.44) was derived with the dot markings of Figure 8.18 which is repeated below as Figure 8.19 for convenience.

\[ v_2 = \frac{di_1}{dt} \quad \text{for both circuits} \]

\[ W_{t_0}^{I_2} = \frac{1}{2}L_1I_1^2 - MI_1I_2 + \frac{1}{2}L_2I_2^2 \]  \hspace{1cm} (8.45)

and relations (8.44) and (8.45) can be combined to a single relation as

\[ W_{t_0}^{I_2} = \frac{1}{2}L_1I_1^2 \pm MI_1I_2 + \frac{1}{2}L_2I_2^2 \]  \hspace{1cm} (8.46)

where the sign of \( M \) is positive if both currents enter the dotted (or undotted) terminals, and it is negative if one current enters the dotted (or undotted) terminal while the other enters the undotted (or dotted) terminal.
The currents \( I_1 \) and \( I_2 \) are assume constants and represent the final values of the instantaneous values of the currents \( i_1 \) and \( i_2 \) respectively. We may express (8.46) in terms of the instantaneous currents as

\[
W_{t_2}^{t_0} = \frac{1}{2}L_1 i_1^2 \pm M i_1 i_2 + \frac{1}{2}L_2 i_2^2
\]  

(8.47)

Obviously, the energy on the left side of (8.47) cannot be negative for any values of \( i_1 \), \( i_2 \), \( L_1 \), \( L_2 \), or \( M \). Let us assume first that \( i_1 \) and \( i_2 \) are either both positive or both negative in which case their product is positive. Then, from (8.47) we see that the energy would be negative if

\[
W_{t_2}^{t_0} = \frac{1}{2}L_1 i_1^2 + \frac{1}{2}L_2 i_2^2 - M i_1 i_2
\]  

(8.48)

and the magnitude of the \( M i_1 i_2 \) is greater than the sum of the other two terms on the right side of that expression. To derive an expression relating the mutual inductance \( M \) to the self-inductances \( L_1 \) and \( L_2 \), we add and subtract the term \( \sqrt{L_1L_2} i_1 i_2 \) on the right side of (8.47), and we complete the square. This expression then becomes

\[
W_{t_2}^{t_0} = \frac{1}{2}(\sqrt{L_1} i_1 - \sqrt{L_2} i_2)^2 + \sqrt{L_1L_2} i_1 i_2 - M i_1 i_2
\]  

(8.49)

We now observe that the first term on the right side of (8.49) could be very small and could approach zero, but it can never be negative. Therefore, for the energy to be positive, the second and third terms on the right side of (8.48) must be such that \( \sqrt{L_1L_2} \geq M \) or

\[
M \leq \sqrt{L_1L_2}
\]  

(8.50)

Expression (8.50) indicates that the mutual inductance can never be larger than the geometric mean of the inductances of the two coils between which the mutual inductance exists.

**Note:** The inequality in (8.49) was derived with the assumption that \( i_1 \) and \( i_2 \) have the same algebraic sign. If their signs are opposite, we select the positive sign of (8.47) and we find that (8.50) holds also for this case.

The ratio \( M / \sqrt{L_1L_2} \) is known as the coefficient of coupling and it is denoted with the letter \( k \), that is,

\[
k = \frac{M}{\sqrt{L_1L_2}}
\]  

(8.51)
Energy Stored in a Pair of Mutually Coupled Inductors

Obviously $k$ must have a value between zero and unity, that is, $0 \leq k \leq 1$. Physically, $k$ provides a measure of the proximity of the primary and secondary coils. If the coils are far apart, we say that they are *loose-coupled* and $k$ has a small value, typically between 0.01 and 0.1. For *close-coupled* circuits, $k$ has a value of about 0.5. Power transformers have a $k$ between 0.90 and 0.95. The value of $k$ is exactly unity only when the two coils are coalesced into a single coil.

**Example 8.5**

For the transformer of Figure 8.21 compute the energy stored at $t = 0$ if:

a. $i_1 = 50 \text{ mA}$ and $i_2 = 25 \text{ mA}$  

b. $i_1 = 0$ and $i_2 = 20 \sin 377t \text{ mA}$  

c. $i_1 = 15 \cos 377t \text{ mA}$ and $i_2 = 40 \sin (377t + 60^\circ) \text{ mA}$

![Transformer for Example 8.5](image)

**Figure 8.21. Transformer for Example 8.5**

**Solution:**

Since the currents enter the dotted terminals, we use (8.45) with the plus (+) sign for the mutual inductance term, that is,

$$W(t) = \frac{1}{2}L_1i_1^2 + Mi_1i_2 + \frac{1}{2}L_2i_2^2$$  

(8.52)

Then,

a.  

$$W|_{t=0} = 0.5 \times 50 \times 10^{-3} \times (50 \times 10^{-3})^2 + 20 \times 10^{-3} \times 50 \times 10^{-3} \times 25 \times 10^{-3}$$

$$+ 0.5 \times 50 \times 10^{-3} \times (25 \times 10^{-3})^2 = 103 \times 10^{-6} \text{ J} = 103 \mu \text{J}$$

b.  

Since $i_1 = 0$ and $i_2 = 20 \sin 377t|_{t=0} = 0$, it follows that
c.
\[ W|_{t=0} = 0 \]

\[ W|_{t=0} = 0.5 \times 50 \times 10^{-3} \times (15 \times 10^{-3})^2 + 20 \times 10^{-3} \times 15 \times 10^{-3} \times 40 \times 10^{-3} \times \sin(60^\circ) \]
\[ + 0.5 \times 50 \times 10^{-3} \times (40 \times 10^{-3} \times \sin(60^\circ))^2 = 46 \times 10^{-6} \text{ J} = 46 \text{ \mu J} \]

### 8.7 Circuits with Linear Transformers

A linear transformer is a four-terminal device in which the voltages and currents in the primary coils are linearly related.

The transformer shown in figure 8.22 is a linear transformer. This transformer contains a voltage source in the primary, a load resistor in the secondary, and the resistors \( R_1 \) and \( R_2 \) represent the resistances of the primary and secondary coils respectively. Moreover, the primary is referenced to directly to ground, but the secondary is referenced to a DC voltage source \( V_0 \) and thus it is said that the secondary of the transformer has a DC isolation.

![Transformer with DC isolation](image)

**Figure 8.22. Transformer with DC isolation**

Application of KVL around the primary and secondary circuits yields the loop equations

\[ v_{in} = R_1 i_1 + L_1 \frac{di_1}{dt} - M \frac{di_2}{dt} \]

\[ 0 = -M \frac{di_1}{dt} + L_2 \frac{di_2}{dt} + (R_2 + R_{LOAD}) \]

and we see that the instantaneous values of the voltages and the currents are not affected by the presence of the DC voltage source \( V_0 \) since we would have obtained the same equations had we let \( V_0 = 0 \).
Example 8.6

For the transformer shown in Figure 8.23, find the total response of $i_2$ for $t > 0$ given that $M = 100 \text{ mH}$ and $i_1(0^-) = i_2(0^-) = 0$. Use MATLAB to sketch $i_2$ for $0 \leq t \leq 5 \text{ s}$.

![Figure 8.23. Transformer for Example 8.6](image)

**Solution:**

The total response consists of the summation of the forced and natural responses, that is,

$$i_{2T} = i_{2f} + i_{2n}$$  \hspace{1cm} (8.54)

and since the applied voltage is constant (DC), no steady-state (forced) voltage is produced in the secondary and thus $i_{2f} = 0$.

For $t > 0$ the $s$-domain circuit is shown in Figure 8.24.

![Figure 8.24. The $s$-domain circuit for the transformer of Example 8.6](image)

The loop equations for this transformer are

$$(3s + 100)I_1(s) - 2sI_2(s) = 24/s$$

$$-2sI_1(s) + (5s + 1200)I_2(s) = 0$$  \hspace{1cm} (8.55)

Since we are interested only in $I_2(s)$, we will use Cramer’s rule.
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\[ I_2(s) = \frac{\begin{bmatrix} 3s + 100 & 24/s \\ -2s & 0 \end{bmatrix}}{\begin{bmatrix} 3s + 100 & -2s \\ -2s & 5s + 1200 \end{bmatrix}} = \frac{48}{11s^2 + 4100s + 120000} = \frac{4.36}{s^2 + 372.73s + 10909.01} \]

or

\[ I_2(s) = \frac{4.36}{(s + 340.71)(s + 32.02)} \]

and by partial fraction expansion,

\[ I_2(s) = \frac{4.36}{(s + 340.71)(s + 32.02)} = \frac{r_1}{s + 340.71} + \frac{r_2}{s + 32.02} \tag{8.56} \]

from which

\[ r_1 = \frac{4.36}{s + 32.02} \bigg|_{s = -340.71} = -0.01 \tag{8.57} \]

\[ r_2 = \frac{4.36}{s + 340.71} \bigg|_{s = -32.02} = 0.01 \tag{8.58} \]

By substitution into (8.56), we get

\[ I_2(s) = \frac{0.01}{s + 32.02} + \frac{-0.01}{s + 340.71} \tag{8.59} \]

and taking the Inverse Laplace of (8.59) we get

\[ i_{2n} = 0.01(e^{-32.02t} - e^{-340.71t}) \tag{8.60} \]

Using the following MATLAB code we get the plot shown on Figure 8.25.

\[ t=0: 0.001: 0.2; i2n=0.01.*(\text{exp}(-32.02*t)-\text{exp}(-340.71.*t)); \text{plot}(t,i2n); \text{grid} \]

Example 8.7

For the transformer of Figure 8.26, find the steady-state (forced) response of \( v_{out} \).

Solution:

The \( s \)-domain equivalent circuit is shown in Figure 8.27.

We could use the same procedure as in the previous example, but it is easier to work with the transfer function \( G(s) \).
The loop equations for the transformer of Figure 8.27 are:
and by Cramer’s rule,

\[
I_2(s) = \frac{
\begin{bmatrix}
3s + 10 + 1/0.1s & V_{in}(s) \\
-(2s + 1/0.1s) & 0
\end{bmatrix}
\begin{bmatrix}
3s + 10 + 1/0.1s \\
-(2s + 1/0.1s)
\end{bmatrix}
- \begin{bmatrix}
3s + 10 + 1/0.1s & -(2s + 1/0.1s) \\
-(2s + 1/0.1s) & 5s + 100 + 1/0.1s
\end{bmatrix}
\begin{bmatrix}
V_{in}(s) \\
0
\end{bmatrix}
}{(3s + 10 + 1/0.1s)(5s + 100 + 1/0.1s) - \left(2s + 1/0.1s\right)^2}
\]

or

\[
I_2(s) = \frac{(2s + 10/s)V_{in}(s)}{11s^2 + 350s + 1040 + 1100/s} = \frac{(2s^2 + 10)V_{in}(s)}{11s^3 + 350s^2 + 1040s + 1100}
\]

\[
= \frac{(0.18s^2 + 0.91)V_{in}(s)}{s^3 + 31.82s^2 + 94.55s + 100}
\]

From Figure 8.27 we see that

\[
V_{out}(s) = 100 \cdot I_2(s) = 100 \cdot \frac{(0.18s^2 + 0.91)V_{in}(s)}{s^3 + 31.82s^2 + 94.55s + 100} = \frac{(18s^2 + 91)V_{in}(s)}{s^3 + 31.82s^2 + 94.55s + 100}
\] (8.62)

and

\[
G(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{18s^2 + 91}{s^3 + 31.82s^2 + 94.55s + 100}
\] (8.63)

The input is a sinusoid, that is,

\[
v_{in} = 170\cos377t \text{ V}
\]

and since we are interested in the steady-state response, we let

\[
s = j\omega = j377
\]

and thus

\[
V_{in}(s) = V_{in}(j\omega) = 170\angle0^\circ
\]

From (8.63) we get:

\[
V_{out}(j\omega) = \frac{-2.56 \times 10^6 + 91}{-j5.36 \times 10^7 - 4.52 \times 10^6 + j3.56 \times 10^4 + 100} = \frac{-4.35 \times 10^8 \angle0^\circ}{-4.52 \times 10^6 - j5.36 \times 10^7}
\]

or
Reflected Impedance in Transformers

\[ V_{out}(j\omega) = \frac{4.35 \times 10^8 \angle 180^\circ}{5.38 \times 10^7 \angle -94.82^\circ} = \frac{43.5 \angle 180^\circ}{5.38 \angle -94.82^\circ} = 8.09 \angle 274.82^\circ = 8.09 \angle -85.18^\circ \]  \hspace{1cm} (8.64)

and in the \( t \)-domain,

\[ v_{out}(t) = 8.09 \cos(377t - 85.18^\circ) \]  \hspace{1cm} (8.65)

The expression of (8.65) indicates that the transformer of this example is a step-down transformer.

8.8 Reflected Impedance in Transformers

In this section, we will see how the load impedance of the secondary can be reflected into the primary.

Let us consider the transformer phasor circuit of Figure 8.28. We assume that the resistance of the primary and secondary coils is negligible.

\[ \begin{align*}
V_S & \quad \cdots \quad I_1 \\
L_1 & \quad \cdots \quad V_1 \\
& \quad \cdots \quad L_2 \\
& \quad \cdots \quad V_2 \\
& \quad \cdots \quad I_2 \\
M & \quad \cdots \quad Z_{LOAD} \\
\end{align*} \]

Figure 8.28. Circuit for the derivation of reflected impedance

By KVL the loops equations in phasor notation are:

\[ j\omega L_1 I_1 - j\omega M I_2 = V_S \]  \hspace{1cm} (8.66)

or

\[ I_2 = \frac{j\omega L_1 I_1 - V_S}{j\omega M} \]  \hspace{1cm} (8.67)

and

\[ -j\omega M I_1 + (j\omega L_2 + Z_{LOAD})I_2 = 0 \]  \hspace{1cm} (8.68)

or

\[ I_2 = \frac{j\omega M I_1}{(j\omega L_2 + Z_{LOAD})} \]  \hspace{1cm} (8.69)
Equating the right sides of (8.67) and (8.69) we get:

\[
\frac{j\omega L_1 I_1 - V_S}{j\omega M} = \frac{j\omega M I_1}{(j\omega L_2 + Z_{\text{LOAD}})}
\]  

(8.70)

Solving for \(V_S\) we get:

\[
V_S = \left[ j\omega L_1 - \frac{(j\omega M)^2}{(j\omega L_2 + Z_{\text{LOAD}})} \right] I_1
\]  

(8.71)

and dividing \(V_S\) by \(I_1\) we obtain the input impedance \(Z_{\text{in}}\) as

\[
Z_{\text{in}} = \frac{V_S}{I_1} = j\omega L_1 + \frac{\omega^2 M^2}{j\omega L_2 + Z_{\text{LOAD}}}
\]  

(8.72)

The first term on the right side of (8.72) represents the reactance of the primary. The second term is a result of the mutual coupling and it is referred to as the \textit{reflected impedance}. It is denoted as \(Z_R\), i.e.,

\[
Z_R = \frac{\omega^2 M^2}{j\omega L_2 + Z_{\text{LOAD}}}
\]  

(8.73)

From (8.73), we make two important observations:

1. The reflected impedance \(Z_R\) does not depend on the dot locations on the transformer. For instance, if either dot in the transformer of the previous page is placed on the opposite terminal, the sign of the mutual term changes from \(M\) to \(-M\). But since \(Z_R\) varies as \(M^2\), its sign remains unchanged.

2. Let \(Z_{\text{LOAD}} = R_{\text{LOAD}} + jX_{\text{LOAD}}\). Then, we can rewrite (8.73) as

\[
Z_R = \frac{\omega^2 M^2}{j\omega L_2 + R_{\text{LOAD}} + jX_{\text{LOAD}}} = \frac{\omega^2 M^2}{R_{\text{LOAD}} + j(X_{\text{LOAD}} + \omega L_2)}
\]  

(8.74)

To express (8.74) as the sum of a real and an imaginary component, we multiply both numerator and denominator by the complex conjugate of the denominator. Then,

\[
Z_R = \frac{\omega^2 M^2 R_{\text{LOAD}}}{R_{\text{LOAD}}^2 + (X_{\text{LOAD}} + \omega L_2)^2} - j\frac{\omega^2 M^2 (X_{\text{LOAD}} + \omega L_2)}{R_{\text{LOAD}}^2 + (X_{\text{LOAD}} + \omega L_2)^2}
\]  

(8.75)

The imaginary part of (8.75) represents the reflected reactance and we see that it is negative. That is, the reflected reactance is opposite to that of the net reactance \(X_{\text{LOAD}} + \omega L_2\) of the secondary.
Reflected Impedance in Transformers

Therefore, if $X_{LOAD}$ is a capacitive reactance whose magnitude is less than $\omega L_2$, or if it is an inductive reactance, then the reflected reactance is capacitive. However, if $X_{LOAD}$ is a capacitive reactance whose magnitude is greater than $\omega L_2$, the reflected reactance is inductive. In the case where the magnitude of $X_{LOAD}$ is capacitive and equal to $\omega L_2$, the reflected reactance is zero and the transformer operates at resonant frequency. In this case, the reflected impedance is purely real since (8.75) reduces to

$$Z_R = \frac{\omega^2 M^2}{R_{LOAD}}$$  

(8.76)

Example 8.8

In the transformer circuit of Figure 8.29, $Z_S$ represents the internal impedance of the voltage source $V_S$.

Find:

a. $Z_{in}$

b. $I_1$

c. $I_2$

d. $V_I$

e. $V_2$

\[\omega = 377 \frac{r}{s}\]

\[V_S = 120 \angle 0^\circ\]

\[Z_{LOAD} = 10 - j\frac{7540}{\omega}\] \(\Omega\)

Figure 8.29. Transformer for Example 8.8
Solution:

a. From (8.72)

\[ Z_{in} = \frac{V_S}{I_1} = j\omega L_1 + \frac{\omega^2 M^2}{j\omega L_2 + Z_{LOAD}} \]

and we must add \( Z_s = 2 \Omega \) to it. Therefore, for the transformer of this example,

\[ Z_{in} = j\omega L_1 + \frac{\omega^2 M^2}{j\omega L_2 + Z_{LOAD}} + 2 = j75.4 + \frac{142129 \times 0.01}{j113.1 + 10 - j20} + 2 \]

\[ = 3.62 + j60.31 = 60.42\angle86.56^\circ \Omega \]

b.

\[ I_1 = \frac{V_S}{Z_{in}} = \frac{120\angle0^\circ}{60.42\angle86.56^\circ \Omega} = 1.98\angle-86.56^\circ A \]

c. By KVL

\[-j\omega M I_1 + (j\omega L_2 + Z_{LOAD}) I_2 = 0\]

or

\[ I_2 = \frac{j\omega M}{j\omega L_2 + Z_{LOAD}} I_1 = \frac{j37.7}{j113.1 + 10 - j20}1.98\angle86.56^\circ = \frac{74.88\angle3.04^\circ}{93.64\angle83.87^\circ} = 0.8\angle-80.83^\circ A \]

d.

\[ V_1 = j\omega L_1 I_1 - j\omega M I_2 = 75.4\angle90^\circ \times 1.98\angle-86.56^\circ - 37.7\angle90^\circ \times 0.8\angle-80.83^\circ \]

\[ = 149.29\angle3.04^\circ - 30.15\angle9.17^\circ = 149.08 + j7.92 - 30.15 - j4.8 = 118.9\angle1.5^\circ V \]

e.

\[ V_2 = Z_{LOAD} \cdot I_2 = (10 - j20)0.8\angle-80.83^\circ = 22.36\angle-63.43^\circ \times 0.8\angle-80.83^\circ = 17.89\angle-144.26^\circ V \]

8.9 The Ideal Transformer

An ideal transformer is one in which the coefficient of coupling is almost unity, and both the primary and secondary inductive reactances are very large in comparison with the load impedances. The primary and secondary coils have many turns wound around a laminated iron-core and are arranged so that the entire flux links all the turns of both coils.

An important parameter of an ideal transformer is the turns ratio \( a \) which is defined as the ratio of the number of turns on the secondary, \( N_2 \), to the number of turns of the primary \( N_1 \), that is,
The flux produced in a winding of a transformer due to a current in that winding is proportional to the product of the current and the number of turns on the winding. Therefore, letting $\alpha$ be a constant of proportionality which depends on the physical properties of the transformer, for the primary and secondary windings we have:

\[
\phi_{11} = \alpha N_1 i_1 \\
\phi_{22} = \alpha N_2 i_2
\]  

(8.78)

The constant $\alpha$ is the same for the primary and secondary windings because we have assumed that the same flux links both coils and thus both flux paths are identical. We recall from (8.8) and (8.14) that

\[
\lambda_1 = N_1 \phi_{11} = L_1 i_1 \\
\lambda_2 = N_2 \phi_{22} = L_2 i_2
\]

(8.79)

Then, from (8.78) and (8.79) we get:

\[
N_1 \phi_{11} = L_1 i_1 = \alpha N_1^2 i_1
\]

(8.80)

or

\[
N_2 \phi_{22} = L_2 i_2 = \alpha N_2^2 i_2
\]

or

\[
L_1 = \alpha N_1^2 \\
L_2 = \alpha N_2^2
\]

(8.81)

Therefore,

\[
\frac{L_2}{L_1} = \left(\frac{N_2}{N_1}\right)^2 = a^2
\]

(8.82)

From (8.69),

\[
I_2 = \frac{j\omega M I_1}{j\omega L_2 + Z_{LOAD}}
\]

(8.83)

or

\[
\frac{I_2}{I_1} = \frac{j\omega M}{j\omega L_2 + Z_{LOAD}}
\]

(8.84)
and since \( j\omega L_2 \gg Z_{LOAD} \), (8.84) reduces to

\[
\frac{I_2}{I_1} = \frac{j\omega M}{j\omega L_2} = \frac{M}{L_2}\]

(8.85)

For the case of unity coupling,

\[
k = \frac{M}{\sqrt{L_1L_2}} = 1
\]

(8.86)

or

\[
M = \sqrt{L_1L_2}
\]

(8.87)

and by substitution of (8.87) into (8.85) we get:

\[
\frac{I_2}{I_1} = \frac{\sqrt{L_1L_2}}{L_2} = \sqrt{\frac{L_1}{L_2}}
\]

(8.88)

From (8.82) and (8.88), we obtain the important relation

\[
\frac{I_2}{I_1} = \frac{1}{a}
\]

(8.89)

Also, from (8.77) and (8.89),

\[
N_1I_1 = N_2I_2
\]

(8.90)

and this relation indicates that if \( N_2 < N_1 \), the current \( I_2 \) is larger than \( I_1 \).

The primary and secondary voltages are also related to the turns ratio \( a \). To find this relation, we define the secondary or load voltage \( V_2 \) as

\[
V_2 = Z_{LOAD}I_2
\]

(8.91)

and the primary voltage \( V_1 \) across \( L_1 \) as

\[
V_1 = Z_{in}I_1
\]

(8.92)

From (8.72),

\[
Z_{in} = \frac{V_s}{I_1} = j\omega L_1 + \frac{\omega^2 M^2}{j\omega L_2 + Z_{LOAD}}
\]

(8.93)
and for \( k = 1 \)

\[
M^2 = L_1 L_2
\]

Then, (8.93) becomes

\[
Z_{in} = j\omega L_1 + \frac{\omega^2 L_1 L_2}{j\omega L_2 + Z_{LOAD}} \tag{8.94}
\]

Next, from (8.82)

\[
L_2 = a^2 L_1 \tag{8.95}
\]

Substitution of (8.95) into (8.94) yields

\[
Z_{in} = j\omega L_1 + \frac{\omega^2 a^2 L_1^2}{j\omega a^2 L_1 + Z_{LOAD}} \tag{8.96}
\]

and if we let \( j\omega L_1 \to \infty \), both terms on the right side of (8.96) become infinite and we get an indeterminate result. To work around this problem, we combine these terms and we get:

\[
Z_{in} = \frac{-\omega^2 a^2 L_1^2 + j\omega L_1 Z_{LOAD} + \omega^2 a^2 L_1^2}{j\omega a^2 L_1 + Z_{LOAD}} = \frac{j\omega L_1 Z_{LOAD}}{j\omega a^2 L_1 + Z_{LOAD}}
\]

and as \( j\omega L_1 \to \infty \),

\[
Z_{in} = \frac{Z_{LOAD}}{a^2} \tag{8.97}
\]

Finally, substitution of (8.97) into (8.92) yields

\[
V_1 = \frac{Z_{LOAD}}{a^2} I_1 \tag{8.98}
\]

and by division of (8.91) by (8.98) we get:

\[
\frac{V_2}{V_1} = \frac{Z_{LOAD} I_2}{(Z_{LOAD}/a^2) I_1} = a^2 \cdot \frac{I_2}{I_1} = a
\]

or

\[
\frac{V_2}{V_1} = a \tag{8.100}
\]
also, from the current and voltage relations of (8.88) and (8.99),

\[ V_2 I_2 = V_1 I_1 \]  \hspace{1cm} (8.101)

that is, the volt-amperes of the secondary and the primary are equal.

An ideal transformer is represented by the network of Figure 8.30.

**Figure 8.30. Ideal transformer representation**

### 8.10 Impedance Matching

An ideal (iron-core) transformer can be used as an impedance level changing device. We recall from basic circuit theory that to achieve maximum power transfer, we must adjust the resistance of the load to make it equal to the resistance of the voltage source. But this is not always possible. A power amplifier for example, has an internal resistance of several thousand ohms. On the other hand, a speaker which is to be connected to the output of a power amplifier has a fixed resistance of just a few ohms. In this case, we can achieve maximum power transfer by inserting an iron-core transformer between the output of the power amplifier and the input of the speaker as shown in Figure 8.31 where \( N_2 < N_1 \).

**Figure 8.31. Transformer used as impedance matching device**
Let us suppose that in Figure 8.31 the amplifier internal impedance is $80000 \ \Omega$ and the impedance of the speaker is only $8 \ \Omega$. We can find the appropriate turns ratio $N_2/N_1 = a$ using (8.97), that is,

$$Z_{in} = \frac{Z_{LOAD}}{a^2}$$

(8.102)

or

$$a = \frac{N_2}{N_1} = \frac{\sqrt{Z_{LOAD}}}{\sqrt{Z_{in}}} = \sqrt{\frac{8}{80000}} = \sqrt{\frac{1}{10000}} = \frac{1}{100}$$

or

$$\frac{N_1}{N_2} = 100$$

(8.103)

that is, the number of turns in the primary must be 100 times the number of the turns in the secondary.

### 8.11 A Simplified Transformer Equivalent Circuit

In analyzing networks containing ideal transformers, it is very convenient to replace the transformer by an equivalent circuit before the analysis. Consider the transformer circuit of Figure 8.32.

![Figure 8.32. Circuit to be simplified](image)

From (8.97)

$$Z_{in} = \frac{Z_{LOAD}}{a^2}$$

The input impedance seen by the voltage source $V_S$ in the circuit of Figure 8.32 is

$$Z_{in} = Z_S + \frac{Z_{LOAD}}{a^2}$$

(8.104)

and thus the circuit of Figure 8.32 can be replaced with the simplified circuit shown in Figure 8.33.
The voltages and currents can now be found from the simple series circuit if Figure 8.33.

8.12 Thevenin Equivalent Circuit

Let us consider again the circuit of Figure 8.32. This time we want to find the Thevenin equivalent to the left of the secondary terminals and replace the primary by its Thevenin equivalent at points $x$ and $y$ as shown in Figure 8.34.

If we open the circuit at points $x$ and $y$ as shown in Figure 8.34, we find the Thevenin voltage as $V_{TH} = V_{OC} = V_{xy}$. Since the secondary is now an open circuit, we have $I_2 = 0$, and also $I_1 = 0$ because $I_1 = aI_2$. Since no voltage appears across $Z_S$, $V_1 = V_S$ and $V_{2\,oc} = aV_1 = aV_S$. Then,

$$V_{TH} = V_{OC} = V_{xy} = aV_S$$

We will find the Thevenin impedance $Z_{TH}$ from the relation

$$Z_{TH} = \frac{V_{OC}}{I_{SC}}$$

The short circuit current $I_{SC}$ is found from
Thevenin Equivalent Circuit

\[ I_{SC} = I_2 = \frac{I_1}{a} = \frac{V_S/Z_S}{a} = \frac{V_S}{aZ_S} \]  
(8.107)

and by substitution into (8.106),

\[ Z_{TH} = \frac{aV_S}{V_S/aZ_S} = a^2 Z_S \]

The Thevenin equivalent circuit with the load connected to it is shown in Figure 8.35.

Figure 8.35. The Thevenin equivalent of the transformer circuit in Figure 8.34

The circuit of Figure 8.35 was derived with the assumption that the dots are placed as shown in Figure 8.34. If either dot is reversed, we simply replace \( a \) by \( -a \).

Example 8.9

For the circuit of Figure 8.36, find \( V_2 \).

Solution:

We will replace the given circuit with its Thevenin equivalent. First, we observe that the dot in the secondary has been reversed, and therefore we will replace \( a \) by \( -a \). The Thevenin equivalent is

Figure 8.36. Circuit for Example 8.9

* Since \( V_2 = 0 \) and \( V_2/V_1 = a \) or \( aV_1 = V_2 \) it follows that \( V_1 = 0 \) also.
obtained by multiplying $V_S$ and the dependent source by $-10$ and the $10 \Omega$ resistor by $(-a)^2 = 100$. With these modifications we obtain the circuit of Figure 8.37.

![Circuit Diagram]

*Figure 8.37. The Thevenin equivalent of the circuit of Example 8.9*

Now, by application of KCL

$$\frac{V_2 - (-80\angle 0^\circ)}{10^3} - (-10^{-3}V_2) + \frac{V_2}{60+j80} = 0$$

$$\frac{V_2}{10^3} + \frac{V_2}{10^3} + \frac{(60-j80)V_2}{10000} = \frac{-80}{10^3}$$

$$2V_2 + (6-j8)V_2 = -80$$

$$8(1-j1)V_2 = 80\angle 180^\circ$$

$$(\sqrt{2}\angle-45^\circ)V_2 = 10\angle 180^\circ$$

or

$$V_2 = \frac{10}{\sqrt{2}}\angle 225^\circ = 5\sqrt{2}\angle -135^\circ$$

Other equivalent circuits can be developed from the equations of the primary and secondary voltages and currents.

Consider, for example the linear transformer circuit of Figure 8.38.

From (8.30), the primary and secondary voltages and currents are:

$$v_1 = L_1\frac{di_1}{dt} + M\frac{di_2}{dt}$$

$$v_2 = M\frac{di_1}{dt} + L_2\frac{di_2}{dt}$$

(8.108)
and these equations are satisfied by the equivalent circuit shown in Figure 8.39.

If we rearrange the equations of (8.108) as

\[
    v_1 = (L_1 - M) \frac{di_1}{dt} + M \left( \frac{di_1}{dt} + \frac{di_2}{dt} \right)
\]

\[
    v_2 = M \left( \frac{di_1}{dt} + \frac{di_2}{dt} \right) + (L_2 - M) \frac{di_2}{dt}
\]

these equations are satisfied by the circuit of Figure 8.40.
8.13 Summary

- Inductance is associated with the magnetic field which is always present when there is an electric current.

- The magnetic field loops are circular in form and are called lines of magnetic flux.

- The magnetic flux is denoted as $\phi$ and the unit of magnetic flux is the weber (Wb).

- If there are $N$ turns and we assume that the flux $\phi$ passes through each turn, the total flux denoted as $\lambda$ is called flux linkage. Then,

$$\lambda = N\phi$$

- A linear inductor one in which the flux linkage is proportional to the current through it, that is,

$$\lambda = Li$$

where the constant of proportionality $L$ is called inductance in webers per ampere.

- Faraday’s law of electromagnetic induction states that

$$v = \frac{d\lambda}{dt}$$

- Lenz’s law states that whenever there is a change in the amount of magnetic flux linking an electric circuit, an induced voltage of value directly proportional to the time rate of change of flux linkages is set up tending to produce a current in such a direction as to oppose the change in flux.

- A linear transformer is a four-terminal device in which the voltages and currents in the primary coils are linearly related.

- In a linear transformer, when there is no current in the secondary winding the voltages are

$$v_1 = L_1 \frac{di_1}{dt} \quad \text{and} \quad v_2 = M_{21} \frac{di_1}{dt}$$

if $i_1 \neq 0$ and $i_2 = 0$

- In a linear transformer, when there is no current in the primary winding, the voltages are

$$v_2 = L_2 \frac{di_2}{dt} \quad \text{and} \quad v_1 = M_{12} \frac{di_2}{dt}$$

if $i_1 = 0$ and $i_2 \neq 0$

- In a linear transformer, when there is a current in both the primary and secondary windings, the voltages are
Summary

The voltage terms are referred to as self-induced voltages.

The voltage terms are referred to as mutual voltages.

The polarity of the mutual voltages is denoted by the dot convention. If a current $i$ entering the dotted (undotted) terminal of one coil induces a voltage across the other coil with positive polarity at the dotted (undotted) terminal of the other coil, the mutual voltage term has a positive sign. If a current $i$ entering the undotted (dotted) terminal of one coil induces a voltage across the other coil with positive polarity at the dotted (undotted) terminal of the other coil, the mutual voltage term has a negative sign.

If the polarity (dot) markings are not given, they can be established by using the right-hand rule which states that if the fingers of the right hand encircle a winding in the direction of the current, the thumb indicates the direction of the flux. Thus, in an ideal transformer with primary and secondary windings $L_1$ and $L_2$ and currents $i_1$ and $i_2$ respectively, we place a dot at the upper end of $L_1$ and assume that the current $i_1$ enters the top end thereby producing a flux in the clockwise direction. Next, we want the current in $L_2$ to enter the end which will produce a flux in the same direction, in this case, clockwise.

The energy stored in a pair of mutually coupled inductors is given by

$$W_{i_0}^{i_2} = \frac{1}{2} L_1 i_1^2 + M i_1 i_2 + \frac{1}{2} L_2 i_2^2$$

where the sign of $M$ is positive if both currents enter the dotted (or undotted) terminals, and it is negative if one current enters the dotted (or undotted) terminal while the other enters the undotted (or dotted) terminal.

The ratio

$$k = \frac{M}{\sqrt{L_1 L_2}}$$
is known as the coefficient of coupling and $k$ provides a measure of the proximity of the primary and secondary coils. If the coils are far apart, we say that they are *loose-coupled*, and $k$ has a small value, typically between $0.01$ and $0.1$. For *close-coupled* circuits, $k$ has a value of about $0.5$. Power transformers have a $k$ between $0.90$ and $0.95$. The value of $k$ is exactly unity only when the two coils are coalesced into a single coil.

- If the secondary of a linear transformer is referenced to a DC voltage source $V_0$, it is said that the secondary has DC isolation.

- In a linear transformer, the load impedance of the secondary can be reflected into the primary using the relation

\[
Z_R = \frac{\omega^2 M^2}{j\omega L_2 + Z_{LOAD}}
\]

where $Z_R$ is referred to as the reflected impedance.

- An ideal transformer is one in which the coefficient of coupling is almost unity, and both the primary and secondary inductive reactances are very large in comparison with the load impedances. The primary and secondary coils have many turns wound around a laminated iron-core and are arranged so that the entire flux links all the turns of both coils.

- In an ideal transformer number of turns on the primary $N_1$ and the number of turns on the secondary $N_2$ are related to the primary and secondary currents $I_1$ and $I_2$ respectively as

\[
N_1 I_1 = N_2 I_2
\]

- An important parameter of an ideal transformer is the turns ratio $a$ which is defined as the ratio of the number of turns on the secondary, $N_2$, to the number of turns of the primary $N_1$, that is,

\[
a = \frac{N_2}{N_1}
\]

- In an ideal transformer the turns ratio $a$ relates the primary and secondary currents as

\[
\frac{I_2}{I_1} = \frac{1}{a}
\]

- In an ideal transformer the turns ratio $a$ relates the primary and secondary voltages as

\[
\frac{V_2}{V_1} = a
\]
Summary

- In an ideal transformer the volt-amperes of the primary and the secondary are equal, that is,

\[ V_2 I_2 = V_1 I_1 \]

- An ideal transformer can be used as an impedance matching device by specifying the appropriate turns ratio \( N_2 / N_1 = a \). Then,

\[ Z_{in} = \frac{Z_{LOAD}}{a^2} \]

- In analyzing networks containing ideal transformers, it is very convenient to replace the transformer by an equivalent circuit before the analysis. One method is presented in Section 8.11.

- An ideal transformer can be replaced by a Thevenin equivalent as discussed in Section 8.12.
8.14 Exercises

1. For the transformer of Figure 8.41, find $v_2$ for $t > 0$.

![Figure 8.41. Circuit for Exercise 1](image1)

2. For the transformer circuit of Figure 8.42, find the phasor currents $I_1$ and $I_2$.

![Figure 8.42. Circuit for Exercise 2](image2)

3. For the network of Figure 8.43, find the transfer function $G(s) = V_{OUT}(s)/V_{IN}(s)$.

![Figure 8.43. Circuit for Exercise 3](image3)
4. For the transformer of Figure 8.44, find the average power delivered to the 4 Ω resistor.

\[ v_s = 4\cos 3t \text{ V} \]

![Figure 8.44. Circuit for Exercise 4](image)

5. Replace the transformer of Figure 8.45 by a Thevenin equivalent and then compute \( V_1, V_2, I_1 \) and \( I_2 \)

6. For the circuit of Figure 8.46, compute the turns ratio \( a \) so that maximum power will be delivered to the 10 KΩ resistor.

![Figure 8.46. Circuit for Exercise 6](image)
8.15 Solutions to Exercises

1.

Application of KVL in the primary yields

\[ 2i_1 + L_1 \frac{di_1}{dt} = 8u_0(t) \]

From (1) we find that \( i_1f \) and \( i_{1n} \) is found from the characteristic equation \( s + 2 = 0 \) from which \( s = -2 \) and thus \( i_{1n} = Ae^{-2t} \). Then,

\[ i_1 = 4 + Ae^{-2t} \quad (2) \]

Since we are not told otherwise, we will assume that \( i_1(0^-) = 0 \) and from (2) \( 0 = 4 + Ae^0 \) or \( A = -4 \) and by substitution into (2)

\[ i_1 = 4(1 - 4e^{-2t}) \]

The voltage \( v_2 \) is found from

\[ v_2 = M \frac{di_1}{dt} + L_2 \frac{di_2}{dt} \]

and since \( i_2 = 0 \),

\[ v_2 = 1 \cdot \frac{di_1}{dt} = \frac{d}{dt}[4(1 - 4e^{-2t})] = 8e^{-2t} \text{ V} \]
2.

The mesh equations for primary and secondary are:

\[(1 + j)I_1 - jI_2 = 10\angle0^\circ\]
\[-jI_1 + (2 - j2)I_2 = 0\]

By Cramer's rule,

\[I_1 = D_1/\Delta \quad I_2 = D_2/\Delta\]

where

\[\Delta = \begin{bmatrix} (1 + j) & -j1 \\ -j1 & (2 - j2) \end{bmatrix} = 5\]
\[D_1 = \begin{bmatrix} 10\angle0^\circ & -j1 \\ 0 & (2 - j2) \end{bmatrix} = 20(1 - j)\]
\[D_2 = \begin{bmatrix} (1 + j) & 10\angle0^\circ \\ -j1 & 0 \end{bmatrix} = j10\]

Thus,

\[I_1 = \frac{20(1 - j)}{5} = 4(1 - j) = 4\sqrt{2} \angle -45^\circ \, A\]
\[I_2 = \frac{j10}{5} = j2 = 2\angle90^\circ \, A\]

Check with MATLAB:

```matlab
Z=[1+j -j; -j 2-j2]; V=[10 0]; l=Z\V;
fprintf('magI1 = %5.2f A \t', abs(l(1))); fprintf('phaseI1 = %5.2f deg \t',angle(l(1))*180/pi);
fprintf(' \n');
fprintf('magI2 = %5.2f A \t', abs(l(2))); fprintf('phaseI2 = %5.2f deg \t',angle(l(2))*180/pi);
```

magI1 = 5.66 A \ phaseI1 = -45.00 deg
magI2 = 2.00 A \ phaseI2 = 90.00 deg
We will find $V_{OUT}(s)$ from $V_{OUT}(s) = (1 \Omega)I_3$. The three mesh equations in matrix form are:

$$
\begin{bmatrix}
(s + 1) & -0.5s & -0.5s \\
-0.5s & (s + 1) & -0.5s \\
-0.5s & -0.5s & (s + 1)
\end{bmatrix}
\begin{bmatrix}
I_1 \\
I_2 \\
I_3
\end{bmatrix}
= \begin{bmatrix}
1 \\
0 \\
0
\end{bmatrix} \cdot V_{IN}(s)
$$

We will use MATLAB to find the determinant $\Delta$ of the $3 \times 3$ matrix.

```matlab
syms s
delta=[s+1 -0.5*s -0.5*s; -0.5*s s+1 -0.5*s; -0.5*s -0.5*s s+1]; det_delta=det(delta)
det_delta = 9/4*s^2+3*s+1
d3=[s+1 -0.5*s -0.5*s; -0.5*s s+1 -0.5*s; 1 0 0]; det_d3=det(d3)
det_d3 = 3/4*s^2+1/2*s
I3=det_d3/det_delta
I3 = (3/4*s^2+1/2*s) / (9/4*s^2+3*s+1)
simplify(I3)
ans = s / (3*s+2)
```

Therefore,

$$
V_{OUT}(s) = I_3 \cdot V_{IN}(s) = \frac{s}{(3s+2)} \cdot V_{IN}(s)
$$

and

$$
G(s) = \frac{V_{OUT}(s)}{V_{IN}(s)} = \frac{s}{(3s+2)}
$$
For this exercise, \( P_{\text{ave}} \, 4 \, \Omega = \frac{1}{2} (I_{4\Omega})^2 4 \) and thus we need to find \( I_{4\Omega} \).

At Node \( A \),

\[
\frac{V_2}{4} + \frac{V_2 - 4 \angle 0^\circ}{8} - I_2 = 0
\]

\[
\frac{3V_2}{8} - I_2 = \frac{1}{2} \quad (1)
\]

From the primary circuit,

\[
2I_1 + V_1 = 4 \quad (2)
\]

Since \( I_2/I_1 = 1/a \), \( V_2/V_1 = a \), and \( a = 2 \), it follows that \( I_1 = 2I_2 \) and \( V_1 = V_2/2 \). By substitution into (2) we get

\[
4I_2 + \frac{V_2}{2} = 4
\]

\[
I_2 + \frac{V_2}{8} = \frac{1}{2} \quad (3)
\]

Addition of (1) and (3) yields

\[
\frac{3V_2}{8} + \frac{V_2}{8} = \frac{1}{2} + \frac{1}{2} = 1
\]

from which \( V_2 = 3 \). Then,

\[
I_{4\Omega} = \frac{V_2}{4} = \frac{3}{4}
\]

and

\[
P_{\text{ave}} \, 4 \, \Omega = \frac{1}{2} \left( \frac{3}{4} \right)^2 4 = \frac{9}{8} \, w
\]
5. Because the dot on the secondary is at the lower end, \( a = -5 \). Then,

\[
aV_S = -5 \times 12 \angle 0^\circ = -60 \angle 0^\circ = 60 \angle 180^\circ
\]

\[
a^2 Z_S = 25(2 + j3) = 50 + j75 = 90.14 \angle 56.31^\circ \Omega
\]

\[
Z_{LOAD} = 100 - j75 = 125 \angle -36.87^\circ \Omega
\]

\[
I_2 = \frac{aV_S}{a^2 Z_S + Z_{LOAD}} = \frac{60 \angle 180^\circ}{50 + j75 + 100 - j75} = \frac{60 \angle 180^\circ}{150} = \frac{2}{5} \angle 180^\circ
\]

and

\[
V_2 = Z_{LOAD} \cdot I_2 = 125 \angle -36.87^\circ \times \frac{2}{5} \angle 180^\circ = 50 \angle 143.13^\circ \text{ V}
\]

6. From (8.102)

\[
Z_{in} = \frac{Z_{LOAD}}{a^2}
\]

Then,

\[
a^2 = \frac{Z_{LOAD}}{Z_{in}} = \frac{10000}{4} = 2500
\]

or

\[
a = 50
\]
Chapter 9

One- and Two-port Networks

This chapter begins with the general principles of one and two-port networks. The $z$, $y$, $h$, and $g$ parameters are defined. Several examples are presented to illustrate their use. It concludes with a discussion on reciprocal and symmetrical networks.

9.1 Introduction and Definitions

Generally, a network has two pairs of terminals; one pair is denoted as the input terminals, and the other as the output terminals. Such networks are very useful in the design of electronic systems, transmission and distribution systems, automatic control systems, communications systems, and others where electric energy or a signal enters the input terminals, it is modified by the network, and it exits through the output terminals.

A port is a pair of terminals in a network at which electric energy or a signal may enter or leave the network. A network that has only one pair of terminals is called a one-port network. In an one-port network, the current that enters one terminal must exit the network through the other terminal. Thus, in Figure 9.1, $i_{in} = i_{out}$

![Figure 9.1. One-port network](image)

Figures 9.2 and 9.3 show two examples of practical one-port networks.

![Figure 9.2. An example of an one-port network](image)
A two-port network has two pairs of terminals, that is, four terminals as shown in Figure 9.4 where \( i_1 = i_3 \) and \( i_2 = i_4 \).

![Diagram of a two-port network](image)

**Figure 9.4. Two-port network**

### 9.2 One-port Driving-point and Transfer Admittances

Let us consider an \( n \)-port network and write the mesh equations for this network in terms of the impedances \( Z \). We assume that the subscript of each current corresponds to the loop number and KVL is applied so that the sign of each \( Z_{ii} \) is positive. The sign of any \( Z_{ij} \) for \( i \neq j \) can be positive or negative depending on the reference directions of \( i_i \) and \( i_j \).

\[
\begin{align*}
Z_{11}i_1 + Z_{12}i_2 + Z_{13}i_3 + \ldots + Z_{1n}i_n &= v_1 \\
Z_{21}i_1 + Z_{22}i_2 + Z_{23}i_3 + \ldots + Z_{2n}i_n &= v_2 \\
&\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\quad\ Quad
One-port Driving-point and Transfer Admittances

\[ \Delta = \begin{bmatrix} Z_{11} & Z_{12} & Z_{13} & \cdots & Z_{1n} \\ Z_{21} & Z_{22} & Z_{23} & \cdots & Z_{2n} \\ Z_{31} & Z_{32} & Z_{33} & \cdots & Z_{3n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ Z_{n1} & Z_{n2} & Z_{n3} & \cdots & Z_{nn} \end{bmatrix} \]  \hspace{1cm} (9.3)

\[ D_1 = \begin{bmatrix} V_1 & Z_{12} & Z_{13} & \cdots & Z_{1n} \\ V_2 & Z_{22} & Z_{23} & \cdots & Z_{2n} \\ V_3 & Z_{32} & Z_{33} & \cdots & Z_{3n} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ V_n & Z_{n2} & Z_{n3} & \cdots & Z_{nn} \end{bmatrix} \]  \hspace{1cm} (9.4)

Next, we recall that the value of the determinant of a matrix \( A \) is the sum of the products obtained by multiplying each element of any row or column by its cofactor*. The cofactor, with the proper sign, is the matrix that remains when both the row and the column containing the element are eliminated. The sign is plus (+) when the sum of the subscripts is even, and it is minus (−) when it is odd. Mathematically, if the cofactor of the element \( a_{qr} \) is denoted as \( A_{qr} \), then

\[ A_{qr} = (-1)^{q+r} M_{qr} \]  \hspace{1cm} (9.5)

where \( M_{qr} \) is the minor of the element \( a_{qr} \). We recall also that the minor is the cofactor without a sign.

**Example 9.1**

Compute the determinant of \( A \) from the elements of the first row and their cofactors given that

\[ A = \begin{bmatrix} 1 & 2 & -3 \\ 2 & -4 & 2 \\ -1 & 2 & -6 \end{bmatrix} \]

Solution:

\[ \text{det}A = 1 \begin{bmatrix} -4 & 2 \\ 2 & -6 \end{bmatrix} - 2 \begin{bmatrix} 2 & 2 \\ -1 & -6 \end{bmatrix} - 3 \begin{bmatrix} 2 & -4 \\ -1 & 2 \end{bmatrix} = 1 \times 20 - 2 \times (-10) - 3 \times 0 = 40 \]

* A detailed discussion on cofactors is included in Appendix C.
Using the cofactor concept, and denoting the cofactor of the element \( a_{ij} \) as \( C_{ij} \), we find that the cofactors of \( Z_{11} \), \( Z_{12} \), and \( Z_{21} \) of (9.1) are respectively,

\[
C_{11} = \begin{bmatrix}
Z_{22} & Z_{23} & \cdots & Z_{2n} \\
Z_{32} & Z_{33} & \cdots & Z_{3n} \\
\vdots & \vdots & \ddots & \vdots \\
Z_{n2} & Z_{n3} & \cdots & Z_{nn}
\end{bmatrix}
\]  \hspace{1cm} (9.6)

\[
C_{12} = -\begin{bmatrix}
Z_{21} & Z_{23} & \cdots & Z_{2n} \\
Z_{31} & Z_{33} & \cdots & Z_{3n} \\
\vdots & \vdots & \ddots & \vdots \\
Z_{n1} & Z_{n3} & \cdots & Z_{nn}
\end{bmatrix}
\]  \hspace{1cm} (9.7)

\[
C_{21} = -\begin{bmatrix}
Z_{12} & Z_{13} & \cdots & Z_{1n} \\
Z_{32} & Z_{33} & \cdots & Z_{3n} \\
\vdots & \vdots & \ddots & \vdots \\
Z_{n2} & Z_{n3} & \cdots & Z_{nn}
\end{bmatrix}
\]  \hspace{1cm} (9.8)

Therefore, we can express (9.2) as

\[
i_1 = \frac{D_1}{\Delta} = \frac{C_{11}v_1}{\Delta} + \frac{C_{21}v_2}{\Delta} + \frac{C_{31}v_3}{\Delta} + \cdots + \frac{C_{n1}v_n}{\Delta} \hspace{1cm} (9.9)
\]

Also,

\[
i_2 = \frac{D_2}{\Delta} = \frac{C_{12}v_1}{\Delta} + \frac{C_{22}v_2}{\Delta} + \frac{C_{32}v_3}{\Delta} + \cdots + \frac{C_{n2}v_n}{\Delta} \hspace{1cm} (9.10)
\]

and the other currents \( i_3, i_4, \) and so on can be written in similar forms.

In network theory the \( y_{ij} \) parameters are defined as

\[
y_{11} = \frac{C_{11}}{\Delta} \quad y_{12} = \frac{C_{21}}{\Delta} \quad y_{13} = \frac{C_{31}}{\Delta} \quad \ldots 
\]

\[
y_{21} = \frac{C_{12}}{\Delta} \quad y_{22} = \frac{C_{22}}{\Delta} \quad y_{23} = \frac{C_{32}}{\Delta} \quad \ldots 
\]

and so on. By substitution of the \( y \) parameters into (9.9) and (9.10) we get:
One-port Driving-point and Transfer Admittances

\[ i_1 = y_{11}v_1 + y_{12}v_2 + y_{13}v_3 + \ldots + y_{1n}v_n \quad (9.13) \]

\[ i_2 = y_{21}v_1 + y_{22}v_2 + y_{23}v_3 + \ldots + y_{2n}v_n \quad (9.14) \]

If the subscripts of the \( y \)-parameters are alike, such as \( y_{11}, y_{22} \) and so on, they are referred to as driving-point admittances. If they are unlike, such as \( y_{12}, y_{21} \) and so on, they are referred to as transfer admittances.

If a network consists of only two loops such as in Figure 9.5,

\[ \begin{align*}
    i_1 &= y_{11}v_1 + y_{12}v_2 \\
    i_2 &= y_{21}v_1 + y_{22}v_2 
\end{align*} \]

From Figure 9.5 we observe that there is only one voltage source, \( v_1 \); there is no voltage source in Loop 2 and thus \( v_2 = 0 \). Then, (9.15) and (9.16) reduce to

\[ \begin{align*}
    i_1 &= y_{11}v_1 \\
    i_2 &= y_{21}v_1 
\end{align*} \]

Relation (9.17) reveals that the driving-point admittance \( y_{11} \) is the ratio \( i_1/v_1 \). That is, the driving-point admittance, as defined by (9.17), is the admittance seen by a voltage source that is present in the respective loop, in this case, Loop 1. Stated in other words, the driving-point admittance is the ratio of the current in a given loop to the voltage source in that loop when there are no voltage sources in any other loops of the network.

Transfer admittance is the ratio of the current in some other loop to the driving voltage source, in this case \( v_1 \). As indicated in (9.18), the transfer admittance \( y_{21} \) is the ratio of the current in Loop 2 to the voltage source in Loop 1.
Example 9.2

For the circuit of Figure 9.6, find the driving-point and transfer admittances and the current through each resistor.

The circuit is shown in Figure 9.6.

![Figure 9.6. Circuit for Example 9.2](image)

**Solution:**

We assign currents as shown in Figure 9.7.

![Figure 9.7. Loop equations for the circuit of Example 9.2](image)

The loop equations are

\[
\begin{align*}
10i_1 - 6i_2 &= 24 \\
-6i_1 + 18i_2 &= 0
\end{align*}
\]  \hspace{1cm} (9.19)

The driving-point admittance is found from (9.11), that is,

\[
y_{11} = \frac{C_{11}}{\Delta} \hspace{1cm} (9.20)
\]

and the transfer admittance from (9.12), that is,

\[
y_{21} = \frac{C_{12}}{\Delta} \hspace{1cm} (9.21)
\]

For this example,
The cofactor $C_{11}$ is obtained by inspection from the matrix of (9.22), that is, eliminating the first row and first column we are left with 18 and thus $C_{11} = 18$. Similarly, the cofactor $C_{12}$ is found by eliminating the first row and second column and changing the sign of $-6$. Then, $C_{12} = 6$. By substitution into (9.20) and (9.21), we obtain

$$y_{11} = \frac{C_{11}}{\Delta} = \frac{18}{144} = \frac{1}{8}$$  \hspace{1cm} (9.23)

and

$$y_{21} = \frac{C_{12}}{\Delta} = \frac{6}{144} = \frac{1}{24}$$  \hspace{1cm} (9.24)

Then, by substitution into (9.17) and (9.18) we get

$$i_1 = y_{11}v_1 = \frac{1}{8} \times 24 = 3 \text{ A}$$  \hspace{1cm} (9.25)

$$i_2 = y_{21}v_1 = \frac{1}{24} \times 24 = 1 \text{ A}$$  \hspace{1cm} (9.26)

Finally, the we observe that the current through the $4 \Omega$ resistor is $3 \text{ A}$, through the $12 \Omega$ is $1 \text{ A}$ and through the $6 \Omega$ is $i_1 - i_2 = 3 - 1 = 2 \text{ A}$

Of course, there are other simpler methods of computing these currents. However, the intent here was to illustrate how the driving-point and transfer admittances are applied. These allow easy computation for complicated network problems.

### 9.3 One-port Driving-point and Transfer Impedances

Now, let us consider an $n$–port network and write the nodal equations for this network in terms of the admittances $Y$. We assume that the subscript of each current corresponds to the loop number and KVL is applied so that the sign of each $Y_{ii}$ is positive. The sign of any $Y_{ij}$ for $i \neq j$ can be positive or negative depending on the reference polarities of $v_i$ and $v_j$.

\[
\begin{align*}
Y_{11}v_1 + Y_{12}v_2 + Y_{13}v_3 + \ldots + Y_{1n}v_n &= i_1 \\
Y_{21}v_1 + Y_{22}v_2 + Y_{23}v_3 + \ldots + Y_{2n}v_n &= i_2 \\
&\hspace{1cm}\vdots \\
Y_{n1}v_1 + Y_{n2}v_2 + Y_{n3}v_3 + \ldots + Y_{nn}v_n &= i_n
\end{align*}
\]  \hspace{1cm} (9.27)
In (9.27), each voltage can be found by Cramer’s rule. For instance, the voltage \( v_1 \) is found by

\[
v_1 = \frac{D_1}{\Delta}
\]

(9.28)

where

\[
\Delta = \begin{vmatrix}
Y_{11} & Y_{12} & Y_{13} & \cdots & Y_{1n} \\
Y_{21} & Y_{22} & Y_{23} & \cdots & Y_{2n} \\
Y_{31} & Y_{32} & Y_{33} & \cdots & Y_{3n} \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
Y_{n1} & Y_{n2} & Y_{n3} & \cdots & Y_{nn}
\end{vmatrix}
\]

(9.29)

\[
D_1 = \begin{vmatrix}
V_1 & Y_{12} & Y_{13} & \cdots & Y_{1n} \\
V_2 & Y_{22} & Y_{23} & \cdots & Y_{2n} \\
V_3 & Y_{32} & Y_{33} & \cdots & Y_{3n} \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
V_n & Y_{n2} & Y_{n3} & \cdots & Y_{nn}
\end{vmatrix}
\]

(9.30)

As in the previous section, we find that the nodal equations of (9.27) can be expressed as

\[
v_1 = z_{11}i_1 + z_{12}i_2 + z_{13}i_3 + \cdots + z_{1n}i_n
\]

(9.31)

\[
v_2 = z_{21}i_1 + z_{22}i_2 + z_{23}i_3 + \cdots + z_{2n}i_n
\]

(9.32)

\[
v_3 = z_{31}i_1 + z_{32}i_2 + z_{33}i_3 + \cdots + z_{3n}i_n
\]

(9.33)

and so on, where

\[
z_{11} = \frac{C_{11}}{\Delta}, \quad z_{12} = \frac{C_{21}}{\Delta}, \quad z_{13} = \frac{C_{31}}{\Delta}, \quad \ldots
\]

(9.34)

\[
z_{21} = \frac{C_{12}}{\Delta}, \quad z_{22} = \frac{C_{22}}{\Delta}, \quad z_{23} = \frac{C_{32}}{\Delta}, \quad \ldots
\]

(9.35)

\[
z_{31} = \frac{C_{13}}{\Delta}, \quad z_{32} = \frac{C_{23}}{\Delta}, \quad z_{33} = \frac{C_{33}}{\Delta}, \quad \ldots
\]

(9.36)

and so on. The matrices \( C_{ij} \) represent the cofactors as in the previous section.
The coefficients of (9.31), (9.32), and (9.33) with like subscripts are referred to as *driving-point impedances*. Thus, $z_{11}$, $z_{22}$ and so on, are driving-point impedances. The remaining coefficients with unlike subscripts, such as $z_{12}$, $z_{21}$ and so on, are called *transfer impedances*.

To understand the meaning of the driving-point and transfer impedances, we examine the network of Figure 9.8 where 0 is the reference node and nodes 1 and 2 are independent nodes. The driving point impedance is the ratio of the voltage across the nodes 1 and 0 to the current that flows through the branch between these nodes. In other words,

$$z_{11} = \frac{v_1}{i_1}$$  \hspace{1cm} (9.37)

![Figure 9.8. Circuit to illustrate the definitions of driving-point and transfer impedances.](image)

The transfer impedance between nodes 2 and 1 is the ratio of the voltage $v_2$ to the current at node 1 when there are no other current (or voltage) sources in the network. That is,

$$z_{21} = \frac{v_2}{i_1}$$  \hspace{1cm} (9.38)

**Example 9.3**

For the network of Figure 9.9, compute the driving-point and transfer impedances and the voltages across each conductance in terms of the current source.

![Figure 9.9. Network for Example 9.3.](image)
Solution:

We assign nodes 0, 1, 2, and 3 as shown in Figure 9.10.

![Node assignment for network of Example 9.3](image)

Figure 9.10. Node assignment for network of Example 9.3

The nodal equations are

\[
10v_1 + 2(v_1 - v_2) + I(v_1 - v_3) = i_1 \\
2(v_2 - v_1) + I(v_2 - v_3) + Iv_2 = 0 \\
I(v_3 - v_1) + I(v_3 - v_2) + Iv_3 = 0
\]  

(9.39)

Simplifying and rearranging we get:

\[
13v_1 - 2v_2 - v_3 = i_1 \\
-2v_1 + 4v_2 - v_3 = 0 \\
-v_1 - v_2 + 3v_3 = 0
\]  

(9.40)

The driving-point impedance \( z_{11} \) is found from (9.34), that is,

\[ z_{11} = \frac{C_{11}}{\Delta} \]  

(9.41)

and the transfer impedances \( z_{21} \) and \( z_{31} \) from (9.35) and (9.36), that is,

\[ z_{21} = \frac{C_{12}}{\Delta} \]  

(9.42)

\[ z_{31} = \frac{C_{13}}{\Delta} \]  

(9.43)

For this example,
The cofactor $C_{11}$ is

$$C_{11} = \begin{bmatrix} 4 & -1 \\ -1 & 3 \end{bmatrix} = 12 - 1 = 11$$  \tag{9.45}$$

Similarly, the cofactors $C_{12}$ and $C_{13}$ are

$$C_{12} = -\begin{bmatrix} -2 & -1 \\ -1 & 3 \end{bmatrix} = -(6 - 1) = 7$$  \tag{9.46}$$

and

$$C_{13} = \begin{bmatrix} -2 & 4 \\ -1 & -1 \end{bmatrix} = 2 + 4 = 6$$  \tag{9.47}$$

By substitution into (9.41), (9.42), and (9.43), we obtain

$$z_{11} = \frac{C_{11}}{\Delta} = \frac{11}{123}$$  \tag{9.48}$$

$$z_{21} = \frac{C_{12}}{\Delta} = \frac{7}{123}$$  \tag{9.49}$$

$$z_{31} = \frac{C_{13}}{\Delta} = \frac{6}{123}$$  \tag{9.50}$$

Then, by substitution into (9.31), (9.32), and (9.33) we get:

$$v_1 = z_{11}i_1 + z_{12}i_2 + z_{13}i_3 = \frac{11}{123}i_1$$  \tag{9.51}$$

$$v_2 = z_{21}i_1 + z_{22}i_2 + z_{23}i_3 = \frac{7}{123}i_1$$  \tag{9.52}$$

$$v_3 = z_{31}i_1 + z_{32}i_2 + z_{33}i_3 = \frac{6}{123}i_1$$  \tag{9.53}$$

Of course, there are other simpler methods of computing these voltages. However, the intent here was to illustrate how the driving-point and transfer impedances are applied. These allow easy computation for complicated network problems.
9.4 Two-Port Networks

Figure 9.11 shows a two-port network with external voltages and currents specified.

![Diagram of a two-port network](https://via.placeholder.com/150)

Here, we assume that \( i_1 = i_3 \) and \( i_2 = i_4 \). We also assume that \( i_1 \) and \( i_2 \) are obtained by the superposition of the currents produced by both \( v_1 \) and \( v_2 \).

Now, we will define the \( y \), \( z \), \( h \), and \( g \) parameters.

9.4.1 The \( y \) Parameters

The two-port network of Figure 9.11 can be described by the following set of equations.

\[
i_1 = y_{11}v_1 + y_{12}v_2 \quad (9.54)
\]

\[
i_2 = y_{21}v_1 + y_{22}v_2 \quad (9.55)
\]

In two-port network theory, the \( y \) coefficients are referred to as the \( y \) parameters.

Let us assume that \( v_2 \) is shorted, that is, \( v_2 = 0 \). Then, (9.54) reduces to

\[
i_1 = y_{11}v_1 \quad (9.56)
\]

or

\[
y_{11} = \frac{i_1}{v_1} \quad (9.57)
\]

and \( y_{11} \) is referred to as the short circuit input admittance at the left port when the right port of Figure 9.11 is short-circuited.

Let us again consider (9.54), that is,

\[
i_1 = y_{11}v_1 + y_{12}v_2 \quad (9.58)
\]

This time we assume that \( v_1 \) is shorted, i.e., \( v_1 = 0 \). Then, (9.58) reduces to

\[
i_1 = y_{12}v_2 \quad (9.59)
\]
Two-Port Networks

or

\[ y_{12} = \frac{i_1}{v_2} \]  
(9.60)

and \( y_{12} \) is referred to as the short circuit transfer admittance when the left port of Figure 9.11 is short-circuited. It represents the transmission from the right to the left port. For instance, in amplifiers where the left port is considered to be the input port and the right to be the output, the parameter \( y_{12} \) represents the internal feedback inside the network.

Similar expressions are obtained when we consider the equation for \( i_2 \), that is,

\[ i_2 = y_{21}v_1 + y_{22}v_2 \]  
(9.61)

In an amplifier, the parameter \( y_{21} \) is also referred to as the short circuit transfer admittance and represents transmission from the left (input) port to the right (output) port. It is a measure of the so-called forward gain.

The parameter \( y_{22} \) is called the short circuit output admittance.

The \( y \) parameters and the conditions under which they are computed are shown in Figures 9.12 through 9.16.

![Figure 9.12. The y parameters for \( v_1 \neq 0 \) and \( v_2 \neq 0 \)](image)

![Figure 9.13. Network for the definition of the \( y_{11} \) parameter](image)
Example 9.4

For the network of Figure 9.17, find the \( y \) parameters.

Solution:

a. The short circuit input admittance \( y_{11} \) is found from the network of Figure 9.18 where we have assumed that \( v_i = 1 \) V and the resistances, for convenience, have been replaced with conductances in mhos.
Two-Port Networks

We observe that the $0.05 \, \Omega^{-1}$ conductance is shorted out and thus the current $i_1$ is the sum of the currents through the $0.2 \, \Omega^{-1}$ and $0.1 \, \Omega^{-1}$ conductances. Then,

$$i_1 = 0.2v_1 + 0.1v_1 = 0.2 \times 1 + 0.1 \times 1 = 0.3 \, A$$

and thus the short circuit input admittance is

$$y_{11} = i_1/v_1 = 0.3/1 = 0.3 \, \Omega^{-1} \quad (9.62)$$

b. The short circuit transfer admittance $y_{12}$ when the left port is short-circuited, is found from the network of Figure 9.19.

We observe that the $0.2 \, \Omega^{-1}$ conductance is shorted out and thus the $0.1 \, \Omega^{-1}$ conductance is in parallel with the $0.05 \, \Omega^{-1}$ conductance. The current $i_1$, with a minus (−) sign, now flows through the $0.1 \, \Omega^{-1}$ conductance. Then,

$$i_1 = -0.1v_2 = -0.1 \times 1 = -0.1 \, A$$

and

$$y_{12} = i_1/v_2 = -0.1/1 = -0.1 \, \Omega^{-1} \quad (9.63)$$
c. The short circuit transfer admittance $y_{21}$ when the right port is short-circuited, is found from the network of Figure 9.20.

![Figure 9.20. Network for computing $y_{21}$](image)

We observe that the $0.05 \, \Omega^{-1}$ conductance is shorted out and thus the $0.1 \, \Omega^{-1}$ conductance is in parallel with the $0.2 \, \Omega^{-1}$ conductance. The current $i_2$, with a minus (−) sign, now flows through the $0.1 \, \Omega^{-1}$ conductance. Then,

$$i_2 = -0.1v_1 = -0.1 \times 1 = -0.1 \, A$$

and

$$y_{21} = i_2/v_1 = -0.1/1 = -0.1 \, \Omega^{-1} \quad (9.64)$$

d. The short circuit output admittance $y_{22}$ at the right port when the left port is short-circuited, is found from the network of 9.21.

![Figure 9.21. Network for computing $y_{22}$](image)

We observe that the $0.2 \, \Omega^{-1}$ conductance is shorted out and thus the current $i_2$ is the sum of the currents through the $0.05 \, \Omega^{-1}$ and $0.1 \, \Omega^{-1}$ conductances. Then,

$$i_2 = 0.05v_2 + 0.1v_2 = 0.05 \times 1 + 0.1 \times 1 = 0.15 \, A$$
Therefore, the two-port network of Figure 9.10 can be described by the following set of equations.

\[ i_1 = y_{11}v_1 + y_{12}v_2 = 0.3v_1 - 0.1v_2 \]  
\[ i_2 = y_{21}v_1 + y_{22}v_2 = -0.1v_1 + 0.3v_2 \]  

(9.66)

Note:

In Example 9.4, we found that the short circuit transfer admittances are equal, that is,

\[ y_{21} = y_{12} = -0.1 \]  

(9.67)

This is not just a coincidence; this is true whenever a two-port network is reciprocal (or bilateral). A network is reciprocal if the reciprocity theorem is satisfied. This theorem states that:

*If a voltage applied in one branch of a linear, two-port passive network produces a certain current in any other branch of this network, the same voltage applied in the second branch will produce the same current in the first branch.*

The reverse is also true, that is, if current applied at one node produces a certain voltage at another, the same current at the second node will produce the same voltage at the first. An example is given at the end of this chapter.

Obviously, if we know that the two-port network is reciprocal, only three computations are required to find the \( y \) parameters.

If in a reciprocal two-port network its ports can be interchanged without affecting the terminal voltages and currents, the network is said to be also symmetric. In a symmetric two-port network,

\[ y_{22} = y_{11} \]  
\[ y_{21} = y_{12} \]  

(9.68)

The network of Figure 9.17 is not symmetric since \( y_{22} \neq y_{11} \)

We will present examples of reciprocal and symmetric two-port networks at the last section of this chapter.

The following example illustrates the applicability of two-port network analysis in more complicated networks.

**Example 9.5**

For the network of Figure 9.22, compute \( v_1, v_2, i_1, \) and \( i_2 \).
Solution:

We recognize the portion of the network enclosed in the dotted square, shown in Figure 9.23, as that of the previous example.

For the network of Figure 9.23, at Node 1,

\[ i_1 = 15 - v_1 / 10 \]  \tag{9.69}

and at Node 2,

\[ i_2 = -v_2 / 4 \]  \tag{9.70}

By substitution of (9.69) and (9.70) into (9.66), we get:

\[ i_1 = y_{11} v_1 + y_{12} v_2 = 0.3 v_1 - 0.1 v_2 = 15 - v_1 / 10 \]  \tag{9.71}

\[ i_2 = y_{21} v_1 + y_{22} v_2 = -0.1 v_1 + 0.3 v_2 = -v_2 / 4 \]

or

\[ 0.4 v_1 - 0.1 v_2 = 15 \]

\[ -0.1 v_1 + 0.4 v_2 = 0 \]  \tag{9.72}

We will use MATLAB to solve the equations of (9.72) to become more familiar with it.
Two-Port Networks

syms v1 v2; [v1 v2]=solve(0.4*v1-0.1*v2-15, -0.1*v1+0.4*v2)

\[ v1 = 40 \]
\[ v2 = 10 \]

and thus

\[ v_1 = 40 \text{ V} \]
\[ v_2 = 10 \text{ V} \] (9.73)

The currents \( i_1 \) and \( i_2 \) are found from (9.69) and (9.70).

\[ i_1 = 15 - 40/10 = 11 \text{ A} \] (9.74)
\[ i_2 = -10/4 = -2.5 \text{ A} \]

9.4.2 The \( z \) parameters

A two-port network such as that of Figure 9.24 can also be described by the following set of equations.

\[ v_1 = z_{11}i_1 + z_{12}i_2 \]
\[ v_2 = z_{21}i_1 + z_{22}i_2 \]

Figure 9.24. The \( z \) parameters for \( i_1 \neq 0 \) and \( i_2 \neq 0 \)

\[ v_1 = z_{11}i_1 + z_{12}i_2 \] (9.75)
\[ v_2 = z_{21}i_1 + z_{22}i_2 \] (9.76)

In two-port network theory, the \( z_{ij} \) coefficients are referred to as the \( z \) parameters or as open circuit impedance parameters.

Let us assume that \( v_2 \) is open, that is, \( i_2 = 0 \) as shown in Figure 9.25.

\[ z_{11} = \frac{v_1}{i_1} \]

Figure 9.25. Network for the definition of the \( z_{11} \) parameter
Then, (9.75) reduces to

\[ v_1 = z_{11}i_1 \quad (9.77) \]

or

\[ z_{11} = \frac{v_1}{i_1} \quad (9.78) \]

and this is the open circuit input impedance when the right port of Figure 9.25 is open.

Let us again consider (9.75), that is,

\[ v_1 = z_{11}i_1 + z_{12}i_2 \quad (9.79) \]

This time we assume that the terminal at \( v_1 \) is open, i.e., \( i_1 = 0 \) as shown in Figure 9.26.

![Figure 9.26. Network for the definition of the \( z_{12} \) parameter](image)

Then, (9.75) reduces to

\[ v_1 = z_{12}i_2 \quad (9.80) \]

or

\[ z_{12} = \frac{v_1}{i_2} \quad (9.81) \]

and this is the open circuit transfer impedance when the left port is open as shown in Figure 9.26.

Similar expressions are obtained when we consider the equation for \( v_2 \), that is,

\[ v_2 = z_{21}i_1 + z_{22}i_2 \quad (9.82) \]

Let us assume that \( v_2 \) is open, that is, \( i_2 = 0 \) as shown in Figure 9.27.

Then, (9.82) reduces to

\[ v_2 = z_{21}i_1 \quad (9.83) \]
The parameter $z_{21}$ is referred to as **open circuit transfer impedance** when the right port is open as shown in Figure 9.27.

Finally, let us assume that the terminal at $v_I$ is open, i.e., $i_I = 0$ as shown in Figure 9.28.

Then, (9.82) reduces to

$$v_2 = z_{22} i_2 \quad (9.85)$$

or

$$z_{22} = \frac{v_2}{i_2} \quad (9.86)$$

The parameter $z_{22}$ is called the **open circuit output impedance**.

We observe that the $z$ parameters definitions are similar to those of the $y$ parameters if we substitute voltages for currents and currents for voltages.

**Example 9.6**

For the network of Figure 9.29, find the $z$ parameters.
Solution:

a. The open circuit input impedance $z_{11}$ is found from the network of Figure 9.30 where we have assumed that $i_1 = 1 \text{ A}$.

![Network for computing $z_{11}$ for the network of Figure 9.29](image)

We observe that the 20 $\Omega$ resistor is in parallel with the series combination of the 5 $\Omega$ and 15 $\Omega$ resistors. Then, by the current division expression, the current through the 20 $\Omega$ resistor is 0.5 A and the voltage across that resistor is

$$v_1 = 20 \times 0.5 = 10 \text{ V}$$

Therefore, the open circuit input impedance $z_{11}$ is

$$z_{11} = v_1/i_1 = 10/1 = 10 \Omega$$

(9.87)

b. The open circuit transfer impedance $z_{12}$ is found from the network of Figure 9.31.

We observe that the 15 $\Omega$ resistance is in parallel with the series combination of the 5 $\Omega$ and 20 $\Omega$ resistances. Then, the current through the 20 $\Omega$ resistance is

$$i_{20\Omega} = \frac{15}{15 + 5 + 20}i_2 = \frac{15}{40} \times 1 = 3/8 \text{ A}$$
and the voltage across this resistor is

\[
\frac{3}{8} \times 20 = \frac{60}{8} = 15/2 \text{ V}
\]

Therefore, the open circuit transfer impedance \( z_{12} \) is

\[
z_{12} = \frac{v_1}{i_2} = \frac{15/2}{1} = 7.5 \text{ }\Omega \tag{9.88}
\]

c. The open circuit transfer impedance \( z_{21} \) is found from the network of Figure 9.32.

In Figure 9.32 the current that flows through the 15 \( \Omega \) resistor is

\[
i_{15 \Omega} = \frac{20}{20 + 5 + 15} i_1 = \frac{20}{40} \times 1 = 1/2 \text{ A}
\]

and the voltage across this resistor is

\[
v_2 = \frac{1}{2} \times 15 = 15/2 \text{ V}
\]

Therefore, the open circuit transfer impedance \( z_{21} \) is

\[
z_{21} = \frac{v_2} {i_1} = \frac{15/2}{1} = 7.5 \text{ }\Omega \tag{9.89}
\]
We observe that
\[
z_{21} = z_{12}
\]  
(9.90)

d. The open circuit output impedance \( z_{22} \) is found from the network of Figure 9.33.

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{network.png}
\caption{Network for computing \( z_{22} \) for the network of Figure 9.29}
\end{figure}

We observe that the 15 Ω resistance is in parallel with the series combination of the 5 Ω and 20 Ω resistances. Then, the current through the 15 Ω resistance is
\[
i_{15} = \frac{20 + 5}{20 + 5 + 15} i_2 = \frac{25}{40} \times 1 = 5/8 \text{ A}
\]

and the voltage across that resistor is
\[
\frac{5}{8} \times 15 = 75/8 \text{ V}
\]

Therefore, the open circuit output impedance \( z_{22} \) is
\[
\frac{v_1}{i_2} = \frac{75/8}{1} = 75/8 \text{ Ω}
\]

(9.91)

9.4.3 The \( h \) Parameters

A two-port network can also be described by the set of equations
\[
v_1 = h_{11} i_1 + h_{12} v_2
\]
(9.92)
\[
i_2 = h_{21} i_1 + h_{22} v_2
\]
(9.93)

as shown in Figure 9.34.

The \( h \) parameters represent an impedance, a voltage gain, a current gain, and an admittance. For this reason they are called hybrid (different) parameters.

Let us assume that \( v_2 = 0 \) as shown in Figure 9.35.
Then, (9.92) reduces to

\[ v_1 = h_{11} i_1 \]

(9.94)

or

\[ h_{11} = \frac{v_1}{i_1} \]

(9.95)

Therefore, the parameter \( h_{11} \) represents the input impedance of a two-port network.

Let us assume that \( i_1 = 0 \) as shown in Figure 9.36.

Then, (9.92) reduces to

\[ v_1 = h_{12} v_2 \]

(9.96)

or

\[ h_{12} = \frac{v_1}{v_2} \]

(9.97)
Therefore, in a two-port network the parameter \( h_{12} \) represents a voltage gain (or loss).

Let us assume that \( v_2 = 0 \) as shown in Figure 9.37.

\[ h_{21} = \left. \frac{i_2}{i_1} \right|_{v_2 = 0} \]

*Figure 9.37. Network for computing \( h_{21} \) for the network of Figure 9.34*

Then, (9.93) reduces to

\[ i_2 = h_{21} i_1 \]

or

\[ h_{21} = \left. \frac{i_2}{i_1} \right|_{v_2 = 0} \]

Therefore, in a two-port network the parameter \( h_{21} \) represents a current gain (or loss).

Finally, let us assume that the terminal at \( v_1 \) is open, i.e., \( i_1 = 0 \) as shown in Figure 9.38.

\[ h_{22} = \left. \frac{i_2}{v_2} \right|_{i_1 = 0} \]

*Figure 9.38. Network for computing \( h_{22} \) for the network of Figure 9.34*

Then, (9.93) reduces to

\[ i_2 = h_{22} v_2 \]

or

\[ h_{22} = \left. \frac{i_2}{v_2} \right|_{i_1 = 0} \]

Therefore, in a two-port network the parameter \( h_{22} \) represents an output admittance.

**Example 9.7**

For the network of Figure 9.39, find the \( h \) parameters.
Solution:

a. The short circuit input impedance $h_{11}$ is found from the network of Figure 9.40 where we have assumed that $i_1 = 1 \, \text{A}$.

![Network for computing $h_{11}$ for the network of Figure 9.39](image)

From the network of Figure 9.40 we observe that the 4 $\Omega$ and 6 $\Omega$ resistors are in parallel yielding an equivalent resistance of 2.4 $\Omega$ in series with the 1 $\Omega$ resistor. Then, the voltage across the current source is

$$v_1 = i_1 \times (1 + 2.4) = 3.4 \, \text{V}$$

Therefore, the short circuit input impedance $h_{11}$ is

$$h_{11} = \frac{v_1}{i_1} = \frac{3.4}{1} = 3.4 \, \Omega$$ \hspace{1cm} (9.98)

b. The voltage gain $h_{12}$ is found from the network of Figure 9.41.

Since no current flows through the 1 $\Omega$ resistor, the voltage $v_1$ is the voltage across the 4 $\Omega$ resistor. Then, by the voltage division expression,

$$v_1 = \frac{4}{6+4}v_2 = \frac{4}{10} \times 1 = 0.4 \, \text{V}$$
Therefore, the voltage gain $h_{12}$ is the dimensionless number

$$h_{12} = \frac{v_I}{v_2} = \frac{0.4}{I} = 0.4$$

\hspace{1cm} (9.99)

c. The current gain $h_{21}$ is found from the network of Figure 9.42.

We observe that the 4 Ω and 6 Ω resistors are in parallel yielding an equivalent resistance of 2.4 Ω. Then, the voltage across the 2.4 Ω parallel combination is

$$v_{2.4Ω} = 2.4 \times I = 2.4 \text{ V}$$

The current $i_2$ is the current through the 6 Ω resistor. Thus,

$$i_2 = \frac{2.4}{6} = -0.4 \text{ A}$$

Therefore, the current gain $h_{21}$ is the dimensionless number

$$h_{21} = \frac{i_2}{I} = \frac{-0.4}{I} = -0.4$$

We observe that
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\[ h_{21} = -h_{12} \]  

(9.100)

and this is a consequence of the fact that the given network is reciprocal.

d. The open circuit admittance \( h_{22} \) is found from the network of Figure 9.43.

\[ i_1 = 0 \quad v_1 \]
\[ i_2 \]
\[ v_2 = 1 \text{ V} \]

**Figure 9.43. Network for computing \( h_{22} \) for the network of Figure 9.39.**

Since no current flows through the \( 1 \Omega \) resistor, the current \( i_2 \) is found by Ohm’s law as

\[ i_2 = \frac{v_2}{6 + 4} = \frac{1}{10} = 0.1 \text{ A} \]

Therefore, the open circuit admittance \( h_{22} \) is

\[ h_{22} = \frac{i_2}{v_2} = \frac{0.1}{1} = 0.1 \Omega^{-1} \]  

(9.101)

**Note:**

The \( h \) parameters and the \( g \) parameters (to be discussed next), are used extensively in networks consisting of transistors *, and feedback networks. The \( h \) parameters are best suited with series-parallel feedback networks, whereas the \( g \) parameters are preferred in parallel-series amplifiers.

**9.4.4 The \( g \) Parameters**

A two-port network can also be described by the set of equations

\[ i_1 = g_{11}v_1 + g_{12}i_2 \]  

(9.102)

\[ v_2 = g_{21}v_1 + g_{22}i_2 \]  

(9.103)

as shown in Figure 9.44.

* Transistors are three-terminal devices. However, they can be represented as large-signal equivalent two-port networks circuits and also as small-signal equivalent two-port networks where linearity can be applied.
The \( g \) parameters, also known as inverse hybrid parameters, represent an admittance, a current gain, a voltage gain and an impedance.

Let us assume that \( i_2 = 0 \) as shown in Figure 9.45.

Then, (9.102) reduces to

\[
i_1 = g_{11} v_1 \quad \text{(9.104)}
\]

or

\[
g_{11} = \left. \frac{i_1}{v_1} \right|_{i_2 = 0} \quad \text{(9.105)}
\]

Therefore, the parameter \( g_{11} \) represents the input admittance of a two-port network.

Let us assume that \( v_1 = 0 \) as shown in Figure 9.46.

Then, (9.102) reduces to

\[
i_1 = g_{12} i_2 \quad \text{(9.106)}
\]
Therefore, in a two-port network the parameter $g_{12}$ represents a current gain (or loss).

Let us assume that $i_2 = 0$ as shown in Figure 9.47.

Then, (9.103) reduces to

$$v_2 = g_{21} v_1$$

or

$$g_{21} = \frac{v_2}{i_j}$$

Therefore, in a two-port network the parameter $g_{21}$ represents a voltage gain (or loss).

Finally, let us assume that $v_1$ is shorted, i.e., $v_1 = 0$ as shown in Figure 9.48.

Then, (9.103) reduces to

$$v_2 = g_{22} i_2$$

or

$$g_{22} = \frac{v_2}{i_2}$$
Thus, in a two-port network the parameter $g_{22}$ represents the output impedance of that network.

**Example 9.8**

For the network of Figure 9.49, find the $g$ parameters.

**Solution:**

a. The open circuit input admittance $g_{11}$ is found from the network of Figure 9.50 where we have assumed that $v_1 = 1$ V.

There is no current through the $4 \, \Omega$ resistor and thus by Ohm’s law,

$$i_1 = \frac{v_1}{1 + 12} = \frac{1}{13} \, A$$

Therefore, the open circuit input admittance $g_{11}$ is

$$g_{11} = \frac{i_1}{v_1} = \frac{1/13}{1} = \frac{1}{13} \, \Omega^{-1} \quad (9.112)$$

b. The current gain $g_{12}$ is found from the network of Figure 9.51.

By the current division expression, the current through the $1 \, \Omega$ resistor is

$$i_1 = -\frac{12}{12 + 1} i_2 = -\frac{12}{13} \times 1 = -\frac{12}{13} \, A$$
Therefore, the current gain \( g_{12} \) is the dimensionless number

\[
g_{12} = \frac{i_1}{i_2} = -\frac{12/13}{1} = -\frac{12}{13}
\]  

(9.113)

c. The voltage gain \( g_{21} \) is found from the network of Figure 9.52.

![Figure 9.52. Network for computing \( g_{21} \) for the network of Figure 9.49.]

Since there is no current through the \( 4 \) \( \Omega \) resistor, the voltage \( v_2 \) is the voltage across the \( 12 \) \( \Omega \) resistor. Then, by the voltage division expression,

\[
v_2 = \frac{12}{1+12} \times 1 = \frac{12}{13} \text{ V}
\]

Therefore, the voltage gain \( g_{21} \) is the dimensionless number

\[
g_{21} = \frac{v_2}{v_1} = \frac{12/13}{1} = \frac{12}{13}
\]

We observe that

\[
g_{21} = -g_{12}
\]  

(9.114)

and this is a consequence of the fact that the given network is reciprocal.

d. The short circuit output impedance \( g_{22} \) is found from the network of Figure 9.53.
Figure 9.53. Network for computing $g_{22}$ for the network of Figure 9.49.

The voltage $v_2$ is the sum of the voltages across the $4 \ \Omega$ resistor and the voltage across the $12 \ \Omega$ resistor. By the current division expression, the current through the $12 \ \Omega$ resistor is

$$i_{12\Omega} = \frac{l}{l + 12} i_2 = \frac{1}{13} \times 1 = 1/13 \ \text{A}$$

(9.115)

Then,

$$v_{12\Omega} = \frac{1}{13} \times 12 = 12/13 \ \text{V}$$

and

$$v_2 = \frac{12}{13} + 4 = 64/13 \ \text{V}$$

Therefore, the short circuit output impedance $g_{22}$ is

$$g_{22} = \frac{v_2}{i_2} = \frac{64/13}{l} = 64/13 \ \Omega$$

(9.116)

### 9.5 Reciprocal Two-Port Networks

If any of the following relationships exist in a two-port network,

$$z_{21} = z_{12}$$
$$y_{21} = y_{12}$$
$$h_{21} = -h_{12}$$
$$g_{21} = -g_{12}$$

(9.117)

the network is said to be reciprocal.

If, in addition to (9.117), any of the following relationship exists
Reciprocal Two-Port Networks

\[ z_{22} = z_{11} \]
\[ y_{22} = y_{11} \]
\[ h_{11} h_{22} - h_{12} h_{21} = 1 \]
\[ g_{11} g_{22} - g_{12} g_{21} = 1 \]

(9.118)

the network is said to be symmetric.

Examples of reciprocal two-port networks are the **tee**, **π**, **bridged (lattice)**, and **bridged tee**. These are shown in Figure 9.54.

Examples of symmetric two-port networks are shown in Figure 9.55.

Let us review the reciprocity theorem and its consequences before we present an example. This theorem states that:

*If a voltage applied in one branch of a linear, two-port passive network produces a certain current in any other branch of this network, the same voltage applied in the second branch will produce the same current in the first branch.*

The reverse is also true, that is, if current applied at one node produces a certain voltage at another, the same current at the second node will produce the same voltage at the first.

It was also stated earlier that if we know that the two-port network is reciprocal, only three computations are required to find the \( y, z, h, \) and \( g \) parameters as shown in (9.117). Furthermore, if we know that the two-port network is symmetric, we only need to make only two computations as shown in (9.118).
Figure 9.55. Examples of symmetric two-port networks.

Example 9.9

In the two-port network of Figure 9.56, the voltage source $v_S$ connected at the left end of the network is set for 15 V, and all impedances are resistive with the values indicated. On the right side of the network is connected a DC ammeter denoted as $A$. Assume that the ammeter is ideal, that is, has no internal resistance.

a. Compute the ammeter reading.

b. Interchange the positions of the voltage source and recompute the ammeter reading.

Solution:

a. Perhaps the easiest method of solution is by nodal analysis since we only need to solve one equation.
The given network is redrawn as shown in Figure 9.57.

![Network Diagram](image)

Figure 9.57. Network for solution of Example 9.9 by nodal analysis

By KCL at node \( a \),

\[
\frac{V_{ab} - 15}{30} + \frac{V_{ab}}{60} + \frac{V_{ab}}{20} = 0
\]

or

\[
\frac{6}{60} V_{ab} = \frac{15}{30}
\]

or

\[ V_{ab} = 5 \text{ V} \]

The current through the ammeter is the sum of the currents \( I_{Z3} \) and \( I_{Z4} \). Thus, denoting the current through the ammeter as \( I_A \) we get:

\[
I_A = I_{Z3} + I_{Z4} = \frac{V_{ab}}{Z_3} + \frac{V}{Z_4} = \frac{5}{20} + \frac{15}{10} = 0.25 + 1.50 = 1.75 \text{ A}
\]  \hspace{1cm} (9.119)

b. With the voltage source and ammeter positions interchanged, the network is as shown in Figure 9.58.

![Network Diagram](image)

Figure 9.58. Network of Figure 9.57 with the voltage source and ammeter interchanged.
Applying KCL for the network of Figure 9.58, we get:

\[ \frac{V_{ab}}{30} + \frac{V_{ab}}{60} + \frac{V_{ab} - 15}{20} = 0 \]

or

\[ \frac{6}{60} V_{ab} = \frac{15}{20} \]

or

\[ V_{ab} = 7.5 \, \text{V} \]

The current through the ammeter this time is the sum of the currents \( I_{Z1} \) and \( I_{Z4} \). Thus, denoting the current through the ammeter as \( I_A \) we get:

\[ I_A = I_{Z1} + I_{Z4} = \frac{V_{ab}}{Z_1} + \frac{V}{Z_4} = \frac{7.5}{30} + \frac{15}{10} = 0.25 + 1.50 = 1.75 \, \text{A} \]  

(9.120)

We observe that (9.119) and (9.120) give the same value and thus we can say that the given network is reciprocal.

9.6 Summary

- A port is a pair of terminals in a network at which electric energy or a signal may enter or leave the network.
- A network that has only one pair a terminals is called a one-port network. In an one-port network, the current that enters one terminal must exit the network through the other terminal.
- A two-port network has two pairs of terminals, that is, four terminals.
- For an \( n \)-port network the \( y \) parameters are defined as

\[
\begin{align*}
    i_1 &= y_{11}v_1 + y_{12}v_2 + y_{13}v_3 + \ldots + y_{1n}v_n \\
    i_2 &= y_{21}v_1 + y_{22}v_2 + y_{23}v_3 + \ldots + y_{2n}v_n \\
    i_3 &= y_{31}v_1 + y_{32}v_2 + y_{33}v_3 + \ldots + y_{3n}v_n \\

e &\quad \ldots
\end{align*}
\]

and so on.
- If the subscripts of the \( y \)-parameters are alike, such as \( y_{11}, y_{22} \) and so on, they are referred to as driving-point admittances. If they are unlike, such as \( y_{12}, y_{21} \) and so on, they are referred to as transfer admittances.
- For a \( 2 \)-port network the \( y \) parameters are defined as
Summary

\[ i_1 = y_{11}v_1 + y_{12}v_2 \]
\[ i_2 = y_{21}v_1 + y_{22}v_2 \]

- In a 2-\textit{port} network where the right port is short-circuited, that is, when \( v_2 = 0 \), the \( y_{11} \) parameter is referred to as the short circuit input admittance. In other words,

\[ y_{11} = \left. \frac{i_1}{v_1} \right|_{v_2 = 0} \]

- In a 2-\textit{port} network where the left port is short-circuited, that is, when \( v_1 = 0 \), the \( y_{12} \) parameter is referred to as the short circuit transfer admittance. In other words,

\[ y_{12} = \left. \frac{i_1}{v_2} \right|_{v_1 = 0} \]

- In a 2-\textit{port} network where the right port is short-circuited, that is, when \( v_2 = 0 \), the \( y_{21} \) parameter is referred to as the short circuit transfer admittance. In other words,

\[ y_{21} = \left. \frac{i_2}{v_1} \right|_{v_2 = 0} \]

- In a 2-\textit{port} network where the left port is short-circuited, that is, when \( v_1 = 0 \), the \( y_{22} \) parameter is referred to as the short circuit output admittance. In other words,

\[ y_{22} = \left. \frac{i_2}{v_1} \right|_{v_2 = 0} \]

- For a \( n \)-\textit{port} network the \( z \) parameters are defined as

\[ v_1 = z_{11}i_1 + z_{12}i_2 + z_{13}i_3 + \ldots + z_{1n}i_n \]
\[ v_2 = z_{21}i_1 + z_{22}i_2 + z_{23}i_3 + \ldots + z_{2n}i_n \]
\[ v_3 = z_{31}i_1 + z_{32}i_2 + z_{33}i_3 + \ldots + z_{3n}i_n \]

and so on.

- If the subscripts of the \( z \)-parameters are alike, such as \( z_{11} \), \( z_{22} \) and so on, they are referred to as driving-point impedances. If they are unlike, such as \( z_{12} \), \( z_{21} \) and so on, they are referred to as transfer impedances.
Chapter 9 One- and Two-port Networks

- For a 2–port network the \( z \) parameters are defined as

\[
v_1 = z_{11}i_1 + z_{12}i_2
\]
\[
v_2 = z_{21}i_1 + z_{22}i_2
\]

- In a 2–port network where the right port is open, that is, when \( i_2 = 0 \), the \( z_{11} \) parameter is referred to as the open circuit input impedance. In other words,

\[
z_{11} = \left. \frac{v_1}{i_1} \right|_{i_2 = 0}
\]

- In a 2–port network where the left port is open, that is, when \( i_1 = 0 \), the \( z_{12} \) parameter is referred to as the open circuit transfer impedance. In other words,

\[
z_{12} = \left. \frac{v_1}{i_2} \right|_{i_1 = 0}
\]

- In a 2–port network where the right port is open, that is, when \( i_2 = 0 \), the \( z_{21} \) parameter is referred to as the open circuit transfer impedance. In other words,

\[
z_{21} = \left. \frac{v_2}{i_1} \right|_{i_2 = 0}
\]

- In a 2–port network where the left port is open, that is, when \( i_1 = 0 \), the \( z_{22} \) parameter is referred to as the open circuit output impedance. In other words,

\[
z_{22} = \left. \frac{v_2}{i_2} \right|_{i_1 = 0}
\]

- A two-port network can also be described in terms of the \( h \) parameters with the equations

\[
v_1 = h_{11}i_1 + h_{12}v_2
\]
\[
i_2 = h_{21}i_1 + h_{22}v_2
\]

- The \( h \) parameters represent an impedance, a voltage gain, a current gain, and an admittance. For this reason they are called hybrid (different) parameters.

- In a 2–port network where the right port is shorted, that is, when \( v_2 = 0 \), the \( h_{11} \) parameter represents the input impedance of the two-port network. In other words,
Summary

\[ h_{11} = \frac{v_1}{i_1} \bigg|_{v_2 = 0} \]

- In a 2\text{–}port network where the left port is open, that is, when \( i_1 = 0 \), the \( h_{12} \) parameter represents a voltage gain (or loss) in the two-port network. In other words,

\[ h_{12} = \frac{v_1}{v_2} \bigg|_{i_1 = 0} \]

- In a 2\text{–}port network where the right port is shorted, that is, when \( v_2 = 0 \), the \( h_{21} \) parameter represents a current gain (or loss). In other words,

\[ h_{21} = \frac{i_2}{i_1} \bigg|_{v_2 = 0} \]

- In a 2\text{–}port network where the left port is open, that is, when \( i_1 = 0 \), the \( h_{22} \) parameter represents an output admittance. In other words,

\[ h_{22} = \frac{i_2}{v_2} \bigg|_{i_1 = 0} \]

- A two-port network can also be described in terms of the \( g \) parameters with the equations

\[ i_1 = g_{11}v_1 + g_{12}i_2 \]

\[ v_2 = g_{21}v_1 + g_{22}i_2 \]

- The \( g \) parameters, also known as inverse hybrid parameters, represent an admittance, a current gain, a voltage gain and an impedance.

- In a 2\text{–}port network where the right port is open, that is, when \( i_2 = 0 \), the \( g_{11} \) parameter represents the input admittance of the two-port network. In other words,

\[ g_{11} = \frac{i_1}{v_1} \bigg|_{i_2 = 0} \]

- In a 2\text{–}port network where the left port is shorted, that is, when \( v_1 = 0 \), the \( g_{12} \) parameter represents a current gain (or loss) in the two-port network. In other words,

\[ g_{12} = \frac{i_1}{i_2} \bigg|_{v_1 = 0} \]
• In a 2-port network where the right port is open, that is, when \( i_2 = 0 \), the \( g_{21} \) parameter represents a voltage gain (or loss). In other words,

\[
g_{21} = \frac{v_2}{v_1} \bigg|_{i_2 = 0}
\]

• In a 2-port network where the left port is shorted, that is, when \( v_1 = 0 \), the \( g_{22} \) parameter represents an output impedance. In other words,

\[
g_{22} = \frac{v_2}{i_2} \bigg|_{v_1 = 0}
\]

• The reciprocity theorem states that if a voltage applied in one branch of a linear, two-port passive network produces a certain current in any other branch of this network, the same voltage applied in the second branch will produce the same current in the first branch. The reverse is also true, that is, if current applied at one node produces a certain voltage at another, the same current at the second node will produce the same voltage at the first.

• A two-port network is said to be reciprocal if any of the following relationships exists.

\[
\begin{align*}
z_{21} &= z_{12} \\
y_{21} &= y_{12} \\
h_{21} &= -h_{12} \\
g_{21} &= -g_{12}
\end{align*}
\]

• A two-port network is said to be symmetrical if any of the following relationships exist.

\[
\begin{align*}
z_{21} &= z_{12} & \text{and} & & z_{22} &= z_{11} \\
y_{21} &= y_{12} & \text{and} & & y_{22} &= y_{11} \\
h_{21} &= -h_{12} & \text{and} & & h_{11} h_{22} - h_{12} h_{21} &= 1 \\
g_{21} &= -g_{12} & \text{and} & & g_{11} g_{22} - g_{12} g_{21} &= 1
\end{align*}
\]
9.7 Exercises

1. For the network of Figure 9.59, find the $z$ parameters.

2. For the network of Figure 9.60, find the $y$ parameters.

3. For the network of Figure 9.61, find the $h$ parameters.

4. For the network of Figure 9.62, find the $g$ parameters.
5. The equations describing the $h$ parameters can be used to represent the network of Figure 9.63. This network is a transistor equivalent circuit for the common-emitter configuration and the $h$ parameters given are typical values for such a circuit. Compute the voltage gain and current gain for this network if a voltage source of $v_1 = cos\omega t \, mV$ in series with $800 \, \Omega$ is connected at the input (left side), and a $5 \, K\Omega$ load is connected at the output (right side).

$$h_{11} = 1.2 \, K\Omega$$

$$h_{12} = 2 \times 10^{-4}$$

$$h_{21} = 50$$

$$h_{22} = 50 \times 10^{-6} \, \Omega^{-1}$$

*Figure 9.63. Network for Exercise 5.*
9.8 Solutions to Exercises

1. \[
\begin{align*}
  z_{11} &= \frac{v_1}{i_1} |_{i_2 = 0} \\
  i_1 &= 1 \text{ A} \\
  i_{5\Omega} &= \frac{(10 + 20)}{(5 + 10 + 20)} i_1 = \frac{30}{35} \times 1 = \frac{6}{7} \text{ A} \\
  v_1 &= 5 i_{5\Omega} = 5 \times \frac{6}{7} = \frac{30}{7} \text{ V} \\
  z_{11} &= \frac{v_1}{i_1} = \frac{\frac{30}{7}}{1} = \frac{30}{7} \Omega \\

  z_{12} &= \frac{v_1}{i_2} |_{i_1 = 0} \\
  i_2 &= 1 \text{ A} \\
  i_{5\Omega} &= \frac{20}{(20 + 5 + 10)} i_2 = \frac{20}{35} \times 1 = \frac{4}{7} \text{ A} \\
  v_1 &= 5 \times \frac{4}{7} = \frac{20}{7} \text{ V} \\
  z_{12} &= \frac{v_1}{i_2} = \frac{\frac{20}{7}}{1} = \frac{20}{7} \Omega \\

  z_{21} &= \frac{v_2}{i_1} |_{i_2 = 0} \\
  i_1 &= 1 \text{ A} \\
  v_1 &= 5 \text{ V} \\
  v_2 &= 20 \text{ V} \\
  z_{21} &= \frac{v_2}{i_1} = \frac{20 \text{ V}}{1 \text{ A}} = 20 \Omega
\end{align*}
\]
Chapter 9 One- and Two-port Networks

\[ i_{20\Omega} = \frac{5}{(5 + 10 + 20)} i_1 = \frac{5}{35} \times 1 = 1/7 \text{ A} \]

\[ v_2 = 20 \times \frac{1}{7} = 20/7 \text{ V} \]

\[ z_{21} = \frac{v_2}{i_1} = \frac{20/7}{1} = 20/7 \text{ \Omega} \]

We observe that

\[ z_{21} = z_{12} \]

\[ z_{22} = \frac{v_2}{i_2} \bigg|_{i_1 = 0} \]

\[ i_{20\Omega} = \frac{(10 + 5)}{(20 + 10 + 5)} i_2 = \frac{15}{35} \times 1 = 3/7 \text{ A} \]

\[ v_2 = 20 \times \frac{3}{7} = 60/7 \text{ V} \]

\[ z_{22} = \frac{v_1}{i_2} = \frac{60/7}{1} = 60/7 \text{ \Omega} \]

2.

\[ y_{11} = \frac{i_1}{v_1} \bigg|_{v_2 = 0} \]

\[ R_{eq} = 5 \parallel 20 = 4 \text{ \Omega} \]

\[ i_1 = \frac{v_1}{R_{eq}} = 1/4 \text{ A} \]

\[ y_{11} = \frac{i_1}{v_1} = \frac{1/4}{1} = 1/4 \text{ \Omega}^{-1} \]
We observe that

\[ v_1 = 0 \]

\[ v_2 = 1 \text{ V} \]

\[ v_{5\Omega} = v_2 = 1 \text{ V} \]

\[ i_1 = -v_{5\Omega}/5 = -1/5 \text{ A} \]

\[ y_{12} = i_1/v_2 = -1/5/1 = -1/5 \Omega^{-1} \]

\[ v_{5\Omega} = v_1 = 1 \text{ V} \]

\[ i_2 = -v_{5\Omega}/5 = -1/5 \text{ A} \]

\[ y_{21} = i_2/v_1 = -1/5/1 = -1/5 \Omega^{-1} \]

\[ y_{21} = y_{12} \]

\[ i_2 = v_2/R_{eq} = 1/(5 \parallel 15) = 1/(75/20) = 4/15 \text{ A} \]

\[ y_{22} = i_2/v_2 = 4/15/1 = 4/15 \Omega^{-1} \]
3. \[ h_{11} = \frac{v_1}{i_1} \bigg|_{v_2=0} \]

\[ i_1 = 1 \, \text{A} \]

\[ v_2 = 0 \]

\[ i_1 \Omega = \frac{4}{(1+4)} i_1 = \frac{4}{5} \times 1 = 4/5 \, \text{A} \]

\[ v_1 = 1 \times i_1 \Omega = 4/5 \, \text{V} \]

\[ h_{11} = \frac{v_1}{i_1} = \frac{4/5}{1} = 4/5 \, \Omega \]

\[ h_{12} = \frac{v_1}{v_2} \bigg|_{v_2=0} \]

\[ i_2 = \frac{v_2}{R_{eq}} = \frac{1}{6 \parallel (4+1)} = \frac{1}{30/11} = 11/30 \, \text{A} \]

\[ v_1 = 1 \times i_1 \Omega = 1 \times \frac{6}{(6+4+1)} \times i_2 = 1 \times \frac{6}{11} \times \frac{11}{30} = 1/5 \, \text{V} \]

\[ h_{12} = \frac{v_1}{v_2} = \frac{1/5}{1} = 1/5 \, (\text{dimensionless}) \]

\[ h_{21} = \frac{i_2}{i_1} \bigg|_{v_2=0} \]

\[ i_2 = \frac{1}{(1+4)} \times (-i_1) = \frac{1}{5} \times (-1) = -1/5 \, \text{A} \]
We observe that

$$h_{21} = \frac{i_2}{i_1} = \frac{-1/5}{1} = -1/5$$

$$h_{21} = -h_{12}$$

4.

$$h_{22} = \frac{i_2}{v_2}$$

$$i_2 = \frac{v_2}{R_{eq}} = \frac{1}{6 \parallel (4 + 1)} = \frac{1}{30/11} = 11/30 \text{ A}$$

$$h_{22} = \frac{i_2}{v_2} = \frac{11/30}{1} = 11/30 \Omega^{-1}$$

$$i_1 = \frac{v_1}{R_{eq}} = \frac{1}{1 \parallel (4 + 6))} = \frac{1}{10/11} = 11/10 \text{ A}$$

$$g_{11} = \frac{i_1}{v_1} = \frac{11/10}{1} = 11/10 \Omega^{-1}$$

$$g_{12} = \frac{i_1}{i_2}$$

$$i_2 = 1 \text{ A}$$
We observe that

\[ i_1 = \left( \frac{6}{6+4} \right) (-i_2) = \frac{6}{10} = -\frac{3}{5} \text{ A} \]

\[ g_{12} = \frac{i_1}{i_2} = \frac{-3/5}{1} = -\frac{3}{5} \text{ (dimensionless)} \]

\[ g_{21} = \frac{v_2}{v_1} \bigg|_{i_2 = 0} \]

\[ i_1 = \frac{v_1}{R_{eq}} = \frac{1}{\frac{1}{1} \parallel (4+6)} = \frac{1}{10/11} = \frac{11}{10} \text{ A} \]

\[ v_2 = 6 \times i_6 \Omega = 6 \times \left( \frac{1}{1+4+6} \cdot \frac{11}{10} \right) = \frac{3}{5} \text{ V} \]

\[ g_{21} = \frac{v_2}{v_1} = \frac{3/5}{1} = \frac{3}{5} \]

We observe that

\[ g_{21} = -g_{12} \]

\[ g_{22} = \frac{v_2}{i_2} \bigg|_{v_1 = 0} \]

\[ v_2 = 6 \times i_6 \Omega = 6 \times \left( \frac{4}{6+4} \times i_2 \right) = \frac{24}{10} \times 1 = \frac{12}{5} \text{ V} \]

\[ g_{22} = \frac{v_2}{i_2} = \frac{\frac{12}{5}}{1} = \frac{12}{5} \Omega \]
5. We recall that
\[ v_1 = h_{11}i_1 + h_{12}v_2 \quad (1) \]
\[ i_2 = h_{21}i_1 + h_{22}v_2 \quad (2) \]

With the voltage source \( v_i = \cos \omega t \text{ mV} \) in series with 800 Ω connected at the input and a 5 KΩ load connected at the output the network is as shown below.

The network above is described by the equations
\[
(800 + 1200)i_1 + 2 \times 10^{-4}v_2 = 10^{-3} \\
50i_1 + 50 \times 10^{-6}v_2 = \frac{-v_2}{5000}
\]
or
\[
2 \times 10^{3}i_1 + 2 \times 10^{-4}v_2 = 10^{-3} \\
50i_1 + 250 \times 10^{-6}v_2 = 0
\]

We write the two equations above in matrix form and use MATLAB for the solution.

\[ A = \begin{bmatrix} 2 \times 10^3 & 2 \times 10^{-4} \\ 50 & 50 \times 10^{-6} \end{bmatrix}; \ B = \begin{bmatrix} 10^3 \ (-3) \ 0 \end{bmatrix}; \ X = A \backslash B; \]
\[ \text{fprintf('i1 = %5.2e A \n',X(1)); fprintf('v2 = %5.2e V',X(2))} \]
\[ i1 = 5.10e-007 \quad V2 = -1.02e-001 \]

Therefore,
\[ i_1 = 0.51 \text{ μA} \quad (3) \]
\[ v_2 = -102 \text{ mV} \quad (4) \]

Next, we use (1) and (2) to find the new values of \( v_1 \) and \( i_2 \)
\[ v_1 = 1.2 \times 10^3 \times 0.51 \times 10^{-6} + 2 \times 10^{-4} \times (-102 \times 10^{-3}) = 0.592 \text{ mV} \]
The voltage gain is

\[ G_V = \frac{v_2}{v_1} = \frac{-102 \text{ mV}}{0.592 \text{ mV}} = -172.3 \]

and the minus (−) sign indicates that the output voltage is 180° out-of-phase with the input.

The current gain is

\[ G_I = \frac{i_2}{i_1} = \frac{20.4 \text{ µA}}{0.51 \text{ µA}} = 40 \]

and the output current is in phase with the input.
Chapter 10

Three-Phase Systems

This chapter is an introduction to three-phase power systems. The advantages of three-phase system operation are listed and computations of three phase systems are illustrated by several examples.

10.1 Advantages of Three-Phase Systems

The circuits and networks we have discussed thus far are known as single-phase systems and can be either DC or AC. We recall that AC is preferable to DC because voltage levels can be changed by transformers. This allows more economical transmission and distribution. The flow of power in a three-phase system is constant rather than pulsating. Three-phase motors and generators start and run more smoothly since they have constant torque. They are also more economical.

10.2 Three-Phase Connections

Figure 10.1 shows three single AC series circuits where, for simplicity, we have assumed that the internal impedance of the voltage sources have been combined with the load impedance. We also have assumed that the voltage sources are $120^\circ$ out-of-phase, the load impedances are the same, and thus the currents $I_a$, $I_b$, and $I_c$ have the same magnitude but are $120^\circ$ out-of-phase with each other as shown in Figure 10.2.

![Figure 10.1. Three circuits with $120^\circ$ out-of-phase voltage sources](image)

![Figure 10.2. Waveforms for three $120^\circ$ out-phase currents](image)
Let us use a single wire for the return current of all three circuits as shown below. This arrangement is known as *four-wire, three-phase system*.

![Four-wire, three-phase system diagram](image)

*Figure 10.3. Four-wire, three-phase system*

This arrangement shown in Figure 10.3 uses only 4 wires instead of the 6 wires shown in Figure 10.1. But now we must find the relative size of the common return wire that it would be sufficient to carry all three currents \( I_a + I_b + I_c \).

We have assumed that the voltage sources are equal in magnitude and \( 120^\circ \) apart, and the loads are equal. Therefore, the currents will be balanced (equal in magnitude and \( 120^\circ \) out-of-phase). These currents are shown in the phasor diagram of Figure 10.4.

![Phasor diagram for three-phase balanced system](image)

*Figure 10.4. Phasor diagram for three-phase balanced system*
From figure 10.4 we observe that the sum of these currents, added vectorially, is zero. Therefore, under ideal (perfect balance) conditions, the common return wire carries no current at all. In a practical situation, however, is not balanced exactly and the sum is not zero. But still it is quite small and in a four-wire three-phase system the return wire is much smaller than the other three. Figure 10.5 shows a four-wire, three-phase $Y$–system where $|V_a| = |V_b| = |V_c|$, the three loads are identical, and $I_n$ is the current in the neutral (fourth) wire.

\[ Y_{\text{system}} - \frac{V_a \cos \omega t}{Z_{\text{LOAD}}} \quad I_a \]

\[ = \frac{V_b \cos(\omega t - 120^\circ)}{Z_{\text{LOAD}}} \quad I_b \]

\[ = \frac{V_c \cos(\omega t - 240^\circ)}{Z_{\text{LOAD}}} \quad I_c \]

\[ I_n \]

Figure 10.5. Four-wire, three-phase $Y$–system

A three-wire three-phase $Y$–system is shown in Figure 10.6 where $|V_a| = |V_b| = |V_c|$, and the three loads are identical.

\[ Y_{\text{system}} - \frac{V_a \cos \omega t}{Z_{\text{LOAD}}} \quad I_a \]

\[ = \frac{V_b \cos(\omega t - 120^\circ)}{Z_{\text{LOAD}}} \quad I_b \]

\[ = \frac{V_c \cos(\omega t - 240^\circ)}{Z_{\text{LOAD}}} \quad I_c \]

This arrangement shown in Figure 10.6 could be used only if all the three voltage sources are perfectly balanced, and if the three loads are perfectly balanced also. This, of course, is a physical impossibility and therefore it is not used.
Chapter 10  Three-Phase Systems

A three-wire three-phase $\Delta - load$ system is shown in Figure 10.7 where $|V_a| = |V_b| = |V_c|$, and the three loads are identical. We observe that while the voltage sources are connected as a $Y$-system, the loads are connected as a $\Delta$-system and hence the name $\Delta - load$.

![Figure 10.7. Three-wire, three-phase $\Delta - load$ system](image)

This arrangement offers the advantage that the $\Delta$-connected loads need not be accurately balanced. However, a $\Delta$-connection with only three voltages is not used for safety reasons, that is, it is a safety requirement to have a connection from the common point to the ground as shown in Figure 10.5.

10.3 Transformer Connections in Three-Phase Systems

Three-phase power systems use transformers to raise or to lower voltage levels. A typical generator voltage, typically 13.2 $KV$, is stepped up to hundreds of kilovolts for transmission over long distances. This voltage is then stepped down; for major distribution may be stepped down at a voltage level anywhere between 15 $KV$ to 50 $KV$, and for local distribution anywhere between 2.4 $KV$ to 12 $KV$. Finally, the electric utility companies furnish power to industrial and commercial facilities at 480 $V$, 120 $V$, and 240 $V$ at residential areas. All voltage levels are in $RMS$ values.

Figure 10.8 shows a bank of three single phase transformers where the primary is $\Delta$-connected, while the secondary is $Y$-connected. This $\Delta - Y$ connection is typical of transformer installations at generating stations.

Figure 10.9 shows a single-phase three-wire system where the middle of the three wires is center-tapped at the transformer secondary winding. As indicated, voltage between the outer wires is 240 $V$ while voltage from either of the two wires to the centered (neutral) wire is 120 $V$. This arrangement is used in residential areas.
Industrial facilities need three-phase power for three-phase motors. Three-phase motors run smoother and have higher efficiency than single-phase motors. A $Y - \Delta$ connection is shown in Figure 10.10 where the secondary provides 240 V three-phase power to the motor, and one of the transformers of the secondary is center-tapped to provide 120 V to the lighting load.

### 10.4 Line-to-Line and Line-to-Neutral Voltages and Currents

We assume that the perfectly balanced $Y$-connected load of Figure 10.11 is perfectly balanced, that is, the three loads are identical. We also assume that the applied voltages are 120° out-of-phase but they have the same magnitude; therefore there is no current flowing from point $n$ to the ground. The currents $I_a$, $I_b$ and $I_c$ are referred to as the line currents and the currents $I_{an}$, $I_{bn}$, and $I_{cn}$ as the phase currents. Obviously, in a $Y$-connected load, the line and phase currents are the same.
Now, we consider the phasor diagram of Figure 10.12.

Figure 10.10. Typical 3-phase distribution system

Figure 10.11. Perfectly balanced Y-connected load

Figure 10.12. Phasor diagram for Y-connected perfectly balanced load
If we choose \( I_a \) as our reference, we have

\[
I_a = I_a \angle 0^\circ \\
I_b = I_a \angle -120^\circ \\
I_c = I_a \angle +120^\circ
\]  

(10.1)  

(10.2)  

(10.3)

These equations define the balance set of currents of positive phase sequence \( a - b - c \).

Next, we consider the voltages. Voltages \( V_{ab} \), \( V_{ac} \), and \( V_{bc} \) are referred to as line-to-line voltages and voltages \( V_{an} \), \( V_{bn} \), and \( V_{cn} \) as phase voltages. We observe that in a \( Y \)-connected load, the line and phase voltages are not the same.

We will now derive the relationships between line and phase voltages in a \( Y \)-connected load.

Arbitrarily, we choose \( V_{an} \) as our reference phase voltage. Then,

\[
V_{an} = V_{an} \angle 0^\circ \\
V_{bn} = V_{an} \angle -120^\circ \\
V_{cn} = V_{an} \angle +120^\circ
\]  

(10.4)  

(10.5)  

(10.6)

These equations define a positive phase sequence \( a - b - c \). These relationships are shown in Figure 10.13.

![Figure 10.13. Phase voltages in a Y-connected perfectly balanced load](image)

The \( Y \)-connected load is repeated in Figure 10.14 for convenience.

From Figure 10.14
These can also be derived from the phasor diagram of Figure 10.15.

From geometry and the law of sines we find that in a balanced three-phase, positive phase sequence $Y$-connected load, the line and phase voltages are related as
The other two line-to-line voltages can be easily obtained from the phasor diagram of the previous page.

Now, let us consider a $\Delta$-connected load shown in Figure 10.16.

![Figure 10.16. Line and phase currents in $\Delta$-connected load](image)

We observe that the line and phase voltages are the same, but the line and phase currents are not the same. To find the relationship between the line and phase currents, we apply KCL at point $a$ and we get:

$$I_{ab} = I_a + I_{ca}$$

or

$$I_a = I_{ab} - I_{ca} \quad (10.11)$$

The line currents $I_b$ and $I_c$ are derived similarly, and the phase-to-line current relationship in a $\Delta$-connected load is shown in the phasor diagram of Figure 10.17.

From geometry and the law of sines we find that a balanced three-phase, positive phase sequence $\Delta$-connected load, the line and phase currents are related as

$$I_a = \sqrt{3} I_{ab} \angle -30^\circ \quad \Delta \text{-connected load} \quad (10.12)$$

The other two line currents can be easily obtained from the phasor diagram of Figure 10.17.
10.5 Equivalent Y and Δ Loads

In this section, we will establish the equivalence between the $Y$ and $Δ$ combinations shown in Figure 10.18.

In the $Y$-connection, the impedance between terminals $B$ and $C$ is

$$Z_{BC \ Y} = Z_b + Z_c \quad (10.13)$$

and in the $Δ$-connection, the impedance between terminals $B$ and $C$ is $Z_2$ in parallel with the sum $Z_1 + Z_3$, that is,

$$Z_{BC \ Δ} = \frac{Z_2(Z_1 + Z_3)}{Z_1 + Z_2 + Z_3} \quad (10.14)$$
Equating (10.13) and (10.14) we get

\[ Z_b + Z_c = \frac{Z_2(Z_1 + Z_3)}{Z_1 + Z_2 + Z_3} \]  

(10.15)

Similar equations for terminals \( AB \) and \( CA \) are derived by rotating the subscripts of (10.15) in a cyclical manner. Then,

\[ Z_a + Z_b = \frac{Z_3(Z_1 + Z_2)}{Z_1 + Z_2 + Z_3} \]  

(10.16)

and

\[ Z_c + Z_a = \frac{Z_1(Z_2 + Z_3)}{Z_1 + Z_2 + Z_3} \]  

(10.17)

Equations (10.15) and (10.17) can be solved for \( Z_a \) by adding (10.16) with (10.17), subtracting (10.15) from this sum, and dividing by two. That is,

\[ 2Z_a + Z_b + Z_c = \frac{2Z_1Z_3 + Z_2Z_3 + Z_1Z_2 + Z_1Z_3}{Z_1 + Z_2 + Z_3} = \frac{2Z_1Z_3 + Z_2Z_3 + Z_1Z_2}{Z_1 + Z_2 + Z_3} \]  

(10.18)

\[ 2Z_a + Z_b + Z_c - Z_b - Z_c = \frac{2Z_1Z_3 + Z_2Z_3 + Z_1Z_2 - Z_1Z_2 - Z_2Z_3}{Z_1 + Z_2 + Z_3} \]  

(10.19)

\[ 2Z_a = \frac{2Z_1Z_3}{Z_1 + Z_2 + Z_3} \]  

(10.20)

\[ Z_a = \frac{Z_1Z_3}{Z_1 + Z_2 + Z_3} \]  

(10.21)

Similar equations for \( Z_b \) and \( Z_c \) are derived by rotating the subscripts of (10.21) in a cyclical manner.

Thus, the three equations that allow us to change any \( \Delta \)-connection of impedances into a \( Y \)-connection are given by (10.22).

\[
\begin{align*}
Z_a &= \frac{Z_1Z_3}{Z_1 + Z_2 + Z_3} \\
Z_b &= \frac{Z_2Z_3}{Z_1 + Z_2 + Z_3} \\
Z_c &= \frac{Z_1Z_2}{Z_1 + Z_2 + Z_3}
\end{align*}
\]

(10.22)

\[ \Delta \rightarrow Y \text{ Conversion} \]
Often, we wish to make the conversion in the opposite direction, that is, from $Y$ to $\Delta$. This conversion is performed as follows:

Consider the $Y$ and $\Delta$ combinations of Figure 10.8 repeated for convenience.

From Figure (a),

\[ V_{AB} = Z_a I_A - Z_b I_B \]  \hspace{1cm} (10.23)
\[ V_{BC} = Z_b I_B - Z_c I_C \]  \hspace{1cm} (10.24)
\[ V_{CA} = Z_c I_C - Z_a I_A \]  \hspace{1cm} (10.25)

If we attempt to solve equations (10.23), (10.24) and (10.25) simultaneously, we will find that the determinant $\Delta$ of these sets of equations is singular, that is, $\Delta = 0$. This can be verified with Cramer’s rule as follows:

\[
\begin{align*}
Z_a I_A - Z_b I_B + 0 &= V_{AB} \\
0 + Z_b I_B - Z_c I_C &= V_{BC} \\
-Z_a I_A + 0 + Z_c I_C &= V_{CA}
\end{align*}
\]  \hspace{1cm} (10.26)

\[
\Delta = \begin{bmatrix}
Z_a & -Z_b & 0 \\
0 & Z_b & -Z_c \\
-Z_a & 0 & Z_c
\end{bmatrix} = Z_a Z_b Z_c - Z_a Z_b Z_c + 0 + 0 + 0 = 0 \]  \hspace{1cm} (10.27)

This result suggests that the equations of (10.26) are not independent and therefore, no solution exists. However, a solution can be found if, in addition to (10.23) through (10.25), we use the equation

\[ I_A + I_B + I_C = 0 \]  \hspace{1cm} (10.28)
Equivalent Y and Δ Loads

Solving (10.28) for $I_C$ we get:

$$I_C = -I_A - I_B$$  \hspace{1cm} (10.29)

and by substitution into (10.25),

$$V_{CA} = -Z_c I_A - Z_c I_B - Z_a I_A = -(Z_a + Z_c) I_A - Z_c I_B$$  \hspace{1cm} (10.30)

From (10.23) and (10.30),

$$Z_a I_A - Z_b I_B = V_{AB}$$

$$-(Z_a + Z_c) I_A - Z_c I_B = V_{CA}$$  \hspace{1cm} (10.31)

and by Cramer’s rule,

$$I_A = \frac{D_1}{\Delta} \quad I_B = \frac{D_2}{\Delta}$$  \hspace{1cm} (10.32)

where

$$\Delta = \begin{bmatrix} Z_a & -Z_b \\ -(Z_a + Z_c) & -Z_c \end{bmatrix} = -Z_c Z_a - Z_a Z_b - Z_b Z_c$$  \hspace{1cm} (10.33)

and

$$D_1 = \begin{bmatrix} V_{AB} & -Z_b \\ V_{CA} & -Z_c \end{bmatrix} = -Z_c V_{AB} + Z_b V_{CA}$$  \hspace{1cm} (10.34)

Then,

$$I_A = \frac{D_1}{\Delta} = \frac{-Z_c V_{AB} + Z_b V_{CA}}{-Z_a Z_b - Z_b Z_c - Z_c Z_a} = \frac{Z_c V_{AB} - Z_b V_{CA}}{Z_a Z_b + Z_b Z_c + Z_c Z_a}$$  \hspace{1cm} (10.35)

Similarly,

$$I_B = \frac{D_2}{\Delta} = \frac{Z_a V_{BC} - Z_c V_{AB}}{Z_a Z_b + Z_b Z_c + Z_c Z_a}$$  \hspace{1cm} (10.36)

and by substitution of $I_A$ and $I_B$ into (10.28),

$$I_C = \frac{Z_b V_{CA} - Z_d V_{BC}}{Z_a Z_b + Z_b Z_c + Z_c Z_a}$$  \hspace{1cm} (10.37)

Therefore, for the Y-connection which is repeated in Figure 10.20 for convenience, we have:
For the Δ-connection, which is also repeated in Figure 10.21 for convenience, the line currents are:

\[
I_A = \frac{V_{AB} - V_{CA}}{Z_3 Z_1} \\
I_B = \frac{V_{BC} - V_{AB}}{Z_2 Z_3} \\
I_C = \frac{V_{CA} - V_{BC}}{Z_1 Z_2}
\]  

(10.39)

Now, the sets of equations of (10.38) and (10.39) are equal if
Equivalent Y and Δ Loads

\[
\frac{Z_c V_{AB} - Z_b V_{CA}}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{V_{AB}}{Z_3} - \frac{V_{CA}}{Z_1} \tag{10.40}
\]

\[
\frac{Z_a V_{BC} - Z_c V_{AB}}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{V_{BC}}{Z_2} - \frac{V_{AB}}{Z_3} \tag{10.41}
\]

\[
\frac{Z_b V_{CA} - Z_a V_{BC}}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{V_{CA}}{Z_1} - \frac{V_{BC}}{Z_2} \tag{10.42}
\]

From (10.40)

\[
\frac{Z_c}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{l}{Z_3} \quad \text{and} \quad \frac{Z_b}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{l}{Z_1} \tag{10.43}
\]

and from (10.41)

\[
\frac{Z_a}{Z_a Z_b + Z_b Z_c + Z_c Z_a} = \frac{l}{Z_2} \tag{10.44}
\]

Rearranging, we get:

\[
Z_1 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_b} \quad Z_2 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_a} \quad Z_3 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_c} \tag{10.45}
\]

Y → Δ Conversion

Example 10.1

For the circuit of Figure 10.22, use the Y → Δ conversion to find the currents in the various branches as indicated.

Solution:

Let us indicate the nodes as \(a\), \(b\), \(c\), and \(d\), and denote the 90 Ω, 90 Ω, and 90 Ω resistances as \(R_a\), \(R_b\), and \(R_c\) respectively as shown in Figure 10.23.

Next, we replace the Y connection formed by \(a\), \(b\), \(c\), and \(d\) with the equivalent Δ connection shown in Figure 10.24.
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Figure 10.22. Circuit (a) for Example 10.1

Figure 10.23. Circuit (b) for Example 10.1

Figure 10.24. Circuit (c) for Example 10.1
Now, with reference to the circuits of Figures 10.23 and 10.24, and the relations of (10.45), we get:

\[ R_1 = \frac{R_a R_b + R_b R_c + R_c R_a}{R_b} = \frac{90 \times 80 + 80 \times 50 + 50 \times 90}{80} = \frac{15700}{80} \approx 196 \ \Omega \]

\[ R_2 = \frac{R_a R_b + R_b R_c + R_c R_a}{R_a} = \frac{15700}{90} \approx 174 \ \Omega \]

\[ R_3 = \frac{R_a R_b + R_b R_c + R_c R_a}{R_c} = \frac{15700}{50} = 314 \ \Omega \]

Combination of parallel resistances in the circuit of Figure 10.24 yields

\[ R_{bd} = \frac{196 \times 60}{196 + 60} \approx 46 \ \Omega \]

and

\[ R_{ad} = \frac{314 \times 70}{314 + 70} \approx 57 \ \Omega \]

The circuit of Figure 10.24 reduces to the circuit of Figure 10.25. The circuit of Figure 10.25 can be further simplified as shown in Figure 10.26.

From the circuit of Figure 10.26,

\[ I_2 = \frac{120}{174} = 0.69 \ A \quad (10.46) \]

\[ I_3 = \frac{120}{103} = 1.17 \ A \quad (10.47) \]

---

![Image](image-url)
By addition of (10.46) and (10.47)

\[ I_1 = I_2 + I_3 = 0.69 + 1.17 = 1.86 \quad (10.48) \]

To compute the other currents, we return to the circuit of Figure 10.25 which, for convenience, is repeated as Figure 10.27 and it is denoted as Circuit (f).

For the circuit of Figure 10.27, by the voltage division expression

\[ V_{ad} = \frac{46}{46 + 57} \times 120 = 53.6 \text{ V} \quad (10.49) \]

\[ V_{db} = \frac{57}{46 + 57} \times 120 = 66.4 \text{ V} \quad (10.50) \]
Next, we return to the circuit of Figure 10.24 which, for convenience, is repeated as Figure 10.28 and denoted as Circuit (g).

![Circuit diagram](image)

**Figure 10.28. Circuit (g) for Example 10.1**

From the circuit of figure 10.28,

\[
I_4 = \frac{V_{ad}}{60} = \frac{66.4}{70} = 0.89 \text{ A} \tag{10.51}
\]

and

\[
I_5 = \frac{V_{bd}}{70} = \frac{53.6}{60} = 0.95 \text{ A} \tag{10.52}
\]

Finally, we return to the circuit of Figure 10.23 which, for convenience, is repeated as Figure 10.29 and denoted as Circuit (h).

![Circuit diagram](image)

**Figure 10.29. Circuit (h) for Example 10.1**
For the circuit of Figure 10.29, by KCL,

\[ I_7 = I_1 - I_4 = 1.86 - 0.89 = 0.97 \, A \]  
(10.53)

\[ I_8 = I_1 - I_5 = 1.86 - 0.95 = 0.91 \, A \]  
(10.54)

and

\[ I_6 = I_5 - I_4 = 0.95 - 0.89 = 0.06 \, A \]  
(10.55)

Of course, we could have found the branch currents with nodal or mesh analysis.

Quite often, the \( Y \) and \( \Delta \) arrangements appear as shown in Figure 10.30 and they are referred to as the tee (T) and pi (\( \pi \)) circuits. Consequently, the formulas we developed for the \( Y \) and \( \Delta \) arrangements can be used with the tee and \( \pi \) arrangements.

![Figure 10.30. T and \( \pi \) circuits](image)

In communications theory, the T and \( \pi \) circuits are symmetrical, i.e., \( Z_a = Z_b \) and \( Z_1 = Z_2 \).

### 10.6 Computation by Reduction to Single Phase

When we want to compute the voltages, currents, and power in a balanced three-phase system, it is very convenient to use the \( Y \)-connection and work with one phase only. The other phases will have corresponding quantities (voltage, current, and power) exactly the same except for a time difference of \( 1/3 \) cycle. Thus, if current is found for phase \( a \), the current in phase \( b \) will be \( 120^\circ \) out-of-phase but it will have the same magnitude as phase \( a \). Likewise, phase \( c \) will be \( 240^\circ \) out-of-phase with phase \( a \).

If the load happens to be \( \Delta \)-connected, we use the \( \Delta \rightarrow Y \) conversion shown in Figure 10.31 and the equations (10.57) on the next page.

Since the system is assumed to be balanced, the loads \( Z_1 = Z_2 = Z_3 \) and \( Z_a = Z_b = Z_c \). Therefore, the first equation in (10.57) reduces to:

\[
Z_a = \frac{Z_1Z_3}{Z_1 + Z_2 + Z_3} = \frac{Z_1^2}{3Z_1} = \frac{Z_1}{3}
\]  
(10.56)
Figure 10.31. $\Delta \rightarrow Y$ conversion

\[
\begin{align*}
Z_a &= \frac{Z_1Z_3}{Z_1 + Z_2 + Z_3} \\
Z_b &= \frac{Z_2Z_3}{Z_1 + Z_2 + Z_3} \\
Z_c &= \frac{Z_1Z_2}{Z_1 + Z_2 + Z_3}
\end{align*}
\]

and the same is true for the other phases.

### 10.7 Three-Phase Power

We can compute the power in a single phase and then multiply by three to find the total power in a three-phase system. Therefore, if a load is $Y$-connected, as in Figure 10.31 (b), the total three-phase power is given by

\[
P_{\text{TOTAL}} = 3|V_{AN}|I_A|\cos\theta
\]

$Y$-connected load

where $V_{AN}$ is the line-to-neutral voltage, $I_A$ is the line current, $\cos\theta$ is the power factor of the load, and $\theta$ is the angle between $V_{AN}$ and $I_A$.

If the load is $\Delta$-connected as in Figure 10.31 (a), the total three-phase power is given by

\[
P_{\text{TOTAL}} = 3|V_{AB}|I_{AB}|\cos\theta
\]

$\Delta$-connected load

We observe that relation (10.59) is given in terms of the line-to-neutral voltage and line current, and relation (10.58) in terms of the line-to-line voltage and phase current.
Chapter 10 Three-Phase Systems

Quite often, the line-to-line voltage and line current of a three-phase systems are given. In this case, we substitute (10.12), i.e., $|I_A| = \sqrt{3}|I_{AB}|$ into (10.59) and we get

$$P_{TOTAL} = \sqrt{3}|V_{AB}||I_A|\cos\theta_{LOAD}$$

(10.60)

$Y$ or $\Delta$ – connected load

It is important to remember that the power factor $\cos\theta_{LOAD}$ in (10.60) refers to the load, that is, the angle $\theta$ is not the angle between $V_{AB}$ and $I_A$.

Example 10.2

The three-phase generator of Figure 10.32 supplies 100 kW at 0.9 lagging power factor to the three-phase load. The line-to-line voltage at the load is 2400 V. The resistance of the line is 4 $\Omega$ per conductor and the inductance and capacitance are negligible. What line-to-line voltage must the generator supply to the line?

Solution:

The load per phase at 0.9 pf is

$$\frac{I}{3} \times 100 = 33.33 \text{ kW}$$

Figure 10.32. Circuit for Example 10.2

From (10.10),

$$V_{ab} = \sqrt{3}V_{an}\angle30^\circ$$

$Y$ – connected load

(10.61)

Then, the magnitude of the line-to-neutral at the load end is

$$|V_{an \text{ load}}| = \frac{|V_{ab \text{ load}}|}{\sqrt{3}} = \frac{2400}{\sqrt{3}} = 1386 \text{ V}$$

(10.62)

and the KVA per phase at the load is

$$\frac{\text{kW/phase}}{\text{pf}} = \frac{33.33}{0.9} = 37.0 \text{ KVA}$$

(10.63)
The line current in each of the three conductors is

\[ I_{\text{LINE}} = \frac{VA}{|V_{\text{an load}}|} = \frac{37000}{1386} = 26.7 \text{ A} \quad (10.64) \]

and the angle by which the line (or phase) current lags the phase voltage is

\[ \theta = \cos^{-1} 0.9 = 25.84^\circ \quad (10.65) \]

Next, let us assume that the line current in phase \( a \) lies on the real axis. Then, the phasor of the line-to-neutral voltage at the load end is

\[ V_{\text{an load}} = |V_{\text{an}}| \angle 25.84^\circ = 1386(\cos 25.84^\circ + j\sin 25.84^\circ) = 1247 + j604 \text{ V} \quad (10.66) \]

The voltage drop across a conductor is in phase with the line current since it resistive in nature. Therefore,

\[ V_{\text{COND}} = I_{\text{LINE}} \times R = 26.7 \times 4 = 106.8 \text{ V} \quad (10.67) \]

Now, the phasor line-to-neutral voltage at the generator end is

\[ V_{\text{an gen}} = V_{\text{an load}} + V_{\text{COND}} = 1247 + j604 + 106.8 = 1354 + j604 \quad (10.68) \]

and its magnitude is

\[ |V_{\text{an gen}}| = \sqrt{1354^2 + 604^2} = 1483 \text{ V} \quad (10.69) \]

Finally, the line-to-line voltage at the generator end is

\[ |V_{\text{line-line gen}}| = \sqrt{3} \times |V_{\text{an gen}}| = \sqrt{3} \times 1483 = 2569 \text{ V} \quad (10.70) \]

### 10.8 Instantaneous Power in Three-Phase Systems

A significant advantage of a three-power system is that the total power in a balanced three-phase system is constant. This is proved as follows:

We assume that the load is purely resistive. Therefore, the voltage and current are always in-phase with each other. Now, let \( V_p \) and \( I_p \) be the peak (maximum) voltage and current respectively, and \( |V| \) and \( |I| \) the magnitude of their RMS values. Then, the instantaneous voltage and current in phase \( a \) are given by

\[ v_a = V_p \cos \omega t = \sqrt{2}|V|\cos \omega t \quad (10.71) \]

\[ i_a = I_p \cos \omega t = \sqrt{2}|I|\cos \omega t \quad (10.72) \]
Multiplication of (10.71) and (10.72) yields the instantaneous power, and using the trigonometric identity

\[
\cos^2 \omega t = (\cos 2\omega t + 1)/2
\]

we get

\[
p_a = v_a \cdot i_a = 2|V||I| \cos^2 \omega t = |V||I|(\cos 2\omega t + 1)
\]

(10.74)

The voltage and current in phase \(b\) are equal in magnitude to those in phase \(a\) but they are \(120^\circ\) out-of-phase. Then,

\[
v_b = \sqrt{2}|V|\cos(\omega t - 120^\circ)
\]

(10.75)

\[
i_b = \sqrt{2}|V|\cos(\omega t - 120^\circ)
\]

(10.76)

\[
p_b = v_b \cdot i_b = 2|V||I| \cos^2(\omega t - 120^\circ) = |V||I|[(\cos(2\omega t - 240^\circ) + 1]
\]

(10.77)

Similarly, the power in phase \(c\) is

\[
p_c = v_c \cdot i_c = 2|V||I| \cos^2(\omega t - 240^\circ) = |V||I|[(\cos(2\omega t - 480^\circ) + 1]
\]

(10.78)

and the total instantaneous power is

\[
p_{total} = p_a + p_b + p_c
\]

\[
= |V||I|[(\cos 2\omega t + \cos(2\omega t - 240^\circ) + \cos(2\omega t - 480^\circ) + 3]
\]

(10.79)

Recalling that

\[
\cos(x - y) = \cos x \cos y + \sin x \sin y
\]

(10.80)

we find that the sum of the three cosine terms in (10.79) is zero. Then,

\[
\boxed{p_{total} = 3|V||I|}
\]

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(10.81)

Therefore, the instantaneous total power is constant and it is equal three times the average power.

The proof can be extended to include any power factor; thus, (10.81) can be also expressed as

\[
p_{total} = 3|V||I|\cos \theta
\]

(10.82)

Example 10.3

Figure 10.33 shows a three-phase feeder with two loads; one consists of a bank of lamps connected line-to neutral and the rating is given in the diagram; the other load is \(\Delta\)-connected and has the
impedance shown. Find the current in the feeder lines and the total power absorbed by the two loads.

\[ Z = 18 + j80 \]

**Figure 10.33. Diagram for Example 10.3**

**Solution:**

To facilitate the computations, we will reduce the given circuit to one phase (phase \(a\)) taken as reference, i.e., at zero degrees, as shown in Figure 10.34

\[ V_{L-L} = 220 \angle 0^\circ \text{ V} \]

**Figure 10.34. Single-phase representation of Figure 10.31**

We first compute the impedance \(Z_Y\). Using (10.56),

\[ Z_Y = \frac{Z_\Delta}{3} = \frac{18 + j80}{3} = \frac{82 \angle 77.32}{3} = 27.33 \angle 77.32 \text{ \Omega} \]

Next, we compute the lamp impedance \(Z_L\)

\[ Z_L = R_{lamp} = \frac{V_{rated}^2}{P_{rated}} = \frac{120^2}{500} = 28.8 \text{ \Omega} \]

The line-to-line voltage is given as \(V_{L-L} = 220 \text{ V}\); therefore, by (10.10), the line-to-neutral voltage \(V_{L-N}\) is
For convenience, we indicate these values in Figure 10.34 which now is as shown in Figure 10.35.

\[ V_{L-N} = \frac{V_{L-L}}{\sqrt{3}} = \frac{220 \angle 0^\circ}{\sqrt{3}} = 127 \angle 0^\circ \text{ V} \]

Figure 10.35. Diagram with computed values, Example 10.3

From Figure 10.35,

\[ I_Z = \frac{V_{L-N}}{Z_Y} = \frac{127 \angle 0^\circ}{27.33 \angle 77.32} = 4.65 \angle -77.32 = 1.02 - j4.54 \]

and

\[ I_L = \frac{V_{L-N}}{Z_L} = \frac{127 \angle 0^\circ}{28.8 \angle 0^\circ} = 4.41 \angle 0^\circ = 4.41 \]

Then,

\[ I_Z + I_L = 1.02 - j4.54 + 4.41 = 5.43 - j4.54 = 7.08 \angle -39.9^\circ \]

and the power delivered by phase \( a \) is

\[ P_A = V_{L-N} \cdot I_A = 127 \times 7.08 \times \cos(-39.9^\circ) = 690 \text{ watts} \]

Finally, the total power delivered to the entire load is three times of \( P_A \), that is,

\[ P_{\text{total}} = 3 \times 690 = 2070 \text{ watts} = 2.07 \text{ Kw} \]

Check:

Each lamp is rated 120 V and 500 w but operates at 127 V. Thus, each lamp absorbs

\[ \left( \frac{V_{\text{oper}}}{V_{\text{rated}}} \right)^2 = \frac{P_{\text{oper}}}{P_{\text{rated}}} \quad P_{\text{oper}} = \left( \frac{127}{120} \right)^2 \times 500 = 560 \text{ w} \]

and the power absorbed by the three lamps is

\[ P_{\text{lamps}} = 3 \times 560 = 1680 \text{ w} \]
The voltage across each impedance $Z$ in the $\Delta$-connected load is (see Figure 10.33) $220 \text{ V}$. Then, the current in each impedance $Z$ is

$$I_Z = \frac{V_{L-L}}{18 + j80} = \frac{220}{82 \angle 77.32} = 2.68 \angle -77.32 \text{ A}$$

and the power absorbed by each impedance $Z$ is

$$P = V_{L-L}I_Z\cos \theta = 220 \times 2.68 \times \cos(-77.32^\circ) = 129.4 \text{ watts}$$

The total power absorbed by the $\Delta$ load is

$$P_{\Delta} = 3 \times 129.4 = 388 \text{ watts}$$

and the total power delivered to the two loads is

$$P_{TOTAL} = P_{\text{lamps}} + P_{\Delta} = 2068 \text{ watts} = 2.068 \text{ kw}$$

This value is in close agreement with the value on the previous page.

### 10.9 Measuring Three-Phase Power

A *wattmeter* is an instrument which measures power in watts or kilowatts. It is constructed with two sets of coils, a current coil and a voltage coil where the interacting magnetic fields of these coils produce a torque which is proportional to the $V \times I$ product. It would appear then that one would need three wattmeters to measure the total power in a three-phase system. This is true in a four-wire system where the current in the neutral (fourth wire) is not zero. However, if the neutral carries no current, it can be eliminated thereby reducing the system to a three-wire three-phase system. In this section, we will show that the total power in a balanced three-wire, three-phase system can be measured with just two wattmeters.

Figure 10.36 shows three wattmeters connected to a $Y$ load* where each wattmeter has its current coil connected in one line, and its potential coil from that line to neutral. With this arrangement, Wattmeters 1, 2, and 3 measure power in phase $a$, $b$, and $c$ respectively.

Figure 10.37 shows a three-wire, three-phase system without a neutral. This arrangement occurs in systems where the load, such as an induction motor, has only three terminals. The lower end of the voltage coils can be connected to any reference point, say $p$. We will now show that with this arrangement, the sum of the three wattmeters gives the correct total power even though the reference point was chosen as any reference point.

* If the load were $\Delta$-connected, each wattmeter would have its current coil in one side of the $\Delta$ and its potential coil from line to line.
We recall that the average power $P_{ave}$ is found from
Measuring Three-Phase Power

\[ P_{ave} = \frac{1}{T} \int_{0}^{T} p \, dt = \frac{1}{T} \int_{0}^{T} v \, idt \]  
\hspace{1in} (10.83)

Then, the total power absorbed by the load of Figure 10.36 is

\[ P_{total} = \frac{1}{T} \int_{0}^{T} (v_{an}i_{a} + v_{bn}i_{b} + v_{cn}i_{c}) \, dt \]  
\hspace{1in} (10.84)

This is the true power absorbed by the load, not power indicated by the wattmeters.

Now, we will compute the total power indicated by the wattmeters. Each wattmeter measures the average of the line current times the voltage to point \( p \). Then,

\[ P_{wattmeters} = \frac{1}{T} \int_{0}^{T} (v_{ap}i_{a} + v_{bp}i_{b} + v_{cp}i_{c}) \, dt \]  
\hspace{1in} (10.85)

But

\[ v_{ap} = v_{an} + v_{np} \]
\[ v_{bp} = v_{bn} + v_{np} \]  
\hspace{1in} (10.86)
\[ v_{cp} = v_{cn} + v_{np} \]

and by substitution of these into (10.85), we get:

\[ P_{wattmeters} = \frac{1}{T} \int_{0}^{T} [(v_{an}i_{a} + v_{bn}i_{b} + v_{cn}i_{c}) + v_{np}(i_{a} + i_{b} + i_{c})] \, dt \]  
\hspace{1in} (10.87)

and since

\[ i_{a} + i_{b} + i_{c} = 0 \]  
\hspace{1in} (10.88)

then (10.87) reduces to

\[ P_{wattmeters} = \frac{1}{T} \int_{0}^{T} (v_{an}i_{a} + v_{bn}i_{b} + v_{cn}i_{c}) \, dt \]  
\hspace{1in} (10.89)

This relation is the same as (10.84); therefore, the power indicated by the wattmeters and the true power absorbed by the load are the same.

Some thought about the location of the arbitrarily selected point \( p \) would reveal a very interesting result. No matter where this point is located, the power relation (10.87) reduces to (10.89). Suppose that we locate point \( p \) on line \( c \). If we do this, the voltage coil of Wattmeter 3 is zero and thus the reading of this wattmeter is zero. Accordingly, we can remove this wattmeter and still obtain the true power with just Wattmeters 1 and 2 as shown in Figure 10.38.
10.10 Summary

- AC is preferable to DC because voltage levels can be changed by transformers. This allows more economical transmission and distribution.

- The flow of power in a three-phase system is constant rather than pulsating. Three-phase motors and generators start and run more smoothly since they have constant torque. They are also more economical.

- If the voltage sources are equal in magnitude and 120° apart, and the loads are also equal, the currents will be balanced (equal in magnitude and 120° out-of-phase).

- Industrial facilities need three-phase power for three-phase motors. Three-phase motors run smoother and have higher efficiency than single-phase motors.

- The equations \( I_a = I_a \angle 0° \), \( I_b = I_a \angle -120° \), \( I_c = I_a \angle +120° \) define a balanced set of currents of positive phase sequence \( a - b - c \).

- The equations \( V_{an} = V_{an} \angle 0° \), \( V_{bn} = V_{an} \angle -120° \), and \( V_{cn} = V_{an} \angle +120° \) also define a balanced set of voltages of positive phase sequence \( a - b - c \).

- In a Y-connected system
  \[
  V_{ab} = \sqrt{3} V_{an} \angle 30°
  \]

- In a Y-connected load, the line and phase currents are the same.

- In a \( \Delta \)-connected system
  \[
  I_a = \sqrt{3} I_{ab} \angle -30°
  \]
Summary

- In a Δ-connected load, the line and phase voltages are the same.

- For Δ → Y Conversion we use the relations

\[
Z_a = \frac{Z_1 Z_3}{Z_1 + Z_2 + Z_3} \\
Z_b = \frac{Z_2 Z_3}{Z_1 + Z_2 + Z_3} \\
Z_c = \frac{Z_1 Z_2}{Z_1 + Z_2 + Z_3}
\]

- For Y → Δ Conversion we use the relations

\[
Z_1 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_b} \\
Z_2 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_a} \\
Z_3 = \frac{Z_a Z_b + Z_b Z_c + Z_c Z_a}{Z_c}
\]

- When we want to compute the voltages, currents, and power in a balanced three-phase system, it is very convenient to use the Y-connection and work with one phase only.

- If a load is Y-connected, the total three-phase power is given by

\[
P_{TOTAL} = 3|V_{AN}||I_A|\cos\theta
\]

Y-connected load

- If the load is Δ-connected the total three-phase power is given by

\[
P_{TOTAL} = 3|V_{AB}||I_{AB}|\cos\theta
\]

Δ-connected load

\[(10.90)\]

- For any load (Y or Δ-connected) the total three-phase power can be computed from

\[
P_{TOTAL} = \sqrt{3}|V_{AB}||I_A|\cos\theta_{LOAD}
\]

Y or Δ-connected load

and it is important to remember that the power factor \(\cos\theta_{LOAD}\) refers to the load, that is, the angle \(\theta\) is not the angle between \(V_{AB}\) and \(I_A\).
10.11 Exercises

1. In the circuit of Figure 10.39, the line-to-line voltage is 100 V, the phase sequence is \(a - b - c\), and each \(Z = 10 \angle 30^\circ\). Compute:
   
a. the total power absorbed by the three-phase load.

   b. the wattmeter reading.

2. In the circuit of Figure 10.40 the lighting load is balanced. Each lamp is rated 500 w at 120 V. Assume constant resistance, that is, each lamp will draw rated current. The three-phase motor draws 5.0 Kw at a power factor of 0.8 lagging. The secondary of the transformer provides balanced 208 V line-to-line. The load is located 1500 feet from the three-phase transformer. The resistance and inductive reactance of the distribution line is 0.403 Ω and 0.143 Ω respectively per 1000 ft of the wire line. Compute line-to-line and line-to-neutral voltages at the load.
10.12 Solutions to Exercises

1. a

From the circuit above

\[ I_{ab} = \frac{V_{ab}}{Z} = \frac{100 \angle 0^\circ}{10 \angle 30^\circ} = 10 \angle -30^\circ = 10 \times \frac{\sqrt{3}}{2} - j10 \times \frac{1}{2} = 5 \sqrt{3} - j5 \]

\[ I_{ca} = \frac{V_{ca}}{Z} = \frac{100 \angle -240^\circ}{10 \angle 30^\circ} = 10 \angle -270^\circ = 10 \angle 90^\circ = j10 \]

\[ I_a = I_{ab} - I_{ca} = 5 \sqrt{3} - j5 - j10 = 5 \sqrt{3} - j15 \]

and with MATLAB

\[ x=5*\sqrt{3}-15j; \text{ fprintf(' |n'); }... \]
\[ \text{ fprintf('mag = %5.2f A | abs(x)); fprintf('phase = %5.2f deg', angle(x)*180/pi) } \]

mag = 17.32 A    phase = -60.00 deg

Thus, \[|I_a| = 17.32 A\]

The phase sequence \[a \rightarrow b \rightarrow c\] implies the phase diagram below.

From (10.59)

\[ P_{\text{total}} = \sqrt{3} |V_{ab}| |I_a| (\text{load pf}) \]
\[ = \sqrt{3} \times 100 \times 17.32 \times \cos 30^\circ = 2,598 \text{ w} \]

b.

The wattmeter reads the product \[V_{ab} \times I_c\] where \[I_c\] is \[240^\circ\] behind \[I_a\] as shown on the phasor diagram. Then, the wattmeter reading is

\[ P_{\text{wattmeter}} = V_{ab} \times I_c = 100 \angle 0^\circ \times 10 \sqrt{3} \times \cos (-60^\circ - 240^\circ) \]
\[ = 100 \times 17.32 \times \cos (-300^\circ) = 866 \text{ w} \]

and, as expected, this value is on-third of the total power.
2. The single-phase equivalent circuit is shown below where

\[ R = 0.403 \, \Omega /1000 \, \text{ft} \times 1500 \, \text{ft} = 0.605 \, \Omega \]

and thus

\[ Z_{\text{LINE}} = 0.605 + j0.215 \]

Also,

\[ I_{\text{lamp}1} = I_{\text{lamp}2} = \frac{P_{\text{rated}}}{V_{\text{rated}}} = \frac{500}{120} = 4.17 \, \text{A} \]

We recall that for a single phase system the real power is given by
\[ P_{\text{real}} = |V_{\text{RMS}}| |I_{\text{RMS}}| \cos \theta \]

where \( \cos \theta = \text{pf} \)

Then, we find the motor current \( I_M \) in terms of the motor voltage \( V_M \) as

\[ |I_M| = \frac{5000 / 3}{0.8 |V_M|} = \frac{2083}{|V_M|} \]

and since \( \cos^{-1} 0.8 = -36.9^\circ \) (lagging pf), the motor current \( I_M \) is expressed as

\[ I_M = \frac{2083}{V_M} \angle -36.9^\circ = \frac{I}{V_M} (1666 - j1251) \]

The total current is

\[ I_{\text{total}} = I_{\text{lamp1}} + I_{\text{lamp2}} + I_M = 2 \times 4.17 + \frac{I}{V_M} (1666 - j1251) = \frac{I}{V_M} (8.34V_M + 1666 - j1251) \]

and the voltage drop across the 1500 ft line is

\[ V_{\text{line}} = I_{\text{total}} \cdot Z_{\text{line}} = \frac{I}{V_M} (8.34V_M + 1666 - j1251) \cdot (0.605 + j0.215) \]

\[ = \frac{I}{V_M} (5.05V_M + j1.79V_M + 1008 + j358.2 - j756.9 + 269.0) \]

\[ = \frac{I}{V_M} [(5.05V_M + 1277) + j(1.79V_M - 398.7)] \]

Next,

\[ V_{an} = 120 \angle 0^\circ = V_{\text{line}} + V_M = \frac{I}{V_M} [(5.05V_M + 1277) + j(1.79V_M - 398.7)] + V_M \]

or

\[ 120V_M = [(5.05V_M + 1277) + j(1.79V_M - 398.7)] + V_M^2 \]

or

\[ V_M^2 - (114.95 - j1.79)V_M + (1277 - j398.7) = 0 \]

We solve this quadratic equation with the following MATLAB code:

\[
p=[1 \ 114.95-1.79j \ 1277-398.7j]; \text{roots}(p)
\]

\[
\text{ans} = \\
1.0000e+002 \ast \\
1.0260 + 0.0238i \\
0.1235 - 0.0417i
\]
Then, $V_{M1} = 102.6 + j2.39 = 102.63\angle1.33^\circ$ and $V_{M2} = 12.35 - j4.17 = 13.4\angle-18.66^\circ$. Of these, the value of $V_{M2}$ is unrealistic and thus it is rejected.

The positive phase angle in $V_{M1}$ is a result of the fact that a motor is an inductive load. But since an inductive load has a lagging power factor, we denote this line-to-neutral of line-to-ground voltage with a negative angle, that is,

$$V_M = V_{load} = 102.63\angle-1.33^\circ V$$

The magnitude of the line-to-line voltage is

$$|V_{l-l}| = \sqrt{3} \times V_M = \sqrt{3} \times 102.63 = 177.76 V$$
Appendix A

Introduction to MATLAB®

This appendix serves as an introduction to the basic MATLAB commands and functions, procedures for naming and saving the user generated files, comment lines, access to MATLAB’s Editor/Debugger, finding the roots of a polynomial, and making plots. Several examples are provided with detailed explanations.

A.1 MATLAB® and Simulink®

MATLAB® and Simulink® are products of The MathWorks, Inc™. These are two outstanding software packages for scientific and engineering computations and are used in educational institutions and in industries including automotive, aerospace, electronics, telecommunications, and environmental applications. MATLAB enables us to solve many advanced numerical problems fast and efficiently. Simulink is a block diagram tool used for modeling and simulating dynamic systems such as controls, signal processing, and communications. In this appendix we will discuss MATLAB only.

A.2 Command Window

To distinguish the screen displays from the user commands, important terms, and MATLAB functions, we will use the following conventions:

Click: Click the left button of the mouse

Courier Font: Screen displays

Helvetica Font: User inputs at MATLAB’s command window prompt >> or EDU>>*

Helvetica Bold: MATLAB functions

Times Bold Italic: Important terms and facts, notes and file names

When we first start MATLAB, we see the toolbar on top of the command screen and the prompt EDU>>. This prompt is displayed also after execution of a command; MATLAB now waits for a new command from the user. It is highly recommended that we use the Editor/Debugger to write our program, save it, and return to the command screen to execute the program as explained below.

To use the Editor/Debugger:

1. From the File menu on the toolbar, we choose New and click on M-File. This takes us to the Edi-

* EDU>> is the MATLAB prompt in the Student Version
tor Window where we can type our code (list of statements) for a new file, or open a previously saved file. We must save our program with a file name which starts with a letter. Important! MATLAB is case sensitive, that is, it distinguishes between upper- and lower-case letters. Thus, \( t \) and \( T \) are two different letters in MATLAB language. The files that we create are saved with the file name we use and the extension .m; for example, myfile01.m. It is a good practice to save the code in a file name that is descriptive of our code content. For instance, if the code performs some matrix operations, we ought to name and save that file as matrices01.m or any other similar name. We should also use a floppy disk to backup our files.

2. Once the code is written and saved as an m-file, we may exit the Editor/Debugger window by clicking on Exit Editor/Debugger of the File menu. MATLAB then returns to the command window.

3. To execute a program, we type the file name without the .m extension at the >> prompt; then, we press <enter> and observe the execution and the values obtained from it. If we have saved our file in drive a or any other drive, we must make sure that it is added to the desired directory in MATLAB’s search path. The MATLAB User’s Guide provides more information on this topic.

Henceforth, it will be understood that each input command is typed after the >> prompt and followed by the <enter> key.

The command help matlab\iofun will display input/output information. To get help with other MATLAB topics, we can type help followed by any topic from the displayed menu. For example, to get information on graphics, we type help matlab\graphics. The MATLAB User’s Guide contains numerous help topics.

To appreciate MATLAB’s capabilities, we type demo and we see the MATLAB Demos menu. We can do this periodically to become familiar with them. Whenever we want to return to the command window, we click on the Close button.

When we are done and want to leave MATLAB, we type quit or exit. But if we want to clear all previous values, variables, and equations without exiting, we should use the command clear. This command erases everything; it is like exiting MATLAB and starting it again. The command clc clears the screen but MATLAB still remembers all values, variables and equations that we have already used. In other words, if we want to clear all previously entered commands, leaving only the >> prompt on the upper left of the screen, we use the clc command.

All text after the % (percent) symbol is interpreted as a comment line by MATLAB, and thus it is ignored during the execution of a program. A comment can be typed on the same line as the function or command or as a separate line. For instance,

\[
\text{conv}(p,q) \quad \% \text{performs multiplication of polynomials } p \text{ and } q.
\]

\[
% \text{The next statement performs partial fraction expansion of } p(x) / q(x)
\]

are both correct.
One of the most powerful features of MATLAB is the ability to do computations involving complex numbers. We can use either $i$ or $j$ to denote the imaginary part of a complex number, such as $3-4i$ or $3-4j$. For example, the statement
\[
z = 3 - 4j
\]
displays
\[
z = 3.0000 - 4.0000i
\]
In the above example, a multiplication (*) sign between 4 and $j$ was not necessary because the complex number consists of numerical constants. However, if the imaginary part is a function, or variable such as $\cos(x)$, we must use the multiplication sign, that is, we must type $\cos(x)*j$ or $j*\cos(x)$ for the imaginary part of the complex number.

### A.3 Roots of Polynomials

In MATLAB, a polynomial is expressed as a row vector of the form $[a_n \ a_{n-1} \ \ldots \ a_2 \ a_1 \ a_0]$. These are the coefficients of the polynomial in descending order. We must include terms whose coefficients are zero.

We find the roots of any polynomial with the `roots(p)` function; $p$ is a row vector containing the polynomial coefficients in descending order.

**Example A.1**

Find the roots of the polynomial

\[
p_1(x) = x^4 - 10x^3 + 35x^2 - 50x + 24
\]

**Solution:**

The roots are found with the following two statements where we have denoted the polynomial as $p_1$, and the roots as `roots_p1`.

\[
p1 = [1 \ -10 \ 35 \ -50 \ 24] \quad \% \ Specify \ and \ display \ the \ coefficients \ of \ p1(x)
\]
\[
p1 =
\begin{array}{cccc}
1 & -10 & 35 & -50 & 24 \\
\end{array}
\]
\[
roots_p1 = roots(p1) \quad \% \ Find \ the \ roots \ of \ p1(x)
\]
\[
roots_p1 =
\begin{array}{cccc}
4.0000 \\
3.0000 \\
2.0000 \\
\end{array}
\]
1.0000

We observe that MATLAB displays the polynomial coefficients as a row vector, and the roots as a column vector.

**Example A.2**

Find the roots of the polynomial

\[ p_2(x) = x^5 - 7x^4 + 16x^2 + 25x + 52 \]

**Solution:**

There is no cube term; therefore, we must enter zero as its coefficient. The roots are found with the statements below, where we have defined the polynomial as \( p2 \), and the roots of this polynomial as \( \text{roots}_p2 \). The result indicates that this polynomial has three real roots, and two complex roots. Of course, complex roots always occur in *complex conjugate* pairs.

\[
p2=[1 \hspace{1em} -7 \hspace{1em} 0 \hspace{1em} 16 \hspace{1em} 25 \hspace{1em} 52]
\]

\[
p2 =
\begin{bmatrix}
1 & -7 & 0 & 16 & 25 & 52
\end{bmatrix}
\]

\[
\text{roots}_p2=\text{roots}(p2)
\]

\[
\text{roots}_p2 =
\begin{bmatrix}
6.5014 \\
2.7428 \\
-1.5711 \\
-0.3366 + 1.3202i \\
-0.3366 - 1.3202i
\end{bmatrix}
\]

**A.4 Polynomial Construction from Known Roots**

We can compute the coefficients of a polynomial, from a given set of roots, with the `poly(r)` function where \( r \) is a row vector containing the roots.

**Example A.3**

It is known that the roots of a polynomial are \( 1, 2, 3, \) and \( 4 \). Compute the coefficients of this polynomial.

---

* By definition, the conjugate of a complex number \( A = a + jb \) is \( A^* = a - jb \)
Solution:

We first define a row vector, say \( r_3 \), with the given roots as elements of this vector; then, we find the coefficients with the \texttt{poly(r)} function as shown below.

\[
r_3 = [1 \ 2 \ 3 \ 4] \quad \% \ \text{Specify the roots of the polynomial}
\]

\[
r_3 = \begin{bmatrix} 1 & 2 & 3 & 4 \end{bmatrix}
\]

\[
poly_r_3 = \text{poly}(r_3) \quad \% \ \text{Find the polynomial coefficients}
\]

\[
poly_r_3 = \begin{bmatrix} 1 & -10 & 35 & -50 & 24 \end{bmatrix}
\]

We observe that these are the coefficients of the polynomial \( p_1(x) \) of Example A.1.

**Example A.4**

It is known that the roots of a polynomial are \(-1, -2, -3, 4+5j \) and \( 4-5j \). Find the coefficients of this polynomial.

**Solution:**

We form a row vector, say \( r_4 \), with the given roots, and we find the polynomial coefficients with the \texttt{poly(r)} function as shown below.

\[
r_4 = [-1 \ -2 \ -3 \ -4+5j \ -4-5j]
\]

\[
r_4 = \begin{bmatrix}
-1.0000 & -2.0000 & -3.0000 & -4.0000+ 5.0000i \\
-4.0000- 5.0000i \\
\end{bmatrix}
\]

\[
poly_r_4 = \text{poly}(r_4)
\]

\[
poly_r_4 = \begin{bmatrix} 1 & 14 & 100 & 340 & 499 & 246 \end{bmatrix}
\]

Therefore, the polynomial is

\[
p_4(x) = x^5 + 14x^4 + 100x^3 + 340x^2 + 499x + 246
\]
A.5 Evaluation of a Polynomial at Specified Values

The \texttt{polyval(p,x)} function evaluates a polynomial $p(x)$ at some specified value of the independent variable $x$.

Example A.5

Evaluate the polynomial

$$p_5(x) = x^6 - 3x^5 + 5x^3 - 4x^2 + 3x + 2$$ \hspace{1cm} (A.1)

at $x = -3$.

Solution:

$p5=[1\ -3\ 0\ 5\ -4\ 3\ 2]$; \% These are the coefficients

\% The semicolon (;) after the right bracket suppresses the display of the row vector
\% that contains the coefficients of $p5$.
\%
val_minus3=polyval(p5, -3) \% Evaluate $p5$ at $x=-3$; no semicolon is used here
\% because we want the answer to be displayed

val_minus3 =

\hspace{1cm} 1280

Other MATLAB functions used with polynomials are the following:

\texttt{conv(a,b)} – multiplies two polynomials \texttt{a} and \texttt{b}

\texttt{[q,r]=deconv(c,d)} – divides polynomial \texttt{c} by polynomial \texttt{d} and displays the quotient \texttt{q} and remainder \texttt{r}.

\texttt{polyder(p)} – produces the coefficients of the derivative of a polynomial \texttt{p}.

Example A.6

Let

$$p_1 = x^5 - 3x^4 + 5x^2 + 7x + 9$$

and

$$p_2 = 2x^6 - 8x^4 + 4x^2 + 10x + 12$$

Compute the product $p_1 \cdot p_2$ using the \texttt{conv(a,b)} function.
Evaluation of a Polynomial at Specified Values

Solution:

\[ p_1 = [1 \ -3 \ 0 \ 5 \ 7 \ 9]; \quad \text{\% The coefficients of } p_1 \]
\[ p_2 = [2 \ 0 \ -8 \ 0 \ 4 \ 10 \ 12]; \quad \text{\% The coefficients of } p_2 \]
\[ p_1p2 = \text{conv}(p1,p2) \quad \text{\% Multiply } p_1 \text{ by } p_2 \text{ to compute coefficients of the product } p_1p2 \]
\[ p1p2 = \]
\[ 2 \ -6 \ -8 \ 34 \ 18 \ -24 \ -74 \ -88 \ 78 \ 166 \ 174 \ 108 \]

Therefore,

\[ p_1 \cdot p_2 = 2x^{11} - 6x^{10} - 8x^9 + 34x^8 + 18x^7 - 24x^6 \]
\[ -74x^5 - 88x^4 + 78x^3 + 166x^2 + 174x + 108 \]

Example A.7

Let

\[ p_3 = x^7 - 3x^5 + 5x^3 + 7x + 9 \]
and

\[ p_4 = 2x^6 - 8x^5 + 4x^2 + 10x + 12 \]

Compute the quotient \( p_3/p_4 \) using the \([q,r]=\text{deconv}(c,d)\) function.

Solution:

\% It is permissible to write two or more statements in one line separated by semicolons
\[ p3=[1 \ 0 \ -3 \ 0 \ 5 \ 7 \ 9]; \quad p4=[2 \ -8 \ 0 \ 0 \ 4 \ 10 \ 12]; \quad [q,r]=\text{deconv}(p3,p4) \]
\[ q = \]
\[ 0.5000 \]
\[ r = \]
\[ 0 \quad 4 \quad -3 \quad 0 \quad 3 \quad 2 \quad 3 \]

Therefore,

\[ q = 0.5 \quad r = 4x^5 - 3x^4 + 3x^2 + 2x + 3 \]
Example A.8

Let

\[ p_5 = 2x^6 - 8x^4 + 4x^2 + 10x + 12 \]

Compute the derivative \( \frac{d}{dx} p_5 \) using the `polyder(p)` function.

Solution:

\[
p5=[2 0 -8 0 4 10 12]; \quad \% \text{The coefficients of } p5
\]
\[
der_p5=polyder(p5) \quad \% \text{Compute the coefficients of the derivative of } p5
\]
\[
der_p5 =
\begin{bmatrix}
12 & 0 & -32 & 0 & 8 & 10
\end{bmatrix}
\]

Therefore,

\[
\frac{d}{dx} p_5 = 12x^5 - 32x^3 + 4x^2 + 8x + 10
\]

A.6 Rational Polynomials

*Rational Polynomials* are those which can be expressed in ratio form, that is, as

\[
R(x) = \frac{\text{Num}(x)}{\text{Den}(x)} = \frac{b_n x^n + b_{n-1} x^{n-1} + b_{n-2} x^{n-2} + \ldots + b_1 x + b_0}{a_m x^m + a_{m-1} x^{m-1} + a_{m-2} x^{m-2} + \ldots + a_1 x + a_0}
\]  

(A.2)

where some of the terms in the numerator and/or denominator may be zero. We can find the roots of the numerator and denominator with the `roots(p)` function as before.

As noted in the comment line of Example A.7, we can write MATLAB statements in one line, if we separate them by commas or semicolons. *Commas will display the results whereas semicolons will suppress the display.*

Example A.9

Let

\[
R(x) = \frac{P_{\text{num}}}{P_{\text{den}}} = \frac{x^5 - 3x^4 + 5x^2 + 7x + 9}{x^6 - 4x^4 + 2x^2 + 5x + 6}
\]

Express the numerator and denominator in factored form, using the `roots(p)` function.
Solution:

\[
\text{num} = [1\ -3\ 0\ 5\ 7\ 9];\ \text{den} = [1\ 0\ -4\ 0\ 2\ 5\ 6];
\]
% Do not display num and den coefficients
\[
\text{roots\_num} = \text{roots(num)},\ \text{roots\_den} = \text{roots(den)}
\]
% Display num and den roots

\[
\text{roots\_num} =
\begin{array}{l}
2.4186 + 1.0712i \\
-0.3370 + 0.9961i \\
\end{array}
\]
\[
\text{roots\_den} =
\begin{array}{l}
1.6760 + 0.4922i \\
-0.2108 + 0.9870i \\
\end{array}
\]

As expected, the complex roots occur in complex conjugate pairs.

For the numerator, we have the factored form

\[
p_{\text{num}} = (x-2.4186-j1.0712)(x-2.4186+j1.0712)(x + 1.1633)
\]
\[
(x + 0.3370-j0.9961)(x + 0.3370+j0.9961)
\]

and for the denominator, we have

\[
p_{\text{den}} = (x-1.6760-j0.4922)(x-1.6760+j0.4922)(x + 1.9304)
\]
\[
(x + 0.2108-j0.9870)(x + 0.2108+j0.9870)(x + 1.0000)
\]

We can also express the numerator and denominator of this rational function as a combination of linear and quadratic factors. We recall that, in a quadratic equation of the form \(x^2 + bx + c = 0\) whose roots are \(x_1\) and \(x_2\), the negative sum of the roots is equal to the coefficient \(b\) of the \(x\) term, that is, \(-x_1 + x_2 = b\), while the product of the roots is equal to the constant term \(c\), that is, \(x_1 \cdot x_2 = c\). Accordingly, we form the coefficient \(b\) by addition of the complex conjugate roots and this is done by inspection; then we multiply the complex conjugate roots to obtain the constant term \(c\) using MATLAB as follows:

\[
(2.4186 + 1.0712i)*(2.4186-1.0712i)
\]
\[
\text{ans} = 6.9971
\]
\[
(-0.3370+ 0.9961i)*(-0.3370-0.9961i)
\]
\[
\text{ans} = 1.1058
\]
\[
(1.6760+ 0.4922i)*(1.6760-0.4922i)
\]
\[
\text{ans} = 3.0512
\]
(-0.2108 + 0.9870i)*(-0.2108 - 0.9870i)

ans = 1.0186

Thus,

\[ R(x) = \frac{p_{num}}{p_{den}} = \frac{(x^2 - 4.8372x + 6.9971)(x^2 + 0.6740x + 1.1058)(x + 1.1633)}{(x^2 - 3.3520x + 3.0512)(x^2 + 0.4216x + 1.0186)(x + 1.0000)(x + 1.9304)} \]

We can check this result with MATLAB’s Symbolic Math Toolbox which is a collection of tools (functions) used in solving symbolic expressions. They are discussed in detail in MATLAB’s Users Manual. For the present, our interest is in using the collect(s) function that is used to multiply two or more symbolic expressions to obtain the result in polynomial form. We must remember that the conv(p,q) function is used with numeric expressions only, that is, polynomial coefficients.

Before using a symbolic expression, we must create one or more symbolic variables such as x, y, t, and so on. For our example, we use the following code:

```matlab
syms x  % Define a symbolic variable and use collect(s) to express numerator in polynomial form
collect((x^2 - 4.8372*x + 6.9971)*(x^2 + 0.6740*x + 1.1058)*(x + 1.1633))
```

```
ans =
x^5-29999/10000*x^4-1323/3125000*x^3+7813277909/1562500000000*x^2+4500454743147/5000000000000
```

and if we simplify this, we find that is the same as the numerator of the given rational expression in polynomial form. We can use the same procedure to verify the denominator.

### A.7 Using MATLAB to Make Plots

Quite often, we want to plot a set of ordered pairs. This is a very easy task with the MATLAB plot(x,y) command that plots y versus x. Here, x is the horizontal axis (abscissa) and y is the vertical axis (ordinate).

**Example A.10**

Consider the electric circuit of Figure A.1, where the radian frequency \( \omega \) (radians/second) of the applied voltage was varied from 300 to 3000 in steps of 100 radians/second, while the amplitude was held constant. The ammeter readings were then recorded for each frequency. The magnitude of the impedance \( |Z| \) was computed as \( |Z| = |V/A| \) and the data were tabulated on Table A.1.
Plot the magnitude of the impedance, that is, $|Z|$ versus radian frequency $\omega$.

**Solution:**

We cannot type $\omega$ (omega) in the MATLAB command window, so we will use the English letter $w$ instead.

If a statement, or a row vector is too long to fit in one line, it can be continued to the next line by typing three or more periods, then pressing `<enter>` to start a new line, and continue to enter data. This is illustrated below for the data of $w$ and $z$. Also, as mentioned before, we use the semicolon (;) to suppress the display of numbers that we do not care to see on the screen.

The data are entered as follows:
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w=[300 400 500 600 700 800 900 1000 1100 1200 1300 1400 1500 1600 1700 1800 1900 .... 
2000 2100 2200 2300 2400 2500 2600 2700 2800 2900 3000];
%
z=[39.339 52.789 71.104 97.665 140.437 222.182 436.056 ....
1014.938 469.830 266.032 187.052 145.751 120.353 103.111 ....
90.603 81.088 73.588 67.513 62.481 58.240 54.611 51.468 ....
48.717 46.286 44.122 42.182 40.432 38.845];

Of course, if we want to see the values of \(w\) or \(z\) or both, we simply type \(w\) or \(z\), and we press \(<\text{enter}>\). To plot \(z\) (y-axis) versus \(w\) (x-axis), we use the \textbf{plot}(x,y) \textbf{command}. For this example, we use \textbf{plot(w,z)}. When this command is executed, MATLAB displays the plot on MATLAB’s \textit{graph screen}. This plot is shown in Figure A.2.

![Plot of impedance \(|z|\) versus frequency \(\omega\) for Example A.10](image)

This plot is referred to as the \textit{amplitude frequency response} of the circuit.

To return to the command window, we press any key, or from the \textit{Window} pull-down menu, we select \textit{MATLAB Command Window}. To see the graph again, we click on the Window pull-down menu, and we select \textit{Figure}.
We can make the above, or any plot, more presentable with the following commands:

grid on: This command adds grid lines to the plot. The grid off command removes the grid. The command grid toggles them, that is, changes from off to on or vice versa. The default* is off.

box off: This command removes the box (the solid lines which enclose the plot), and box on restores the box. The command box toggles them. The default is on.

title('string'): This command adds a line of the text string (label) at the top of the plot.

xlabel('string') and ylabel('string') are used to label the x- and y-axis respectively.

The amplitude frequency response is usually represented with the x-axis in a logarithmic scale. We can use the semilogx(x,y) command that is similar to the plot(x,y) command, except that the x-axis is represented as a log scale, and the y-axis as a linear scale. Likewise, the semilogy(x,y) command is similar to the plot(x,y) command, except that the y-axis is represented as a log scale, and the x-axis as a linear scale. The loglog(x,y) command uses logarithmic scales for both axes.

Throughout this text it will be understood that \( \log \) is the common (base 10) logarithm, and \( \ln \) is the natural (base e) logarithm. We must remember, however, the function \( \log(x) \) in MATLAB is the natural logarithm, whereas the common logarithm is expressed as \( \log_{10}(x) \), and the logarithm to the base 2 as \( \log_2(x) \).

Let us now redraw the plot with the above options by adding the following statements:

semilogx(w,z); grid; % Replaces the plot(w,z) command

After execution of these commands, our plot is as shown in Figure A.3.

If the y-axis represents power, voltage or current, the x-axis of the frequency response is more often shown in a logarithmic scale, and the y-axis in dB (decibels). The decibel unit is defined in Chapter 4.

---

* A default is a particular value for a variable that is assigned automatically by an operating system and remains in effect unless canceled or overridden by the operator.
To display the voltage $v$ in a dB scale on the y-axis, we add the relation $dB = 20 \cdot \log_{10}(v)$, and we replace the `semilogx(w,z)` command with `semilogx(w,dB)`.

The command `gtext('string')` \* switches to the current Figure Window, and displays a cross-hair that can be moved around with the mouse. For instance, we can use the command `gtext('Impedance $|Z|$ versus Frequency')`, and this will place a cross-hair in the Figure window. Then, using the mouse, we can move the cross-hair to the position where we want our label to begin, and we press <enter>.

The command `text(x,y,'string')` is similar to `gtext('string')`. It places a label on a plot in some specific location specified by $x$ and $y$, and `string` is the label which we want to place at that location. We will illustrate its use with the following example that plots a 3-phase sinusoidal waveform.

The first line of the code below has the form

\* With MATLAB Versions 6 and 7 we can add text, lines and arrows directly into the graph using the tools provided on the Figure Window.
linspace(first_value, last_value, number_of_values)

This function specifies \textit{the number of data points} but not the increments between data points. An alternate function is

\texttt{x=first: increment: last}

and this specifies \textit{the increments between points} but not the number of data points.

The code for the 3-phase plot is as follows:

\begin{verbatim}
x=linspace(0, 2*pi, 60); % pi is a built-in function in MATLAB; % we could have used x=0:0.02*pi:2*pi or x = (0: 0.02: 2)*pi instead;
y=sin(x); u=sin(x+2*pi/3); v=sin(x+4*pi/3);
plot(x,y,x,u,x,v); % The x-axis must be specified for each function grid on, box on, % turn grid and axes box on
text(0.75, 0.65, 'sin(x)'); text(2.85, 0.65, 'sin(x+2*pi/3)'); text(4.95, 0.65, 'sin(x+4*pi/3)')
\end{verbatim}

These three waveforms are shown on the same plot of Figure A.4.

![Figure A.4. Three-phase waveforms](image-url)
In our previous examples, we did not specify line styles, markers, and colors for our plots. However, MATLAB allows us to specify various line types, plot symbols, and colors. These, or a combination of these, can be added with the `plot(x,y,s)` command, where `s` is a character string containing one or more characters shown on the three columns of Table A.2. MATLAB has no default color; it starts with blue and cycles through the first seven colors listed in Table A.2 for each additional line in the plot. Also, there is no default marker; no markers are drawn unless they are selected. The default line is the solid line.

### TABLE A.2 Styles, colors, and markets used in MATLAB

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Color</th>
<th>Symbol</th>
<th>Marker</th>
<th>Symbol</th>
<th>Line Style</th>
</tr>
</thead>
<tbody>
<tr>
<td>b</td>
<td>blue</td>
<td>.</td>
<td>point</td>
<td>—</td>
<td>solid line</td>
</tr>
<tr>
<td>g</td>
<td>green</td>
<td>o</td>
<td>circle</td>
<td>:</td>
<td>dotted line</td>
</tr>
<tr>
<td>r</td>
<td>red</td>
<td>x</td>
<td>x-mark</td>
<td>.</td>
<td>dash-dot line</td>
</tr>
<tr>
<td>c</td>
<td>cyan</td>
<td>+</td>
<td>plus</td>
<td>——</td>
<td>dashed line</td>
</tr>
<tr>
<td>m</td>
<td>magenta</td>
<td>*</td>
<td>star</td>
<td></td>
<td></td>
</tr>
<tr>
<td>y</td>
<td>yellow</td>
<td>s</td>
<td>square</td>
<td></td>
<td></td>
</tr>
<tr>
<td>k</td>
<td>black</td>
<td>d</td>
<td>diamond</td>
<td></td>
<td></td>
</tr>
<tr>
<td>w</td>
<td>white</td>
<td>▽</td>
<td>triangle down</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>▽</td>
<td>triangle up</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>▽</td>
<td>triangle left</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>▽</td>
<td>triangle right</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>p</td>
<td>pentagram</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>h</td>
<td>hexagram</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

For example, `plot(x,y,'m*:')` plots a magenta dotted line with a star at each data point, and `plot(x,y,'rs')` plots a red square at each data point, but does not draw any line because no line was selected. If we want to connect the data points with a solid line, we must type `plot(x,y,'rs-')`. For additional information we can type `help plot` in MATLAB’s command screen.

The plots we have discussed thus far are two-dimensional, that is, they are drawn on two axes. MATLAB has also a three-dimensional (three-axes) capability and this is discussed next.

The `plot3(x,y,z)` command plots a line in 3-space through the points whose coordinates are the elements of `x`, `y`, and `z`, where `x`, `y`, and `z` are three vectors of the same length.

The general format is `plot3(x_1,y_1,z_1,s_1,x_2,y_2,z_2,s_2,x_3,y_3,z_3,s_3,...)` where `x_n`, `y_n` and `z_n` are vectors or matrices, and `s_n` are strings specifying color, marker symbol, or line style. These strings are the same as those of the two-dimensional plots.
Example A.11

Plot the function

\[ z = -2x^3 + x + 3y^2 - 1 \] \hspace{1cm} (A.3)

Solution:

We arbitrarily choose the interval (length) shown on the code below.

```matlab
x = -10: 0.5: 10; \quad \% \text{Length of vector } x
y = x; \quad \% \text{Length of vector } y \text{ must be same as } x
z = -2.*x.^3 + x + 3.*y.^2 - 1; \quad \% \text{Vector } z \text{ is function of both } x \text{ and } y
plot3(x,y,z); grid
```

The three-dimensional plot is shown in Figure A.5.

![Three-dimensional plot for Example A.11](image)

Figure A.5. Three dimensional plot for Example A.11

In a two-dimensional plot, we can set the limits of the x- and y-axes with the \texttt{axis([xmin xmax ymin ymax])} command. Likewise, in a three-dimensional plot we can set the limits of all three axes

* This statement uses the so called dot multiplication, dot division, and dot exponentiation where the multiplication, division, and exponential operators are preceded by a dot. These operations will be explained in Section A.8.

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with the \texttt{axis([xmin xmax ymin ymax zmin zmax])} command. It must be placed after the \texttt{plot(x,y)} or \texttt{plot3(x,y,z)} commands, or on the same line without first executing the \texttt{plot} command. This must be done for each plot. The three-dimensional \texttt{text(x,y,z,'string')} command will place \textit{string} beginning at the co-ordinate \((x,y,z)\) on the plot.

For three-dimensional plots, \texttt{grid on} and \texttt{box off} are the default states.

We can also use the \texttt{mesh(x,y,z)} command with two vector arguments. These must be defined as \(\text{length}(x) = n\) and \(\text{length}(y) = m\) where \([m,n] = \text{size}(Z)\). In this case, the vertices of the mesh lines are the triples \(\{x(j), y(i), Z(i,j)\}\). We observe that \(x\) corresponds to the columns of \(Z\), and \(y\) corresponds to the rows.

To produce a mesh plot of a function of two variables, say \(z = f(x, y)\), we must first generate the \(X\) and \(Y\) matrices that consist of repeated rows and columns over the range of the variables \(x\) and \(y\). We can generate the matrices \(X\) and \(Y\) with the \texttt{[X,Y]=meshgrid(x,y)} function that creates the matrix \(X\) whose rows are copies of the vector \(x\), and the matrix \(Y\) whose columns are copies of the vector \(y\).

Example A.12

The volume \(V\) of a right circular cone of radius \(r\) and height \(h\) is given by

\[
V = \frac{1}{3} \pi r^2 h
\]  

(A.4)

Plot the volume of the cone as \(r\) and \(h\) vary on the intervals \(0 \leq r \leq 4\) and \(0 \leq h \leq 6\) meters.

Solution:

The volume of the cone is a function of both the radius \(r\) and the height \(h\), that is,

\[
V = f(r, h)
\]

The three-dimensional plot is created with the following MATLAB code where, as in the previous example, in the second line we have used the dot multiplication, dot division, and dot exponentiation. This will be explained in Section A.8.

\[
\begin{align*}
[R,H]= & \text{meshgrid}(0:4, 0:6); & \quad \% \text{Creates } R \text{ and } H \text{ matrices from vectors } r \text{ and } h \\
V= & (pi .* R .^ 2 .* H) ./ 3; & \quad \text{mesh}(R, H, V) \\
exlabel('x-axis, radius r (meters)'); & \text{ylabel('y-axis, altitude h (meters)');} \\
xlabel('z-axis, volume (cubic meters)'); & \text{title('Volume of Right Circular Cone');} \\
\end{align*}
\]

The three-dimensional plot of Figure A.6, shows how the volume of the cone increases as the radius and height are increased.
A.8 Subplots

MATLAB can display up to four windows of different plots on the Figure window using the command `subplot(m,n,p)`. This command divides the window into an $m \times n$ matrix of plotting areas and chooses the $p$th area to be active. No spaces or commas are required between the three integers $m$, $n$ and $p$. The possible combinations are shown in Figure A.7.

We will illustrate the use of the `subplot(m,n,p)` command following the discussion on multiplication, division and exponentiation that follows.

![Figure A.6. Volume of a right circular cone.](image)

This, and the plot of Figure A.5, are rudimentary; MATLAB can generate very sophisticated three-dimensional plots. The MATLAB User’s manual contains more examples.
A.9 Multiplication, Division and Exponentiation

MATLAB recognizes two types of multiplication, division, and exponentiation. These are the matrix multiplication, division, and exponentiation, and the element-by-element multiplication, division, and exponentiation. They are explained in the following paragraphs.

In Section A.2, the arrays \([a \ b \ c \ \ldots]\), such as those that contained the coefficients of polynomials, consisted of one row and multiple columns, and thus are called row vectors. If an array has one column and multiple rows, it is called a column vector. We recall that the elements of a row vector are separated by spaces. To distinguish between row and column vectors, the elements of a column vector must be separated by semicolons. An easier way to construct a column vector is to write it first as a row vector, and then transpose it into a column vector. MATLAB uses the single quotation character (’), to transpose a vector. Thus, a column vector can be written either as \(b=[-1; 3; 6; 11]\) or as \(b=[-1 \ 3 \ 6 \ 11]'\). MATLAB produces the same display with either format as shown below.

\[
b=[-1; 3; 6; 11]\\
b = \\
-1 \\
3 \\
6 \\
11 \\
b=[-1 \ 3 \ 6 \ 11]'\\
b = \\
-1 \\
3 \\
6 \\
11
\]

We will now define Matrix Multiplication and Element-by-Element multiplication.

1. **Matrix Multiplication** (multiplication of row by column vectors)

Let

\[
A = [a_1 \ a_2 \ a_3 \ \ldots \ a_n]
\]

and

\[
B = [b_1 \ b_2 \ b_3 \ \ldots \ b_n]'
\]

be two vectors. We observe that \(A\) is defined as a row vector whereas \(B\) is defined as a column vector, as indicated by the transpose operator (’). Here, multiplication of the row vector \(A\) by the column
Multiplication, Division and Exponentiation

vector B, is performed with the matrix multiplication operator (*). Then,

$$A * B = [a_1 b_1 + a_2 b_2 + a_3 b_3 + \ldots + a_n b_n] = \text{single value} \quad (A.5)$$

For example, if

$$A = [1 \ 2 \ 3 \ 4 \ 5]$$

and

$$B = [-2 \ 6 \ -3 \ 8 \ 7]'$$

the matrix multiplication $A * B$ produces the single value 68, that is,

$$A * B = 1 \times (-2) + 2 \times 6 + 3 \times (-3) + 4 \times 8 + 5 \times 7 = 68$$

and this is verified with MATLAB as

```matlab
A=[1 2 3 4 5]; B=[-2 6 -3 8 7];
A*B
```

ans =

68

Now, let us suppose that both $A$ and $B$ are row vectors, and we attempt to perform a row-by-row multiplication with the following MATLAB statements.

```matlab
A=[1 2 3 4 5]; B=[-2 6 -3 8 7];
A*B
```

When these statements are executed, MATLAB displays the following message:

```matlab
??? Error using ==> *
Inner matrix dimensions must agree.
```

Here, because we have used the matrix multiplication operator (*) in $A * B$, MATLAB expects vector B to be a column vector, not a row vector. It recognizes that B is a row vector, and warns us that we cannot perform this multiplication using the matrix multiplication operator (*). Accordingly, we must perform this type of multiplication with a different operator. This operator is defined below.

2. **Element-by-Element Multiplication** (multiplication of a row vector by another row vector)

Let

$$C = [c_1 \ c_2 \ c_3 \ \ldots \ c_n]$$

and

$$D = [d_1 \ d_2 \ d_3 \ \ldots \ d_n]$$

be two row vectors. Here, multiplication of the row vector $C$ by the row vector $D$ is performed with the **dot multiplication operator** (.*). There is no space between the dot and the multiplication sym-
bol. Thus,

\[
C \cdot D = [c_1d_1 \ c_2d_2 \ c_3d_3 \ \ldots \ c_nd_n]
\]  \hspace{1cm} (A.6)

This product is another row vector with the same number of elements, as the elements of \( C \) and \( D \).

As an example, let

\[
C = [1 \ 2 \ 3 \ 4 \ 5]
\]

and

\[
D = [-2 \ 6 \ -3 \ 8 \ 7]
\]

Dot multiplication of these two row vectors produce the following result.

\[
C \cdot D = 1 \times (-2) \quad 2 \times 6 \quad 3 \times (-3) \quad 4 \times 8 \quad 5 \times 7 = -2 \quad 12 \quad -9 \quad 32 \quad 35
\]

Check with MATLAB:

```matlab
C=[1 2 3 4 5]; \% Vectors C and D must have
D=[-2 6 -3 8 7]; \% same number of elements
C.*D \% We observe that this is a dot multiplication
ans =
    -2    12    -9    32    35
```

Similarly, the division (\(/\)) and exponentiation (\(^\)\)) operators, are used for matrix division and exponentiation, whereas dot division (\(./\)) and dot exponentiation (\(.^\)) are used for element-by-element division and exponentiation.

We must remember that no space is allowed between the dot (\(\cdot\)) and the multiplication, division, and exponentiation operators.

Note: A dot (\(\cdot\)) is never required with the plus (+) and minus (\(-\)) operators.

Example A.13

Write the MATLAB code that produces a simple plot for the waveform defined as

\[
y = f(t) = 3e^{-4t}cos5t - 2e^{-3t}sin2t + \frac{t^2}{t + 1}
\]  \hspace{1cm} (A.7)

in the \( 0 \leq t \leq 5 \) seconds interval.

Solution:

The MATLAB code for this example is as follows:

```matlab
t=0:0.01:5 \% Define t-axis in 0.01 increments
y=3.*exp(-4.*t).*cos(5.*t)-2.*exp(-3.*t).*sin(2.*t) + t.^2 ./ (t+1);
```
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plot(t,y); grid; xlabel('t'); ylabel('y=f(t)'); title('Plot for Example A.13')

Figure A.8 shows the plot for this example.

![Plot for Example A.13](image)

Had we, in this example, defined the time interval starting with a negative value equal to or less than \(-1\), say as \(-3 \leq t \leq 3\), MATLAB would have displayed the following message:

Warning: Divide by zero.

This is because the last term (the rational fraction) of the given expression, is divided by zero when \(t = -1\). To avoid division by zero, we use the special MATLAB function \texttt{eps}, which is a number approximately equal to \(2.2 \times 10^{-16}\). It will be used with the next example.

The command \texttt{axis([xmin xmax ymin ymax])} scales the current plot to the values specified by the arguments \texttt{xmin, xmax, ymin and ymax}. There are no commas between these four arguments. This command must be placed \textit{after} the plot command and must be repeated for each plot.

The following example illustrates the use of the dot multiplication, division, and exponentiation, the \texttt{eps} number, the \texttt{axis([xmin xmax ymin ymax])} command, and also MATLAB’s capability of displaying up to four windows of different plots.

**Example A.14**

Plot the functions

\[ y = \sin^2 x, \quad z = \cos^2 x, \quad w = \sin^2 x \cdot \cos^2 x, \quad v = \sin^2 x / \cos^2 x \]
in the interval $0 \leq x \leq 2\pi$ using 100 data points. Use the `subplot` command to display these functions on four windows on the same graph.

**Solution:**

The MATLAB code to produce the four subplots is as follows:

```matlab
x=linspace(0,2*pi,100); % Interval with 100 data points
y=(sin(x).^ 2);  z=(cos(x).^ 2);
w=y.* z;
v=y./(z+eps); % add eps to avoid division by zero
subplot(221);% upper left of four subplots
plot(x,y);  axis([0 2*pi 0 1]);
title('y=(sinx)^2');
subplot(222); % upper right of four subplots
plot(x,z);  axis([0 2*pi 0 1]);
title('z=(cosx)^2');
subplot(223); % lower left of four subplots
plot(x,w);  axis([0 2*pi 0 0.3]);
title('w=(sinx)^2*(cosx)^2');
subplot(224); % lower right of four subplots
plot(x,v);  axis([0 2*pi 0 400]);
title('v=(sinx)^2/(cosx)^2');
```

These subplots are shown in Figure A.9.

![Figure A.9. Subplots for the functions of Example A.14](image)
The next example illustrates MATLAB’s capabilities with imaginary numbers. We will introduce the \texttt{real(z)} and \texttt{imag(z)} functions that display the real and imaginary parts of the complex quantity $z = x + iy$, the \texttt{abs(z)}, and the \texttt{angle(z)} functions that compute the absolute value (magnitude) and phase angle of the complex quantity $z = x + iy = r\angle \theta$. We will also use the \texttt{polar(theta,r)} function that produces a plot in polar coordinates, where $r$ is the magnitude, \texttt{theta} is the angle in radians, and the \texttt{round(n)} function that rounds a number to its nearest integer.

Example A.15

Consider the electric circuit of Figure A.10.

With the given values of resistance, inductance, and capacitance, the impedance $Z_{ab}$ as a function of the radian frequency $\omega$ can be computed from the following expression:

\begin{equation}
Z_{ab} = Z = 10 + \frac{10^4 - j(10^6/\omega)}{10 + j(0.1 \omega - 10^5/\omega)} \tag{A.8}
\end{equation}

a. Plot $Re\{Z\}$ (the real part of the impedance $Z$) versus frequency $\omega$

b. Plot $Im\{Z\}$ (the imaginary part of the impedance $Z$) versus frequency $\omega$

c. Plot the impedance $Z$ versus frequency $\omega$ in polar coordinates.

Solution:

The MATLAB code below computes the real and imaginary parts of $Z_{ab}$ that is, for simplicity, denoted as $z$, and plots these as two separate graphs (parts a & b). It also produces a polar plot (part c).

```matlab
w=0: 1: 2000;  \% Define interval with one radian interval
z=(10+(10 .^ 4 - j .* 10 .^ 6 ./ (w+eps)) ./ (10 + j .* (0.1 .* w -10. .^ 5./ (w+eps)))));
\%  
\% The first five statements (next two lines) compute and plot Re\{z\}
real_part=real(z);  plot(w,real_part);  grid;
xlabel('radian frequency w');  ylabel('Real part of Z');
\%
```

---

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% The next five statements (next two lines) compute and plot Im\{z\}
imag_part=imag(z);  plot(w,imag_part);  grid;
xlabel('radian frequency w');  ylabel('Imaginary part of Z');
% The last six statements (next six lines) below produce the polar plot of z
mag=abs(z); % Computes |Z|
rndz=round(abs(z)); % Rounds |Z| to read polar plot easier
theta=angle(z); % Computes the phase angle of impedance Z
polar(theta,rndz);  % Angle is the first argument
grid;
ylabel('Polar Plot of Z');

The real, imaginary, and polar plots are shown in Figures A.11, A.12, and A.13 respectively.

Example A.15 clearly illustrates how powerful, fast, accurate, and flexible MATLAB is.

### A.10 Script and Function Files

MATLAB recognizes two types of files: script files and function files. Both types are referred to as m-files since both require the .m extension.

A script file consists of two or more built-in functions such as those we have discussed thus far. Thus, the code for each of the examples we discussed earlier, make up a script file. Generally, a script file is one which was generated and saved as an m-file with an editor such as the MATLAB’s Editor/Debugger.

![Figure A.11. Plot for the real part of the impedance in Example A.15](image-url)
A function file is a user-defined function using MATLAB. We use function files for repetitive tasks. The first line of a function file must contain the word function, followed by the output argument, the equal sign (=), and the input argument enclosed in parentheses. The function name and file name must be the same, but the file name must have the extension .m. For example, the function file consisting of the two lines below

```matlab
function Y = myFunction(X)
    Y = X + 1;
end
```
function y = myfunction(x)
y=x. ^ 3 + cos(3.* x)
is a function file and must be saved as myfunction.m

For the next example, we will use the following MATLAB functions.

**fzero(f,x)** tries to find a zero of a function of one variable, where \( f \) is a string containing the name of a real-valued function of a single real variable. MATLAB searches for a value near a point where the function \( f \) changes sign, and returns that value, or returns **NaN** if the search fails.

**Important:** We must remember that we use **roots(p)** to find the roots of polynomials only, such as those in Examples A.1 and A.2.

**fmin(f,x1,x2)** minimizes a function of one variable. It attempts to return a value of \( x \) where \( f(x) \) is minimum in the interval \( x_1 < x < x_2 \). The string \( f \) contains the name of the function to be minimized.

**Note:** MATLAB does not have a function to maximize a function of one variable, that is, there is no **fmax(f,x1,x2)** function in MATLAB; but since a maximum of \( f(x) \) is equal to a minimum of \( -f(x) \), we can use **fmin(f,x1,x2)** to find both minimum and maximum values of a function.

**fplot(fcn,lims)** plots the function specified by the string \( fcn \) between the x-axis limits specified by \( \text{lims} = [\text{xmin} \ \text{xmax}] \). Using \( \text{lims} = [\text{xmin} \ \text{xmax} \ \text{ymin} \ \text{ymax}] \) also controls the y-axis limits. The string \( fcn \) must be the name of an **m-file** function or a string with variable \( x \).

**Note:** **NaN** (Not-a-Number) is not a function; it is MATLAB’s response to an undefined expression such as \( 0/0, \infty/\infty \), or inability to produce a result as described on the next paragraph. We can avoid division by zero using the **eps** number, that we mentioned earlier.

**Example A.16**

Find the zeros, maxima and minima of the function

\[
f(x) = \frac{1}{(x - 0.1)^2 + 0.01} + \frac{1}{(x - 1.2)^2 + 0.04} - 10
\]

**Solution:**

We first plot this function to observe the approximate zeros, maxima, and minima using the following code.

\[
x=1.5: 0.01: 1.5;
y=1./ ((x-0.1).^ 2 + 0.01) -1./ ((x-1.2).^ 2 + 0.04) -10;
\]

plot(x,y); grid

The plot is shown in Figure A.14.
The roots (zeros) of this function appear to be in the neighborhood of \( x = -0.2 \) and \( x = 0.3 \). The maximum occurs at approximately \( x = 0.1 \) where, approximately, \( y_{\text{max}} = 90 \), and the minimum occurs at approximately \( x = 1.2 \) where, approximately, \( y_{\text{min}} = -34 \).

Next, we define and save \( f(x) \) as the `funczero01.m` function m-file with the following code:

```matlab
function y=funczero01(x)
% Finding the zeros of the function shown below
y=1/((x-0.1)^2+0.01)-1/((x-1.2)^2+0.04)-10;
```

Now, we can use the `fplot(fcn,lims)` command to plot \( f(x) \) as follows.

```matlab
fplot('funczero01', [-1.5 1.5]); grid
```

This plot is shown in Figure A.15. As expected, this plot is identical to the plot of Figure A.14 that was obtained with the `plot(x,y)` command.

---

**Figure A.14. Plot for Example A.16 using the plot command**
We will use the \texttt{fzero(f,x)} function to compute the roots of \( f(x) \) in (A.20) more precisely. The code below must be saved with a file name, and then invoked with that file name.

\begin{verbatim}
x1 = fzero('funczero01', -0.2);
x2 = fzero('funczero01', 0.3);
fprintf('The roots (zeros) of this function are r1= %3.4f', x1);
fprintf(' and r2= %3.4f \n', x2)
\end{verbatim}

MATLAB displays the following:

The roots (zeros) of this function are \( r1= -0.1919 \) and \( r2= 0.3788 \)

Whenever we use the \texttt{fmin(f,x1,x2)} function, we must remember that this function searches for a minimum and it may display the values of local minima\footnote{Local maxima or local minima, are the maximum or minimum values of a function within a restricted range of values in the independent variable. When the entire range is considered, the maxima and minima are considered be to the maximum and minimum values in the entire range in which the function is defined.}, if any, before displaying the function minimum. It is, therefore, advisable to plot the function with either the \texttt{plot(x,y)} or the \texttt{fplot(fcn,lims)} command to find the smallest possible interval within which the function minimum lies. For this example, we specify the range \( 0 \leq x \leq 1.5 \) rather than the interval \(-1.5 \leq x \leq 1.5\).

The minimum of \( f(x) \) is found with the \texttt{fmin(f,x1,x2)} function as follows.

\begin{verbatim}
min_val = fmin('funczero01', 0, 1.5)
\end{verbatim}

\texttt{min_val} = 1.2012
This is the value of $x$ at which $y = f(x)$ is minimum. To find the value of $y$ corresponding to this value of $x$, we substitute it into $f(x)$, that is,

$$x=1.2012; \quad y=1 / ((x-0.1)^2 + 0.01) - 1 / ((x-1.2)^2 + 0.04) - 10$$

$$y = -34.1812$$

To find the maximum value, we must first define a new function $m$-file that will produce $-f(x)$. We define it as follows:

```matlab
function y=minusfunczero01(x)
    % It is used to find maximum value from -f(x)
    y=-(1/((x-0.1)^2+0.01)-1/((x-1.2)^2+0.04)-10);
end
```

We have placed the minus ($-$) sign in front of the right side of the last expression above, so that the maximum value will be displayed. Of course, this is equivalent to the negative of the `funczero01` function.

Now, we execute the following code to get the value of $x$ where the maximum $y = f(x)$ occurs.

```matlab
max_val=fmin('minusfunczero01', 0,1)
```

```
max_val = 0.0999
```

$x=0.0999; \%$ Using this value find the corresponding value of $y$

$$y=1 / ((x-0.1)^2 + 0.01) - 1 / ((x-1.2)^2 + 0.04) - 10$$

$$y = 89.2000$$

### A.11 Display Formats

MATLAB displays the results on the screen in integer format without decimals if the result is an integer number, or in short floating point format with four decimals if it a fractional number. The format displayed has nothing to do with the accuracy in the computations. MATLAB performs all computations with accuracy up to 16 decimal places.

The output format can changed with the `format` command. The available formats can be displayed with the `help format` command as follows:

```matlab
help format
```

FORMAT Set output format.

All computations in MATLAB are done in double precision.

FORMAT may be used to switch between different output display formats as follows:

<table>
<thead>
<tr>
<th>FORMAT</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>Default</td>
<td>Same as SHORT</td>
</tr>
<tr>
<td>SHORT</td>
<td></td>
</tr>
<tr>
<td>FORMATS</td>
<td></td>
</tr>
</tbody>
</table>
FORMAT SHORT Scaled fixed point format with 5 digits.
FORMAT LONG Scaled fixed point format with 15 digits.
FORMAT SHORT E Floating point format with 5 digits.
FORMAT LONG E Floating point format with 15 digits.
FORMAT SHORT G Best of fixed or floating point format with 5 digits.
FORMAT LONG G Best of fixed or floating point format with 15 digits.
FORMAT HEX Hexadecimal format.
FORMAT + The symbols +, - and blank are printed for positive, negative and zero elements. Imaginary parts are ignored.
FORMAT BANK Fixed format for dollars and cents.
FORMAT RAT Approximation by ratio of small integers.

Spacing:
FORMAT COMPACT Suppress extra line-feeds.
FORMAT LOOSE Puts the extra line-feeds back in.

Some examples with different format displays age given below.

```
format short
33.3335  Four decimal digits (default)
format long
33.33333333333334 16 digits
format short e
3.3333e+01  Four decimal digits plus exponent
format short g
33.333  Better of format short or format format short e
format bank
33.33  two decimal digits
format + only + or – or zero are printed
format rat
100/3  rational approximation
```

The disp(X) command displays the array X without printing the array name. If X is a string, the text is displayed.

The fprintf(format,array) command displays and prints both text and arrays. It uses specifiers to indicate where and in which format the values would be displayed and printed. Thus, if %f is used, the values will be displayed and printed in fixed decimal format, and if %e is used, the values will be displayed and printed in scientific notation format. With these commands only the real part of each parameter is processed.
Appendix B

Differential Equations

This appendix is a review of ordinary differential equations. Some definitions, topics, and examples are not applicable to introductory circuit analysis but are included for continuity of the subject, and for reference to more advanced topics in electrical engineering such as state variables. These are denoted with an asterisk and may be skipped.

B.1 Simple Differential Equations

In this section we present two simple examples to show the importance of differential equations in engineering applications.

Example B.1

A 1 F capacitor is being charged by a constant current $I$. Find the voltage $v_C$ across this capacitor as a function of time given that the voltage at some reference time $t = 0$ is $V_0$.

Solution:

It is given that the current, as a function of time, is constant, that is,

$$i_C(t) = I = \text{constant} \quad (B.1)$$

We know that the current and voltage in a capacitor are related by

$$i_C(t) = C \frac{dv_C}{dt} \quad (B.2)$$

and for our example, $C = 1$. Then, by substitution of (B.2) into (B.1) we get

$$\frac{dv_C}{dt} = I$$

By separation of the variables,

$$dv_C = Idt \quad (B.3)$$

and by integrating both sides of (B.3) we get

$$v_C(t) = It + k \quad (B.4)$$

where $k$ represents the constants of integration of both sides.
We can find the value of the constant $k$ by making use of the initial condition, i.e., at $t = 0$, $v_C = V_0$ and (B.4) then becomes

$$V_0 = 0 + k$$  \hspace{1cm} (B.5)

or $k = V_0$, and by substitution into (B.4),

$$v_C(t) = It + V_0$$  \hspace{1cm} (B.6)

This example shows that when a capacitor is charged with a constant current, a linear voltage is produced across the terminals of the capacitor.

**Example B.2**

Find the current $i_L(t)$ through an inductor whose slope at the coordinate $(t, i_L)$ is $cost$ and the current $i_L$ passes through the point $(\pi/2, I)$.

**Solution:**

We are given that

$$\frac{di_L}{dt} = cost$$  \hspace{1cm} (B.7)

By separating the variables we get

$$di_L = costdt$$  \hspace{1cm} (B.8)

and integrating both sides we get

$$i_L(t) = sint + k$$  \hspace{1cm} (B.9)

where $k$ represents the constants of integration of both sides.

We find the value of the constant $k$ by making use of the initial condition. For this example, $\omega = 1$ and thus at $\omega t = t = \pi/2$, $i_L = I$. With these values (B.9) becomes

$$I = sin\frac{\pi}{2} + k$$  \hspace{1cm} (B.10)

or $k = 0$, and by substitution into (B.9),

$$i_L(t) = sint$$  \hspace{1cm} (B.11)
Classification

B.2 Classification

Differential equations are classified by:

1. **Type** – Ordinary or Partial
2. **Order** – The highest order derivative which is included in the differential equation
3. **Degree** – The exponent of the highest power of the highest order derivative after the differential equation has been cleared of any fractions or radicals in the dependent variable and its derivatives

For example, the differential equation

\[
\frac{d^4y}{dx^4} + 5\frac{d^3y}{dx^3} + 6\frac{d^2y}{dx^2} + 3\left(\frac{dy}{dx}\right)^8 + \frac{y^2}{x^3 + 1} = ye^{-2x}
\]

is an ordinary differential equation of order 4 and degree 2.

If the dependent variable \( y \) is a function of only a single variable \( x \), that is, if \( y = f(x) \), the differential equation which relates \( y \) and \( x \) is said to be an ordinary differential equation and it is abbreviated as ODE.

The differential equation

\[
\frac{d^2y}{dt^2} + 3\frac{dy}{dt} + 2 = 5\cos4t
\]

is an ODE with constant coefficients.

The differential equation

\[
x^2 \frac{d^2y}{dt^2} + x \frac{dy}{dt} + (x^2 - n^2) = 0
\]

is an ODE with variable coefficients.

If the dependent variable \( y \) is a function of two or more variables such as \( y = f(x, t) \), where \( x \) and \( t \) are independent variables, the differential equation that relates \( y \), \( x \), and \( t \) is said to be a partial differential equation and it is abbreviated as PDE.

An example of a partial differential equation is the well-known one-dimensional wave equation shown below.

\[
\frac{\partial^2y}{\partial t^2} = a^2\frac{\partial^2y}{\partial x^2}
\]

Most of the electrical engineering problems are solved with ordinary differential equations with constant coefficients; however, partial differential equations provide often quick solutions to some practical applications as illustrated with the following three examples.
Example B.3

The equivalent resistance $R_T$ of three resistors $R_1$, $R_2$, and $R_3$ in parallel is given by

$$\frac{1}{R_T} = \frac{1}{R_1} + \frac{1}{R_2} + \frac{1}{R_3}$$

Given that initially $R_1 = 5 \, \Omega$, $R_2 = 20 \, \Omega$, and $R_3 = 4 \, \Omega$ compute the change in $R_T$ if $R_2$ is increased by 10% and $R_3$ is decreased by 5% while $R_1$ does not change.

Solution:

The initial value of the equivalent resistance is $R_T = 5 \parallel 20 \parallel 4 = 2 \, \Omega$.

Now, we treat $R_2$ and $R_3$ as constants and differentiating $R_T$ with respect to $R_1$ we get

$$-\frac{1}{R_T^2} \frac{\partial R_T}{\partial R_1} = -\frac{1}{R_1^2} \text{ or } \frac{\partial R_T}{\partial R_1} = \left( \frac{R_T}{R_1} \right)^2$$

Similarly,

$$\frac{\partial R_T}{\partial R_2} = \left( \frac{R_T}{R_2} \right)^2 \text{ and } \frac{\partial R_T}{\partial R_3} = \left( \frac{R_T}{R_3} \right)^2$$

and the total differential $dR_T$ is

$$dR_T = \frac{\partial R_T}{\partial R_1} dR_1 + \frac{\partial R_T}{\partial R_2} dR_2 + \frac{\partial R_T}{\partial R_3} dR_3 = \left( \frac{R_T}{R_1} \right)^2 dR_1 + \left( \frac{R_T}{R_2} \right)^2 dR_2 + \left( \frac{R_T}{R_3} \right)^2 dR_3$$

By substitution of the given numerical values we get

$$dR_T = \left( \frac{2}{5} \right)^2 (0) + \left( \frac{2}{20} \right)^2 (2) + \left( \frac{2}{4} \right)^2 (0.2) = 0.02 - 0.05 = -0.03$$

Therefore, the equivalent resistance decreases by 3%.

Example B.4

In a series RC circuit that is excited by a sinusoidal voltage, the magnitude of the impedance $Z$ is computed from $Z = \sqrt{R^2 + X_C^2}$. Initially, $R = 4 \, \Omega$ and $X_C = 3 \, \Omega$. Find the change in the impedance $Z$ if the resistance $R$ is increased by 0.25 Ω (6.25%) and the capacitive reactance $X_C$ is decreased by 0.125 Ω (–4.167%).
Solution:

We will first find the partial derivatives \( \frac{\partial Z}{\partial R} \) and \( \frac{\partial Z}{\partial X_C} \); then we compute the change in impedance from the total differential \( dZ \). Thus,

\[
\frac{\partial Z}{\partial R} = \frac{R}{\sqrt{R^2 + X_C^2}} \quad \text{and} \quad \frac{\partial Z}{\partial X_C} = \frac{X_C}{\sqrt{R^2 + X_C^2}}
\]

and

\[
dZ = \frac{\partial Z}{\partial R} \, dR + \frac{\partial Z}{\partial X_C} \, dX_C = \frac{R \, dR + X_C \, dX_C}{\sqrt{R^2 + X_C^2}}
\]

and by substitution of the given values

\[
dZ = \frac{4 \, (0.25) + 3 \, (-0.125)}{\sqrt{4^2 + 3^2}} = \frac{1 - 0.375}{5} = 0.125
\]

Therefore, if \( R \) increases by 6.25% and \( X_C \) decreases by 4.167%, the impedance \( Z \) increases by 4.167%.

Example B.5

A light bulb is rated at 120 volts and 75 watts. If the voltage decreases by 5 volts and the resistance of the bulb is increased by 8 \( \Omega \), by how much will the power change?

Solution:

At \( V = 120 \) volts and \( P = 75 \) watts, the bulb resistance is

\[
R = \frac{V^2}{P} = \frac{120^2}{75} = 192 \, \Omega
\]

and since

\[
P = \frac{V^2}{R} \quad \text{then} \quad \frac{\partial P}{\partial V} = \frac{2V}{R} \quad \text{and} \quad \frac{\partial P}{\partial R} = -\frac{V^2}{R^2}
\]

and the total differential is

\[
dP = \frac{\partial P}{\partial V} \, dV + \frac{\partial P}{\partial R} \, dR = \frac{2V}{R} \, dV - \frac{V^2}{R^2} \, dR
\]

\[
= \frac{2(120)}{192} (-5) - \frac{120^2}{192^2} (8) = -9.375
\]

That is, the power will decrease by 9.375 watts.
B.3 Solutions of Ordinary Differential Equations (ODE)

A function \( y = f(x) \) is a solution of a differential equation if the latter is satisfied when \( y \) and its derivatives are replaced throughout by \( f(x) \) and its corresponding derivatives. Also, the initial conditions must be satisfied.

For example a solution of the differential equation

\[
\frac{d^2y}{dx^2} + y = 0
\]

is

\[
y = k_1 \sin x + k_2 \cos x
\]

since \( y \) and its second derivative satisfy the given differential equation.

Any linear, time-invariant electric circuit can be described by an ODE which has the form

\[
a_n \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + \ldots + a_1 \frac{dy}{dt} + a_0 y = 0
\]

For excitation

\[
b_m \frac{d^m x}{dt^m} + b_{m-1} \frac{d^{m-1} x}{dt^{m-1}} + \ldots + b_1 \frac{dx}{dt} + b_0 x
\]

If the excitation in (B.12) is not zero, that is, if \( x(t) \neq 0 \), the ODE is called a non-homogeneous ODE. If \( x(t) = 0 \), it reduces to:

\[
a_n \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + \ldots + a_1 \frac{dy}{dt} + a_0 y = 0
\]

The differential equation of (B.13) above is called a homogeneous ODE and has \( n \) different linearly independent solutions denoted as \( y_1(t), y_2(t), y_3(t), \ldots, y_n(t) \).

We will now prove that the most general solution of (B.13) is:

\[
y_H(t) = k_1 y_1(t) + k_2 y_2(t) + k_3 y_3(t) + \ldots + k_n y_n(t)
\]

where the subscript \( H \) on the left side is used to emphasize that this is the form of the solution of the homogeneous ODE and \( k_1, k_2, k_3, \ldots, k_n \) are arbitrary constants.
Proof:

Let us assume that \( y_1(t) \) is a solution of (B.13); then by substitution,

\[
a_n \frac{d^n y_1}{dt^n} + a_{n-1} \frac{d^{n-1} y_1}{dt^{n-1}} + \ldots + a_1 \frac{dy_1}{dt} + a_0 y_1 = 0
\]

(B.15)

A solution of the form \( k_1 y_1(t) \) will also satisfy (B.13) since

\[
a_n \frac{d^n}{dt^n}(k_1 y_1) + a_{n-1} \frac{d^{n-1}}{dt^{n-1}}(k_1 y_1) + \ldots + a_1 \frac{d}{dt}(k_1 y_1) + a_0 (k_1 y_1)
\]

\[= k_1 \left( a_n \frac{d^n y_1}{dt^n} + a_{n-1} \frac{d^{n-1} y_1}{dt^{n-1}} + \ldots + a_1 \frac{dy_1}{dt} + a_0 y_1 \right) = 0 \]

(B.16)

If \( y = y_1(t) \) and \( y = y_2(t) \) are any two solutions, then \( y = y_1(t) + y_2(t) \) will also be a solution since

\[
a_n \frac{d^n y_1}{dt^n} + a_{n-1} \frac{d^{n-1} y_1}{dt^{n-1}} + \ldots + a_1 \frac{dy_1}{dt} + a_0 y_1 = 0
\]

and

\[
a_n \frac{d^n y_2}{dt^n} + a_{n-1} \frac{d^{n-1} y_2}{dt^{n-1}} + \ldots + a_1 \frac{dy_2}{dt} + a_0 y_2 = 0
\]

Therefore, (B.17)

\[
a_n \frac{d^n}{dt^n} (y_1 + y_2) + a_{n-1} \frac{d^{n-1}}{dt^{n-1}} (y_1 + y_2) + \ldots + a_1 \frac{d}{dt} (y_1 + y_2) + a_0 (y_1 + y_2)
\]

\[= a_n \frac{d^n}{dt^n} y_1 + a_{n-1} \frac{d^{n-1}}{dt^{n-1}} y_1 + \ldots + a_1 \frac{d}{dt} y_1 + a_0 y_1
\]

\[+ a_n \frac{d^n}{dt^n} y_2 + a_{n-1} \frac{d^{n-1}}{dt^{n-1}} y_2 + \ldots + a_1 \frac{d}{dt} y_2 + a_0 y_2 = 0
\]

In general, if

\[y = k_1 y_1(t), k_2 y_1(t), k_3 y_3(t), \ldots, k_n y_n(t)\]

are the \( n \) solutions of the homogeneous ODE of (B.13), the linear combination

\[y = k_1 y_1(t) + k_2 y_1(t) + k_3 y_3(t) + \ldots + k_n y_n(t)\]

is also a solution.

In our subsequent discussion, the solution of the homogeneous ODE, i.e., the complementary solution, will be referred to as the *natural response*, and will be denoted as \( y_N(t) \) or simply \( y_N \). The particular solution of a non-homogeneous ODE will be referred to as the *forced response*, and will be
denoted as \( y_F(t) \) or simply \( y_F \). Accordingly, we express the total solution of the non-homogeneous ODE of (B.12) as:

\[
y(t) = y_{\text{Natural}} + y_{\text{Forced}} = y_N + y_F
\]

(B.18)

The natural response \( y_N \) contains arbitrary constants and these can be evaluated from the given initial conditions. The forced response \( y_F \), however, contains no arbitrary constants. It is imperative to remember that the arbitrary constants of the natural response must be evaluated from the total response.

### B.4 Solution of the Homogeneous ODE

Let the solutions of the homogeneous ODE

\[
a_n \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + \ldots + a_1 \frac{dy}{dt} + a_0 y = 0
\]

(B.19)

be of the form

\[
y = ke^{st}
\]

(B.20)

Then, by substitution of (B.20) into (B.19) we get

\[
a_n kse^{st} + a_{n-1} kse^{st} + \ldots + a_1 kse^{st} + a_0 ke^{st} = 0
\]

or

\[
(a_n s^n + a_{n-1} s^{n-1} + \ldots + a_1 s + a_0) ke^{st} = 0
\]

(B.21)

We observe that (B.21) can be satisfied when

\[
(a_n s^n + a_{n-1} s^{n-1} + \ldots + a_1 s + a_0) = 0 \quad \text{or} \quad k = 0 \quad \text{or} \quad s = -\infty
\]

(B.22)

but the only meaningful solution is the quantity enclosed in parentheses since the latter two yield trivial (meaningless) solutions. We, therefore, accept the expression inside the parentheses as the only meaningful solution and this is referred to as the characteristic (auxiliary) equation, that is,

\[
\frac{(a_n s^n + a_{n-1} s^{n-1} + \ldots + a_1 s + a_0) = 0}{\text{Characteristic Equation (B.23)}}
\]

Since the characteristic equation is an algebraic equation of an \( n \)-th power polynomial, its solutions are \( s_1, s_2, s_3, \ldots, s_n \), and thus the solutions of the homogeneous ODE are:

\[
y_1 = k_1 e^{s_1 t}, \quad y_2 = k_2 e^{s_2 t}, \quad y_3 = k_3 e^{s_3 t}, \ldots, \quad y_n = k_n e^{s_n t}
\]

(B.24)
Case I – Distinct Roots

If the roots of the characteristic equation are distinct (different from each other), the $n$ solutions of (B.23) are independent and the most general solution is:

$$y_N = k_1 e^{s_1 t} + k_2 e^{s_2 t} + \ldots + k_n e^{s_n t} \quad \text{(B.25)}$$

**FOR DISTINCT ROOTS**

Case II – Repeated Roots

If two or more roots of the characteristic equation are repeated (same roots), then some of the terms of (B.24) are not independent and therefore (B.25) does not represent the most general solution. If, for example, $s_1 = s_2$, then,

$$k_1 e^{s_1 t} + k_2 e^{s_2 t} = k_1 e^{s_1 t} + k_2 e^{s_1 t} = (k_1 + k_2) e^{s_1 t} = k_3 e^{s_1 t}$$

and we see that one term of (B.25) is lost. In this case, we express one of the terms of (B.25), say $k_2 e^{s_1 t}$ as $k_2 t e^{s_1 t}$. These two represent two independent solutions and therefore the most general solution has the form:

$$y_N = (k_1 + k_2 t) e^{s_1 t} + k_3 e^{s_2 t} + \ldots + k_n e^{s_n t} \quad \text{(B.26)}$$

If there are $m$ equal roots the most general solution has the form:

$$y_N = (k_1 + k_2 t + \ldots + k_m t^{m-1}) e^{s_1 t} + k_{n-1} e^{s_2 t} + \ldots + k_n e^{s_n t} \quad \text{(B.27)}$$

**FOR M EQUAL ROOTS**

Case III – Complex Roots

If the characteristic equation contains complex roots, these occur as complex conjugate pairs. Thus, if one root is $s_1 = -\alpha + j\beta$ where $\alpha$ and $\beta$ are real numbers, then another root is $s_1 = -\alpha - j\beta$. Then,

$$k_1 e^{s_1 t} + k_2 e^{s_2 t} = k_1 e^{-\alpha t + \beta t} + k_2 e^{-\alpha t - \beta t} = e^{-\alpha t} (k_1 e^{j\beta t} + k_2 e^{-j\beta t})$$

$$= e^{-\alpha t} (k_1 \cos \beta t + j k_1 \sin \beta t + k_2 \cos \beta t - j k_2 \sin \beta t)$$

$$= e^{-\alpha t} [(k_1 + k_2) \cos \beta t + j (k_1 - k_2) \sin \beta t]$$

$$= e^{-\alpha t} (k_3 \cos \beta t + k_4 \sin \beta t) = e^{-\alpha t} k_5 \cos (\beta t + \phi) \quad \text{(B.28)}$$

**FOR TWO COMPLEX CONJUGATE ROOTS**
If (B.28) is to be a real function of time, the constants $k_1$ and $k_2$ must be complex conjugates. The other constants $k_3$, $k_4$, $k_5$, and the phase angle $\phi$ are real constants.

The forced response can be found by

a. *The Method of Undetermined Coefficients* or

b. *The Method of Variation of Parameters*

We will study the Method of Undetermined Coefficients first.

**B.5 Using the Method of Undetermined Coefficients for the Forced Response**

For simplicity, we will only consider ODEs of *order 2*. Higher order ODEs are discussed in differential equations textbooks.

Consider the non-homogeneous ODE

$$a\frac{d^2y}{dt^2} + b\frac{dy}{dt} + cy = f(x)$$

(B.29)

where $a$, $b$, and $c$ are real constants.

We have learned that the total (complete) solution consists of the summation of the natural and forced responses.

For the natural response, if $y_1$ and $y_2$ are any two solutions of (B.29), the linear combination $y_3 = k_1y_1 + k_2y_2$, where $k_1$ and $k_2$ are arbitrary constants, is also a solution, that is, if we know the two solutions, we can obtain the most general solution by forming the linear combination of $y_1$ and $y_2$. To be certain that there exist no other solutions, we examine the Wronskian Determinant defined below.

$$W(y_1, y_2) \equiv \begin{vmatrix} y_1 & y_2 \\ \frac{dy_1}{dx} & \frac{dy_2}{dx} \end{vmatrix} = y_1 \frac{d}{dx}y_2 - y_2 \frac{d}{dx}y_1 \neq 0$$

(B.30)

**WRONSKIAN DETERMINANT**

If (B.30) is true, we can be assured that all solutions of (B.29) are indeed the linear combination of $y_1$ and $y_2$.

The forced response is, in most circuit analysis problems, obtained by observation of the right side of the given ODE as it is illustrated by the examples that follow.
Example B.6

Find the total solution of the ODE

\[ \frac{d^2y}{dt^2} + 4 \frac{dy}{dt} + 3y = 0 \]  

subject to the initial conditions \( y(0) = 3 \) and \( y'(0) = 4 \) where \( y' = \frac{dy}{dt} \)

Solution:

This is a homogeneous ODE and its total solution is just the natural response found from the characteristic equation \( s^2 + 4s + 3 = 0 \) whose roots are \( s_1 = -1 \) and \( s_2 = -3 \). The total response is:

\[ y(t) = y_N(t) = k_1e^{-t} + k_2e^{-3t} \]  

(B.32)

The constants \( k_1 \) and \( k_2 \) are evaluated from the given initial conditions. For this example,

\[ y(0) = 3 = k_1e^0 + k_2e^0 \]

or

\[ k_1 + k_2 = 3 \]  

(B.33)

Also,

\[ y'(0) = 4 = \frac{dy}{dt}\bigg|_{t=0} = -k_1e^{-t} - 3k_2e^{-3t}\bigg|_{t=0} \]

or

\[ -k_1 - 3k_2 = 4 \]  

(B.34)

Simultaneous solution of (B.33) and (B.34) yields \( k_1 = 6.5 \) and \( k_2 = -3.5 \). By substitution into (B.32), we get

\[ y(t) = y_N(t) = 6.5e^{-t} - 3.5e^{-3t} \]  

(B.35)

Check with MATLAB:

```markdown
y = dsolve('D2y+4*Dy+3*y=0', 'y(0)=3', 'Dy(0)=4')
y = (-7/2*exp(-3*t)*exp(t)+13/2)/exp(t)
pretty(y)
- 7/2 exp(-3 t) exp(t) + 13/2
-------------------------------------
exp(t)
```

Circuit Analysis II with MATLAB Applications

Orchard Publications
The function \( y = f(t) \) is shown in Figure B.1 plotted with the MATLAB command `ezplot(y,[0 10]).`

![Figure B.1. Plot for the function \( y = f(t) \) of Example B.6.](image)

**Example B.7**

Find the total solution of the ODE

\[
\frac{d^2 y}{dt^2} + 4\frac{dy}{dt} + 3y = 3e^{-2t}
\]  
(B.36)

subject to the initial conditions \( y(0) = 1 \) and \( y'(0) = -1 \)

**Solution:**

The left side of (B.36) is the same as that of Example B.6. Therefore,

\[
y_N(t) = k_1 e^{-t} + k_2 e^{-3t}
\]  
(B.37)

(We must remember that the constants \( k_1 \) and \( k_2 \) must be evaluated from the total response).

To find the forced response, we assume a solution of the form

\[
y_F = Ae^{-2t}
\]  
(B.38)

We can find out whether our assumption is correct by substituting (B.38) into the given ODE of (B.36). Then,

\[
4Ae^{-2t} - 8Ae^{-2t} + 3Ae^{-2t} = 3e^{-2t}
\]  
(B.39)
Using the Method of Undetermined Coefficients for the Forced Response

from which \( A = -3 \) and the total solution is

\[
y(t) = y_N + y_F = k_1 e^{-t} + k_2 e^{-3t} - 3e^{-2t}
\]

(B.40)

The constants \( k_1 \) and \( k_2 \) are evaluated from the given initial conditions. For this example,

\[
y(0) = 1 = k_1 e^0 + k_2 e^0 - 3 e^0
\]

or

\[
k_1 + k_2 = 4 \tag{B.41}
\]

Also,

\[
y'(0) = -1 = \left. \frac{dy}{dt} \right|_{t=0} = -k_1 e^{-t} - 3k_2 e^{-3t} + 6 e^{-2t} \bigg|_{t=0}
\]

or

\[
-k_1 - 3k_2 = -7
\]

Simultaneous solution of (B.41) and (B.42) yields \( k_1 = 2.5 \) and \( k_2 = 1.5 \). By substitution into (B.40), we get

\[
y(t) = y_N + y_F = 2.5 e^{-t} + 1.5 e^{-3t} - 3e^{-2t} \tag{B.42}
\]

Check with MATLAB:

\[
y = \text{dsolve('D2y+4*Dy+3*y=3*exp(-2*t)','y(0)=1','Dy(0)=-1')}
\]

\[
y = (-3*exp(-2*t))*exp(t)+3/2*exp(-3*t)*exp(t)+5/2)/exp(t)
\]

\[
\text{pretty(y)}
\]

\[
-3 \exp(-2 \ t) \ \exp(t) + 3/2 \ \exp(-3 \ t) \ \exp(t) + 5/2
\]

\[
\text{-----------------------------------------------}
\]

\[
\exp(t)
\]

\[
\text{ezplot(y,[0 8])}
\]

The plot is shown in Figure B.2

Example B.8

Find the total solution of the ODE

\[
\frac{d^2 y}{dt^2} + 6 \frac{dy}{dt} + 9y = 0 \tag{B.43}
\]

subject to the initial conditions \( y(0) = -1 \) and \( y'(0) = 1 \)
Solution:
This is a homogeneous ODE and therefore its total solution is just the natural response found from the characteristic equation \( s^2 + 6s + 9 = 0 \) whose roots are \( s_1 = s_2 = -3 \) (repeated roots). Thus, the total response is

\[
y(t) = y_N = k_1 e^{-3t} + k_2 t e^{-3t}
\]  

(B.44)

Next, we evaluate the constants \( k_1 \) and \( k_2 \) from the given initial conditions. For this example,

\[
y(0) = -1 = k_1 e^0 + k_2 (0)e^0
\]

or

\[
k_1 = -1
\]

(B.45)

Also,

\[
y'(0) = 1 = \left. \frac{dy}{dt} \right|_{t=0} = -3k_1 e^{-3t} + k_2 e^{-3t} - 3k_2 t e^{-3t}\left|_{t=0}
\]

or

\[
-3k_1 + k_2 = 1
\]

(B.46)

From (B.45) and (B.46) we get yields \( k_1 = -1 \) and \( k_2 = -2 \). By substitution into (B.44),

\[
y(t) = -e^{-3t} - 2te^{-3t}
\]

(B.47)
Using the Method of Undetermined Coefficients for the Forced Response

Check with MATLAB:
\[
y = \text{dsolve}('D2y+6*Dy+9*y=0', 'y(0)=-1', 'Dy(0)=1')
\]
\[
y = -\exp(-3*t)-2*\exp(-3*t)*t
\]
\[
\text{ezplot}(y,[0 4])
\]
The plot is shown in Figure B.3.

![Figure B.3. Plot for the function $y = f(t)$ of Example B.8.](image)

Example B.9

Find the total solution of the ODE
\[
\frac{d^2y}{dt^2} + 5\frac{dy}{dt} + 6y = 3e^{-2t}
\]
(B.48)

Solution:

No initial conditions are given; therefore, we will express the solution in terms of the constants $k_1$ and $k_2$. By inspection, the roots of the characteristic equation of (B.48) are $s_1 = -2$ and $s_2 = -3$ and thus the natural response has the form
\[
y_N = k_1e^{-2t} + k_2e^{-3t}
\]
(B.49)

Next, we find the forced response by assuming a solution of the form
\[
y_F = Ae^{-2t}
\]
(B.50)
We can find out whether our assumption is correct by substitution of (B.50) into the given ODE of (B.48). Then,

\[ 4Ae^{-2t} - 10Ae^{-2t} + 6Ae^{-2t} = 3e^{-2t} \]  

(B.51)

but the sum of the three terms on the left side of (B.52) is zero whereas the right side can never be zero unless we let \( t \to \infty \) and this produces a meaningless result.

The problem here is that the right side of the given ODE of (B.48) has the same form as one of the terms of the natural response of (B.49), namely the term \( k_i e^{-2t} \).

To work around this problem, we assume that the forced response has the form

\[ y_F = Ate^{-2t} \]  

(B.52)

that is, we multiply (B.50) by \( t \) in order to eliminate the duplication of terms in the total response. Then, by substitution of (B.52) into (B.48) and equating like terms, we find that \( A = 3 \). Therefore, the total response is

\[ y(t) = y_N + y_F = k_1e^{-2t} + k_2e^{-3t} + 3te^{-2t} \]  

(B.53)

Check with MATLAB:

```matlab
y=dsolve('D2y+5*Dy+6*y=3*exp(-2*t)')
y =
-3*exp(-2*t)+3*t*exp(-2*t)+C1*exp(-3*t)+C2*exp(-2*t)
```

Example B.10

Find the total solution of the ODE

\[ \frac{d^2y}{dt^2} + 5\frac{dy}{dt} + 6y = 4\cos 5t \]  

(B.54)

Solution:

No initial conditions are given; therefore, we will express solution in terms of the constants \( k_1 \) and \( k_2 \). We observe that the left side of (B.54) is the same of that of Example B.9. Therefore, the natural response is the same, that is, it has the form

\[ y_N = k_1e^{-2t} + k_2e^{-3t} \]  

(B.55)

Next, to find the forced response and we assume a solution of the form

\[ y_F = A\cos 5t \]  

(B.56)
Using the Method of Undetermined Coefficients for the Forced Response

We can find out whether our assumption is correct by substitution of the assumed solution of (B.56) into the given ODE of (B.55). Then,

\[-25A \cos 5t - 25A \sin 5t + 6A \cos 5t = -19A \cos 5t - 25A \sin 5t = 4 \cos 5t\]

but this relation is invalid since by equating cosine and sine terms, we find that \(A = -4/19\) and also \(A = 0\). This inconsistency is a result of our failure to recognize that the derivatives of \(A \cos 5t\) produce new terms of the form \(B \sin 5t\) and these terms must be included in the forced response. Accordingly, we let

\[y_F = k_3 \sin 5t + k_4 \cos 5t\]  \hspace{1cm} \text{(B.57)}

and by substitution into (B.54) we get

\[-25k_3 \sin 5t - 25k_4 \cos 5t + 25k_3 \cos 5t - 25k_4 \sin 5t + 6k_3 \sin 5t + 6k_4 \cos 5t = 4 \cos 5t\]

Collecting like terms and equating sine and cosine terms, we obtain the following set of equations

\[19k_3 + 25k_4 = 0\]
\[25k_3 - 19k_4 = 4\]

(B.58)

We use MATLAB to solve (B.58)

\[
\text{format rat; } \begin{bmatrix} k3 \\ k4 \end{bmatrix} = \text{solve}(19*x+25*y, 25*x-19*y-4)
\]

\[
k3 = 50/493
\]
\[
k4 = -38/493
\]

Therefore, the total solution is

\[y(t) = y_N + y_F(t) = k_1 e^{-2t} + k_2 e^{-3t} + \frac{50}{493} \sin 5t + \frac{-38}{493} \cos 5t\]  \hspace{1cm} \text{(B.59)}

Check with MATLAB.

\[
y=\text{dsolve('D2y+5*Dy+6*y=4*cos(5*t)'); } y=\text{simple}(y)
\]

\[
y = -38/493 * \cos (5*t) + 50/493 * \sin (5*t) + \text{C1*exp}(-3*t) + \text{C2*exp}(-2*t)
\]

In most engineering problems the right side of the non-homogeneous ODE consists of elementary functions such as \(k\) (constant), \(x^n\) where \(n\) is a positive integer, \(e^{kx}\), \(\cos kx\), \(\sin kx\), and linear combinations of these. Table B.1 summarizes the forms of the forced response for a second order ODE with constant coefficients.
TABLE B.1 Form of the forced response for 2nd order differential equations

<table>
<thead>
<tr>
<th>Forced Response of the ODE $\frac{d^2y}{dt^2} + b\frac{dy}{dt} + cy = f(t)$</th>
<th>Form of Forced Response $y_F(t)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$k$ (constant)</td>
<td>$K$ (constant)</td>
</tr>
<tr>
<td>$kt^n$ ($n$ = positive integer)</td>
<td>$K_0 t^n + K_1 t^{n-1} + \ldots + K_{n-1} t + K_n$</td>
</tr>
<tr>
<td>$ke^{rt}$ ($r$ = real or complex)</td>
<td>$Ke^{rt}$</td>
</tr>
<tr>
<td>$k\cos\alpha t$ or $k\sin\alpha t$ ($\alpha$ = constant)</td>
<td>$K_1 \cos\alpha t + K_2 \sin\alpha t$</td>
</tr>
</tbody>
</table>
| $kt^n e^{rt} \cos\alpha t$ or $kt^n e^{rt} \sin\alpha t$ | $(K_0 t^n + K_1 t^{n-1} + \ldots + K_{n-1} t + K_n)e^{rt} \cos\alpha t$  
$+ (K_0 t^n + K_1 t^{n-1} + \ldots + K_{n-1} t + K_n)e^{rt} \sin\alpha t$ |

We must remember that if $f(t)$ is the sum of several terms, the most general form of the forced response $y_F(t)$ is the linear combination of these terms. Also, if a term in $y_F(t)$ is a duplicate of a term in the natural response $y_N(t)$, we must multiply $y_F(t)$ by the lowest power of $t$ that will eliminate the duplication.

Example B.11

Find the total solution of the ODE

$$\frac{d^2y}{dt^2} + 4\frac{dy}{dt} + 4y = te^{-2t} - e^{-2t} \quad (B.60)$$

Solution:

No initial conditions are given; therefore we will express solution in terms of the constants $k_1$ and $k_2$. The roots of the characteristic equation are equal, that is, $s_1 = s_2 = -2$, and thus the natural response has the form

$$y_N = k_1 e^{-2t} + k_2 te^{-2t} \quad (B.61)$$

To find the forced response (particular solution), we refer to the table of the previous page and from the last row we choose the term $kt^n e^{rt} \cos\alpha t$. This term with $n = 1$, $r = -2$, and $\alpha = 0$, reduces to $kte^{-2t}$. Therefore the forced response will have the form...
Using the Method of Undetermined Coefficients for the Forced Response

\[ y_F = (k_3 t + k_4) e^{-2t} \]  \hspace{1cm} (B.62)

But the terms \( e^{-2t} \) and \( te^{-2t} \) are also present in (B.61); therefore, we multiply (B.62) by \( t^2 \) to obtain a suitable form for the forced response which now is

\[ y_F = (k_3 t^3 + k_4 t^2) e^{-2t} \]  \hspace{1cm} (B.63)

Now, we need to evaluate the constants \( k_3 \) and \( k_4 \). This is done by substituting (B.63) into the given ODE of (B.60) and equating with the right side. We use MATLAB to do the computations as shown below.

```matlab
syms t k3 k4
f0 = (k3*t^3 + k4*t^2)*exp(-2*t); % Forced response (B.64)
f1 = diff(f0); f1 = simple(f1) % Compute and simplify first derivative
f1 = -t*exp(-2*t)*(-3*k3*t-2*k4+2*k3*t^2+2*k4*t)
f2 = diff(f0,2); f2 = simple(f2) % Compute and simplify second derivative
f2 = 2*exp(-2*t)*(3*k3*t+k4-6*k3*t^2-4*k4*t+2*k3*t^3+2*k4*t^2)
f = f2+4*f1+4*f0; f = simple(f) % Form and simplify the left side of the given ODE
f = 2*(3*k3*t+k4)*exp(-2*t)
```

Finally, we equate \( f \) above with the right side of the given ODE, that is

\[ 2(3k_3 t + k_4)e^{-2t} = te^{-2t} - e^{-2t} \]  \hspace{1cm} (B.64)

and we find \( k_3 = 1/6 \) and \( k_4 = -1/2 \). By substitution of these values into (B.64) and combining the forced response with the natural response, we get the total solution

\[ y(t) = k_1 e^{-2t} + k_2 te^{-2t} + \frac{1}{6} t^3 e^{-2t} - \frac{1}{2} t^2 e^{-2t} \]  \hspace{1cm} (B.65)

We verify this solution with MATLAB

```matlab
z = dsolve('D2y+4*Dy+4*y=t*exp(-2*t)-exp(-2*t)')
```

\[ z = \]

\[ 1/6*exp(-2*t)*t^3-1/2*exp(-2*t)*t^2+C1*exp(-2*t)+C2*t*exp(-2*t) \]
B.6 Using the Method of Variation of Parameters for the Forced Response

In certain non-homogeneous ODEs, the right side \( f(t) \) cannot be determined by the method of undetermined coefficients. For these ODEs we must use the method of variation of parameters. This method will work with all linear equations including those with variable coefficients such as

\[
\frac{d^2y}{dt^2} + \alpha(t)\frac{dy}{dt} + \beta(t)y = f(t) \tag{B.66}
\]

provided that the general form of the natural response is known.

Our discussion will be restricted to second order ODEs with constant coefficients.

The method of variation of parameters replaces the constants \( k_1 \) and \( k_2 \) by two variables \( u_1 \) and \( u_2 \) that satisfy the following three relations:

\[
y = u_1 y_1 + u_2 y_2 \tag{B.67}
\]

\[
\frac{du_1}{dt} y_1 + \frac{du_2}{dt} y_2 = 0 \tag{B.68}
\]

\[
\frac{du_1}{dt} \cdot \frac{dy_1}{dt} + \frac{du_2}{dt} \cdot \frac{dy_2}{dt} = f(t) \tag{B.69}
\]

Simultaneous solution of \( \text{(B.68)} \) and \( \text{(B.69)} \) will yield the values of \( du_1/\, dt \) and \( du_2/\, dt \); then, integration of these will produce \( u_1 \) and \( u_2 \), which when substituted into \( \text{(B.67)} \) will yield the total solution.

Example B.12

Find the total solution of

\[
\frac{d^2y}{dt^2} + 4\frac{dy}{dt} + 3y = 12 \tag{B.70}
\]

in terms of the constants \( k_1 \) and \( k_2 \) by the

a. method of undetermined coefficients

b. method of variation of parameters

Solution:

With either method, we must first find the natural response. The characteristic equation yields the roots \( s_1 = -1 \) and \( s_2 = -3 \). Therefore, the natural response is
Using the Method of Variation of Parameters for the Forced Response

\[ y_N = k_1 e^{-t} + k_2 e^{-3t} \]  \hspace{1cm} (B.71)

a. Using the method of undetermined coefficients we let \( y_F = k_3 \) (a constant). Then, by substitution into (B.70) we get \( k_3 = 4 \) and thus the total solution is

\[ y(t) = y_N + y_F = k_1 e^{-t} + k_2 e^{-3t} + 4 \]  \hspace{1cm} (B.72)

b. With the method of variation of parameters we start with the natural response found above as (B.71) and we let the solutions \( y_1 \) and \( y_2 \) be represented as

\[ y_1 = e^{-t} \text{ and } y_2 = e^{-3t} \]  \hspace{1cm} (B.73)

Then by (B.67), the total solution is

\[ y = u_1 y_1 + u_2 y_2 \]

or

\[ y = u_1 e^{-t} + u_2 e^{-3t} \]  \hspace{1cm} (B.74)

Also, from (B.68),

\[ \frac{d u_j}{d t} y_1 + \frac{d u_2}{d t} y_2 = 0 \]

or

\[ \frac{d u_1}{d t} e^{-t} + \frac{d u_2}{d t} e^{-3t} = 0 \]  \hspace{1cm} (B.75)

and from (B.69),

\[ \frac{d u_1}{d t} \cdot \frac{d y_1}{d t} + \frac{d u_2}{d t} \cdot \frac{d y_2}{d t} = f(t) \]

or

\[ \frac{d u_1}{d t} (-e^{-t}) + \frac{d u_2}{d t} (-3e^{-3t}) = 12 \]  \hspace{1cm} (B.76)

Next, we find \( du_1 / dt \) and \( du_2 / dt \) by Cramer’s rule as follows:

\[ \frac{d u_1}{d t} = \begin{vmatrix} 0 & e^{-3t} \\ 12 & -3e^{-3t} \\ e^{-t} & e^{-3t} \end{vmatrix} = \frac{-12e^{-3t}}{-2e^{-4t}} = \frac{6e^t}{1} \]  \hspace{1cm} (B.77)

and
Now, integration of (B.77) and (B.78) and substitution into (B.75) yields

\[ u_1 = 6 \int e^t \, dt = 6e^t + k_1 \quad u_2 = -6 \int e^{-3t} \, dt = -2e^{-3t} + k_2 \quad (B.79) \]

\[
y = u_1 e^{-t} + u_2 e^{-3t} = ((6e^t + k_1)e^{-t} + (-2e^{-3t} + k_2)e^{-3t})
= (6 + k_1 e^{-t} - 2 + k_2 e^{-3t})
= (k_1 e^{-t} + k_2 e^{-3t} + 4) \quad (B.80)
\]

We observe that the last expression in (B.80) is the same as (B.72) of part (a).

Check with MATLAB:

\[
y = \text{dsolve('D2y+4*Dy+3*y=12')}\]

\[
y = (4*\exp(t)+C1*\exp(-3*t)*\exp(t)+C2)/\exp(t)\]

**Example B.13**

Find the total solution of

\[
\frac{d^2y}{dt^2} + 4y = \tan 2t \quad (B.81)
\]

in terms of the constants \(k_1\) and \(k_2\) by any method.

**Solution:**

This ODE cannot be solved by the method of undetermined coefficients; therefore, we will use the method of variation of parameters.

The characteristic equation is \(s^2 + 4 = 0\) from which \(s = \pm j2\) and thus the natural response is

\[
y_N = k_1 e^{j2t} + k_2 e^{-j2t} \quad (B.82)\]

We let

\[
y_1 = \cos 2t \quad \text{and} \quad y_2 = \sin 2t \quad (B.83)\]
Using the Method of Variation of Parameters for the Forced Response

Then, by (B.67) the solution is

$$y = u_1y_1 + u_2y_2 = u_1\cos 2t + u_2\sin 2t$$  \hspace{1cm} (B.84)

Also, from (B.68),

$$\frac{du_1}{dt}y_1 + \frac{du_2}{dt}y_2 = 0$$

or

$$\frac{du_1}{dt}\cos 2t + \frac{du_2}{dt}\sin 2t = 0$$  \hspace{1cm} (B.85)

and from (B.69),

$$\frac{du_1}{dt}\cdot \frac{dy_1}{dt} + \frac{du_2}{dt}\cdot \frac{dy_2}{dt} = f(t) = \frac{du_1}{dt}(-2\sin 2t) + \frac{du_2}{dt}(2\cos 2t) = \tan 2t$$  \hspace{1cm} (B.86)

Next, we find $\frac{du_1}{dt}$ and $\frac{du_2}{dt}$ by Cramer's rule as follows:

$$\frac{du_1}{dt} = \begin{vmatrix} 0 & \sin 2t \\ \tan 2t & 2\cos 2t \end{vmatrix} = \frac{\sin 2t}{\cos 2t} = \frac{-\sin^2 2t}{2\cos 2t}$$  \hspace{1cm} (B.87)

and

$$\frac{du_2}{dt} = \begin{vmatrix} \cos 2t & 0 \\ -2\sin 2t & \tan 2t \end{vmatrix} = \frac{\sin 2t}{2}$$  \hspace{1cm} (B.88)

Now, integration of (B.87) and (B.88) and substitution into (B.84) yields

$$u_1 = -\frac{1}{2} \int \frac{\sin^2 2t}{\cos 2t} dt = \frac{\sin 2t}{4} - \frac{1}{4} \ln(\sec 2t + \tan 2t) + k_1$$  \hspace{1cm} (B.89)

$$u_2 = \frac{1}{2} \int \sin 2t dt = -\frac{\cos 2t}{4} + k_2$$  \hspace{1cm} (B.90)

$$y = u_1y_1 + u_2y_2 = \frac{\sin 2t\cos 2t}{4} - \frac{1}{4} \cos 2t \ln(\sec 2t + \tan 2t) + k_1\cos 2t - \frac{\sin 2t\cos 2t}{4} + k_2\sin 2t$$  \hspace{1cm} (B.91)

Check with MATLAB:

```matlab
y=dsolve('D2y+4*y=tan(2*t)')
```

$$y = -\frac{1}{4} \cos(2t) \times \log((1+\sin(2t)) / \cos(2t)) + C1 \times \cos(2t) + C2 \times \sin(2t)$$
B.7 Exercises

Solve the following ODEs by any method.

1. \[
\frac{d^2y}{dt^2} + 4\frac{dy}{dt} + 3y = t - 1
\]
   Answer: \( y = k_1e^{-t} + k_2e^{-3t} + \frac{1}{3}t - \frac{7}{9} \)

2. \[
\frac{d^2y}{dt^2} + 4\frac{dy}{dt} + 3y = 4e^{-t}
\]
   Answer: \( y = k_1e^{-t} + k_2e^{-3t} + 2te^{-t} \)

3. \[
\frac{d^2y}{dt^2} + 2\frac{dy}{dt} + y = \cos^2 t \quad \text{Hint: Use } \cos^2 t = \frac{1}{2}(\cos 2t + 1)
\]
   Answer: \( y = k_1e^{-t} + k_2te^{-t} + \frac{1}{2} - \frac{3\cos 2t - 4\sin 2t}{50} \)

4. \[
\frac{d^2y}{dt^2} + y = \sec t
\]
   Answer: \( y = k_1\cos t + k_2\sin t + tsint + \cos(\ln \cos t) \)
Appendix C

Matrices and Determinants

This appendix is an introduction to matrices and matrix operations. Determinants, Cramer’s rule, and Gauss’s elimination method are reviewed. Some definitions and examples are not applicable to subsequent material presented in this text, but are included for subject continuity, and reference to more advance topics in matrix theory. These are denoted with a dagger (†) and may be skipped.

C.1 Matrix Definition

A matrix is a rectangular array of numbers such as those shown below.

\[
\begin{bmatrix}
2 & 3 & 7 \\
1 & -1 & 5
\end{bmatrix}
\quad \text{or} \quad
\begin{bmatrix}
1 & 3 & 1 \\
-2 & 1 & -5 \\
4 & -7 & 6
\end{bmatrix}
\]

In general form, a matrix \( A \) is denoted as

\[
A = \begin{bmatrix}
a_{11} & a_{12} & a_{13} & \ldots & a_{1n} \\
a_{21} & a_{22} & a_{23} & \ldots & a_{2n} \\
a_{31} & a_{32} & a_{33} & \ldots & a_{3n} \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
a_{m1} & a_{m2} & a_{m3} & \ldots & a_{mn}
\end{bmatrix}
\quad (C.1)
\]

The numbers \( a_{ij} \) are the elements of the matrix where the index \( i \) indicates the row, and \( j \) indicates the column in which each element is positioned. Thus, \( a_{43} \) indicates the element positioned in the fourth row and third column.

A matrix of \( m \) rows and \( n \) columns is said to be of \( m \times n \) order matrix.

If \( m = n \), the matrix is said to be a square matrix of order \( m \) (or \( n \)). Thus, if a matrix has five rows and five columns, it is said to be a square matrix of order 5.

In a square matrix, the elements \( a_{11}, a_{22}, a_{33}, \ldots, a_{nn} \) are called the main diagonal elements. Alternatively, we say that the matrix elements \( a_{11}, a_{22}, a_{33}, \ldots, a_{nn} \), are located on the main diagonal.
† The sum of the diagonal elements of a square matrix $A$ is called the *trace* of $A$.

† A matrix in which every element is zero, is called a *zero matrix*.

### C.2 Matrix Operations

Two matrices $A = \begin{bmatrix} a_{ij} \end{bmatrix}$ and $B = \begin{bmatrix} b_{ij} \end{bmatrix}$ are equal, that is, $A = B$, if and only if

$$a_{ij} = b_{ij} \quad i = 1, 2, 3, \ldots, m \quad j = 1, 2, 3, \ldots, n$$

(C.2)

Two matrices are said to be *conformable for addition (subtraction)*, if they are of the same order $m \times n$.

If $A = \begin{bmatrix} a_{ij} \end{bmatrix}$ and $B = \begin{bmatrix} b_{ij} \end{bmatrix}$ are conformable for addition (subtraction), their sum (difference) will be another matrix $C$ with the same order as $A$ and $B$, where each element of $C$ is the sum (difference) of the corresponding elements of $A$ and $B$, that is,

$$C = A \pm B = \begin{bmatrix} a_{ij} \pm b_{ij} \end{bmatrix}$$

(C.3)

**Example C.1**

Compute $A + B$ and $A - B$ given that

$$A = \begin{bmatrix} 1 & 2 & 3 \\ 0 & 1 & 4 \end{bmatrix} \quad \text{and} \quad B = \begin{bmatrix} 2 & 3 & 0 \\ -1 & 2 & 5 \end{bmatrix}$$

**Solution:**

$$A + B = \begin{bmatrix} 1+2 & 2+3 & 3+0 \\ 0-1 & 1+2 & 4+5 \end{bmatrix} = \begin{bmatrix} 3 & 5 & 3 \\ -1 & 3 & 9 \end{bmatrix}$$

and

$$A - B = \begin{bmatrix} 1-2 & 2-3 & 3-0 \\ 0+1 & 1-2 & 4-5 \end{bmatrix} = \begin{bmatrix} -1 & -1 & 3 \\ 1 & -1 & -1 \end{bmatrix}$$

Check with MATLAB:

```matlab
A=[1 2 3; 0 1 4]; B=[2 3 0; -1 2 5]; % Define matrices A and B
A+B % Add A and B
```

* Henceforth, all paragraphs and topics preceded by a dagger ($†$) may be skipped. These are discussed in matrix theory textbooks.*
Matrix Operations

ans =
3 5 3
-1 3 9

A–B                           % Subtract B from A
ans =
-1 -1 3
1 -1 -1

If \( k \) is any scalar (a positive or negative number), and not \([k]\) which is a \(1 \times 1\) matrix, then multiplication of a matrix \(A\) by the scalar \(k\) is the multiplication of every element of \(A\) by \(k\).

Example C.2

Multiply the matrix

\[
A = \begin{bmatrix}
1 & -2 \\
2 & 3
\end{bmatrix}
\]

by

a. \( k_1 = 5 \)

b. \( k_2 = -3 + j2 \)

Solution:

a. 
\[
k_1 \cdot A = 5 \times \begin{bmatrix}
1 & -2 \\
2 & 3
\end{bmatrix} = \begin{bmatrix}
5 \times 1 & 5 \times (-2) \\
5 \times 2 & 5 \times 3
\end{bmatrix} = \begin{bmatrix}
5 & -10 \\
10 & 15
\end{bmatrix}
\]

b. 
\[
k_2 \cdot A = (-3 + j2) \times \begin{bmatrix}
1 & -2 \\
2 & 3
\end{bmatrix} = \begin{bmatrix}
(-3 + j2) \times 1 & (-3 + j2) \times (-2) \\
(-3 + j2) \times 2 & (-3 + j2) \times 3
\end{bmatrix} = \begin{bmatrix}
-3 + j2 & 6 - j4 \\
-6 + j4 & -9 + j6
\end{bmatrix}
\]

Check with MATLAB:

\[
k1=5; k2=(-3 + 2*j); \% Define scalars \(k_1\) and \(k_2\)
A=[1 -2; 2 3]; \% Define matrix \(A\)
k1*A \% Multiply matrix \(A\) by constant \(k_1\)
\]

ans =
5 -10
10 15
Two matrices $A$ and $B$ are said to be *conformable for multiplication* $A \cdot B$ in that order, only when the number of columns of matrix $A$ is equal to the number of rows of matrix $B$. That is, the product $A \cdot B$ (but not $B \cdot A$) is conformable for multiplication only if $A$ is an $m \times p$ matrix and matrix $B$ is an $p \times n$ matrix. The product $A \cdot B$ will then be an $m \times n$ matrix. A convenient way to determine if two matrices are conformable for multiplication is to write the dimensions of the two matrices side-by-side as shown below.

\[
\text{Shows that } A \text{ and } B \text{ are conformable for multiplication}
\]

\[
A \xleftarrow{m \times p} \quad B \xrightarrow{p \times n}
\]

\[
\text{Indicates the dimension of the product } A \cdot B
\]

For the product $B \cdot A$ we have:

\[
\text{Here, } B \text{ and } A \text{ are not conformable for multiplication}
\]

\[
B \xleftarrow{p \times n} \quad A \xrightarrow{m \times p}
\]

For matrix multiplication, the operation is row by column. Thus, to obtain the product $A \cdot B$, we multiply each element of a row of $A$ by the corresponding element of a column of $B$; then, we add these products.

**Example C.3**

Matrices $C$ and $D$ are defined as

\[
C = \begin{bmatrix} 2 & 3 & 4 \end{bmatrix} \quad \text{and} \quad D = \begin{bmatrix} 1 \\ -1 \\ 2 \end{bmatrix}
\]

Compute the products $C \cdot D$ and $D \cdot C$
Solution:

The dimensions of matrices $C$ and $D$ are respectively $1 \times 3$ and $3 \times 1$; therefore the product $C \cdot D$ is feasible, and will result in a $1 \times 1$, that is,

$$C \cdot D = \begin{bmatrix} 2 & 3 & 4 \\ 1 \\ 2 \end{bmatrix} \begin{bmatrix} 1 \\ -1 \\ 2 \end{bmatrix} = (2) \cdot (1) + (3) \cdot (-1) + (4) \cdot (2) = 7$$

The dimensions for $D$ and $C$ are respectively $3 \times 1$ and $1 \times 3$ and therefore, the product $D \cdot C$ is also feasible. Multiplication of these will produce a $3 \times 3$ matrix as follows:

$$D \cdot C = \begin{bmatrix} 1 \\ -1 \\ 2 \end{bmatrix} \begin{bmatrix} 2 & 3 & 4 \\ 1 \\ 2 \end{bmatrix} = \begin{bmatrix} (1) \cdot (2) & (1) \cdot (3) & (1) \cdot (4) \\ (-1) \cdot (2) & (-1) \cdot (3) & (-1) \cdot (4) \\ (2) \cdot (2) & (2) \cdot (3) & (2) \cdot (4) \end{bmatrix} = \begin{bmatrix} 2 & 3 & 4 \\ -2 & -3 & -4 \\ 4 & 6 & 8 \end{bmatrix}$$

Check with MATLAB:

```matlab
C=[2 3 4]; D=[1; -1; 2];  \% Define matrices C and D
C*D  \% Multiply C by D
ans =
    7

D*C  \% Multiply D by C
ans =
   2    3    4
  -2   -3   -4
   4    6    8
```

Division of one matrix by another is not defined. However, an equivalent operation exists, and it will become apparent later in this chapter, when we discuss the inverse of a matrix.

C.3 Special Forms of Matrices

† A square matrix is said to be upper triangular when all the elements below the diagonal are zero. The matrix $A$ of (C.4) is an upper triangular matrix.

In an upper triangular matrix, not all elements above the diagonal need to be non-zero.
A square matrix is said to be \textit{lower triangular}, when all the elements above the diagonal are zero. The matrix $B$ of (C.5) is a lower triangular matrix.

\[
B = \begin{bmatrix}
    a_{11} & 0 & 0 & \ldots & 0 \\
    0 & a_{22} & 0 & \ldots & 0 \\
    0 & 0 & \ddots & \ddots & \vdots \\
    \vdots & \ddots & \ddots & \ddots & 0 \\
    0 & \ldots & \ldots & 0 & a_{nn}
\end{bmatrix}
\]  
(C.5)

In a lower triangular matrix, not all elements below the diagonal need to be non-zero.

A square matrix is said to be \textit{diagonal}, if all elements are zero, except those in the diagonal. The matrix $C$ of (C.6) is a diagonal matrix.

\[
C = \begin{bmatrix}
    a_{11} & 0 & 0 & \ldots & 0 \\
    0 & a_{22} & 0 & \ldots & 0 \\
    0 & 0 & \ddots & \ddots & \vdots \\
    0 & \ldots & \ddots & \ddots & 0 \\
    0 & 0 & \ldots & 0 & a_{nn}
\end{bmatrix}
\]  
(C.6)

A diagonal matrix is called a \textit{scalar matrix}, if $a_{11} = a_{22} = a_{33} = \ldots = a_{nn} = k$ where $k$ is a scalar. The matrix $D$ of (C.7) is a scalar matrix with $k = 4$.

\[
D = \begin{bmatrix}
    4 & 0 & 0 & 0 \\
    0 & 4 & 0 & 0 \\
    0 & 0 & 4 & 0 \\
    0 & 0 & 0 & 4
\end{bmatrix}
\]  
(C.7)

A scalar matrix with $k = 1$, is called an \textit{identity matrix} $I$. Shown below are $2 \times 2$, $3 \times 3$, and $4 \times 4$ identity matrices.
Special Forms of Matrices

\[
\begin{bmatrix}
1 & 0 \\
0 & 1
\end{bmatrix}
\quad
\begin{bmatrix}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{bmatrix}
\quad
\begin{bmatrix}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{bmatrix}
\]

(C.8)

The MATLAB `eye(n)` function displays an \( n \times n \) identity matrix. For example,

\`
>> eye(4) % Display a 4 by 4 identity matrix
ans =
\`
\[
\begin{bmatrix}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1
\end{bmatrix}
\]

Likewise, the `eye(size(A))` function, produces an identity matrix whose size is the same as matrix \( A \). For example, let matrix \( A \) be defined as

\`
A=[1 3 1; -2 1 -5; 4 -7 6] % Define matrix A
```
\[
A =
\begin{bmatrix}
1 & 3 & 1 \\
-2 & 1 & -5 \\
4 & -7 & 6
\end{bmatrix}
\]

then,

\`
eye(size(A))
```

displays

\`
ans =
\`
\[
\begin{bmatrix}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{bmatrix}
\]

† The transpose of a matrix \( A \), denoted as \( A^T \), is the matrix that is obtained when the rows and columns of matrix \( A \) are interchanged. For example, if

\[
A = \begin{bmatrix}
1 & 2 & 3 \\
4 & 5 & 6
\end{bmatrix}
\]

then

\[
A^T = \begin{bmatrix}
1 & 4 \\
2 & 5 \\
3 & 6
\end{bmatrix}
\]

(C.9)
In MATLAB we use the apostrophe (') symbol to denote and obtain the transpose of a matrix. Thus, for the above example,

\[
A=\begin{bmatrix}1 & 2 & 3 \\ 4 & 5 & 6 \end{bmatrix} \quad \% \text{ Define matrix } A
\]

\[
A =
\begin{bmatrix}
1 & 2 & 3 \\
4 & 5 & 6
\end{bmatrix}
\]

\[A' \quad \% \text{ Display the transpose of } A\]

\[
\text{ans} =
\begin{bmatrix}
1 & 4 \\
2 & 5 \\
3 & 6
\end{bmatrix}
\]

† A **symmetric matrix** \( A \) is a matrix such that \( A^T = A \), that is, the transpose of a matrix \( A \) is the same as \( A \). An example of a symmetric matrix is shown below.

\[
A = \begin{bmatrix}1 & 2 & 3 \\ 2 & 4 & -5 \\ 3 & -5 & 6 \end{bmatrix} \quad A^T = \begin{bmatrix}1 & 2 & 3 \\ 2 & 4 & -5 \\ 3 & -5 & 6 \end{bmatrix} = A
\]

† If a matrix \( A \) has complex numbers as elements, the matrix obtained from \( A \) by replacing each element by its conjugate, is called the **conjugate of \( A \)**, and it is denoted as \( A^* \).

An example is shown below.

\[
A = \begin{bmatrix}1+j2 & j \\ 3 & 2-j3 \end{bmatrix} \quad A^* = \begin{bmatrix}1-j2 & -j \\ 3 & 2+j3 \end{bmatrix}
\]

MATLAB has two built-in functions which compute the complex conjugate of a number. The first, \texttt{conj(x)}, computes the complex conjugate of any complex number, and the second, \texttt{conj(A)}, computes the conjugate of a matrix \( A \). Using MATLAB with the matrix \( A \) defined as above, we get

\[
A = [1+2j \quad j; \quad 3 \quad 2-3j] \quad \% \text{ Define and display matrix } A
\]

\[
A =
\begin{bmatrix}
1.0000+ 2.0000i \\
3.0000
\end{bmatrix}
\begin{bmatrix}
0+ 1.0000i \\
2.0000- 3.0000i
\end{bmatrix}
\]

\[
\text{conj}_A=\text{conj}(A) \quad \% \text{ Compute and display the conjugate of } A
\]

\[
\text{conj}_A =
\begin{bmatrix}
1.0000- 2.0000i \\
0- 1.0000i
\end{bmatrix}
\]
3.0000  2.0000+ 3.0000i

† A square matrix \( A \) such that \( A^T = -A \) is called skew-symmetric. For example,

\[
A = \begin{bmatrix}
0 & 2 & -3 \\
-2 & 0 & -4 \\
3 & 4 & 0
\end{bmatrix} \quad A^T = \begin{bmatrix}
0 & -2 & 3 \\
2 & 0 & 4 \\
-3 & -4 & 0
\end{bmatrix} = -A
\]

Therefore, matrix \( A \) above is skew symmetric.

† A square matrix \( A \) such that \( A^{\ast} = A \) is called Hermitian. For example,

\[
A = \begin{bmatrix}
1 & 1-j & 2 \\
1+j & 3 & j \\
2 & -j & 0
\end{bmatrix} \quad A^T = \begin{bmatrix}
1 & 1+j & 2 \\
1-j & 3 & -j \\
2 & j & 0
\end{bmatrix} = A
\]

Therefore, matrix \( A \) above is Hermitian.

† A square matrix \( A \) such that \( A^{\ast} = -A \) is called skew-Hermitian. For example,

\[
A = \begin{bmatrix}
j & 1-j & 2 \\
-1-j & 3 & j \\
-2 & j & 0
\end{bmatrix} \quad A^T = \begin{bmatrix}
j & -1-j & -2 \\
1-j & 3 & j \\
2 & j & 0
\end{bmatrix} \quad A^{\ast} = \begin{bmatrix}
-j & -1+j & -2 \\
1+j & -3j & -j \\
2 & -j & 0
\end{bmatrix} = -A
\]

Therefore, matrix \( A \) above is skew-Hermitian.

**C.4 Determinants**

Let matrix \( A \) be defined as the square matrix

\[
A = \begin{bmatrix}
a_{11} & a_{12} & a_{13} & \ldots & a_{1n} \\
a_{21} & a_{22} & a_{23} & \ldots & a_{2n} \\
a_{31} & a_{32} & a_{33} & \ldots & a_{3n} \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
a_{n1} & a_{n2} & a_{n3} & \ldots & a_{nn}
\end{bmatrix}
\]  \hspace{1cm} (C.11)

then, the determinant of \( A \), denoted as \( \det A \), is defined as

\[
\det A = a_{11}a_{22}a_{33}\ldots a_{nn} + a_{12}a_{23}a_{34}\ldots a_{n1} + a_{13}a_{24}a_{35}\ldots a_{n2} + \ldots \\
- a_{n1}\ldots a_{22}a_{13}\ldots a_{n3}a_{14} - a_{n2}\ldots a_{23}a_{14} - a_{n3}\ldots a_{24}a_{15} - \ldots
\]  \hspace{1cm} (C.12)
Appendix C Matrices and Determinants

The determinant of a square matrix of order \( n \) is referred to as determinant of order \( n \).

Let \( A \) be a determinant of order 2, that is,

\[
A = \begin{bmatrix}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{bmatrix}
\]  \hspace{1cm} (C.13)

Then,

\[
detA = a_{11}a_{22} - a_{12}a_{21}
\]  \hspace{1cm} (C.14)

Example C.4

Matrices \( A \) and \( B \) are defined as

\[
A = \begin{bmatrix}1 & 2 \\ 3 & 4\end{bmatrix} \quad \text{and} \quad B = \begin{bmatrix}2 & -1 \\ 2 & 0\end{bmatrix}
\]

Compute \( detA \) and \( detB \).

Solution:

\[
detA = 1 \cdot 4 - 3 \cdot 2 = 4 - 6 = -2
\]
\[
detB = 2 \cdot 0 - 2 \cdot (-1) = 0 - (-2) = 2
\]

Check with MATLAB:

\begin{verbatim}
A=[1 2; 3 4]; B=[2 -1; 2 0]; % Define matrices A and B
det(A) % Compute the determinant of A
ans =
-2
det(B) % Compute the determinant of B
ans =
2
\end{verbatim}

Let \( A \) be a matrix of order 3, that is,

\[
A = \begin{bmatrix}
a_{11} & a_{12} & a_{13} \\
a_{21} & a_{22} & a_{23} \\
a_{31} & a_{32} & a_{33}
\end{bmatrix}
\]  \hspace{1cm} (C.15)

then, \( detA \) is found from
A convenient method to evaluate the determinant of order 3, is to write the first two columns to the right of the $3 \times 3$ matrix, and add the products formed by the diagonals from upper left to lower right; then subtract the products formed by the diagonals from lower left to upper right as shown on the diagram of the next page. When this is done properly, we obtain (C.16) above.

\[ \begin{align*}
&\begin{pmatrix}
  a_{11} & a_{12} & a_{13} \\
  a_{21} & a_{22} & a_{23} \\
  a_{31} & a_{32} & a_{33}
\end{pmatrix} \quad \begin{pmatrix}
  \text{+} \\
  \text{-}
\end{pmatrix} \\
\end{align*}\]

This method works only with second and third order determinants. To evaluate higher order determinants, we must first compute the cofactors; these will be defined shortly.

**Example C.5**

Compute \(detA\) and \(detB\) if matrices \(A\) and \(B\) are defined as

\[
A = \begin{bmatrix}
2 & 3 & 5 \\
1 & 0 & 1 \\
2 & 1 & 0
\end{bmatrix}
\quad \text{and} \quad
B = \begin{bmatrix}
2 & -3 & -4 \\
1 & 0 & -2 \\
0 & -5 & -6
\end{bmatrix}
\]

**Solution:**

\[detA = \begin{pmatrix}
2 & 3 & 5 \\
1 & 0 & 1 \\
2 & 1 & 0
\end{pmatrix} \quad \text{and} \quad
B = \begin{pmatrix}
2 & -3 & -4 \\
1 & 0 & -2 \\
0 & -5 & -6
\end{pmatrix}
\]

or

\[detA = (2 \times 0 \times 0) + (3 \times 1 \times 1) + (5 \times 1 \times 1) - (2 \times 0 \times 5) - (1 \times 1 \times 2) - (0 \times 1 \times 3) = 11 - 2 = 9\]

Likewise,

\[detB = \begin{pmatrix}
2 & -3 & -4 \\
1 & 0 & -2 \\
0 & -5 & -6
\end{pmatrix}
\]

or

\[detB = [2 \times 0 \times (-6)] + [(-3) \times (-2) \times 0] + [(-4) \times 1 \times (-5)] - [0 \times 0 \times (-4)] - [(-5) \times (-2) \times 2] - [(-6) \times 1 \times (-3)] = 20 - 38 = -18\]

Check with MATLAB:
A = [2 3 5; 1 0 1; 2 1 0]; det(A) % Define matrix A and compute detA
ans = 9

B = [2 -3 -4; 1 0 -2; 0 -5 -6]; det(B) % Define matrix B and compute detB
ans = -18

C.5 Minors and Cofactors

Let matrix $A$ be defined as the square matrix of order $n$ as shown below.

$$A = \begin{bmatrix}
  a_{11} & a_{12} & a_{13} & \ldots & a_{1n} \\
  a_{21} & a_{22} & a_{23} & \ldots & a_{2n} \\
  a_{31} & a_{32} & a_{33} & \ldots & a_{3n} \\
  \vdots & \vdots & \vdots & \ddots & \vdots \\
  a_{n1} & a_{n2} & a_{n3} & \ldots & a_{nn}
\end{bmatrix}$$  \hspace{1cm} (C.17)

If we remove the elements of its $i$th row, and $j$th column, the remaining $n-1$ square matrix is called the *minor of $A$*, and it is denoted as $[M_{ij}]$.

The signed minor $(-1)^{i+j}[M_{ij}]$ is called the *cofactor of $a_{ij}$* and it is denoted as $\alpha_{ij}$.

**Example C.6**

Matrix $A$ is defined as

$$A = \begin{bmatrix}
  a_{11} & a_{12} & a_{13} \\
  a_{21} & a_{22} & a_{23} \\
  a_{31} & a_{32} & a_{33}
\end{bmatrix}$$  \hspace{1cm} (C.18)

Compute the minors $[M_{11}]$, $[M_{12}]$, $[M_{13}]$ and the cofactors $\alpha_{11}$, $\alpha_{12}$ and $\alpha_{13}$.

**Solution:**
Minors and Cofactors

\[ M_{11} = \begin{bmatrix} a_{22} & a_{23} \\ a_{32} & a_{33} \end{bmatrix}, \quad M_{12} = \begin{bmatrix} a_{21} & a_{23} \\ a_{31} & a_{33} \end{bmatrix}, \quad M_{13} = \begin{bmatrix} a_{21} & a_{22} \\ a_{31} & a_{32} \end{bmatrix} \]

and

\[ \alpha_{11} = (-1)^{1+1} M_{11}, \quad \alpha_{12} = (-1)^{1+2} M_{12} = -M_{12}, \quad \alpha_{13} = (-1)^{1+3} M_{13} \]

The remaining minors

\[ M_{21}, \quad M_{22}, \quad M_{23}, \quad M_{31}, \quad M_{32}, \quad M_{33} \]

and cofactors

\[ \alpha_{21}, \quad \alpha_{22}, \quad \alpha_{23}, \quad \alpha_{31}, \quad \alpha_{32}, \quad \alpha_{33} \]

are defined similarly.

**Example C.7**

Compute the cofactors of matrix \( A \) defined as

\[ A = \begin{bmatrix} 1 & 2 & -3 \\ 2 & -4 & 2 \\ -1 & 2 & -6 \end{bmatrix} \]  

**Solution:**

\[ \alpha_{11} = (-1)^{1+1} \begin{bmatrix} -4 & 2 \\ 2 & -6 \end{bmatrix} = 20 \quad \alpha_{12} = (-1)^{1+2} \begin{bmatrix} 2 & 2 \\ -1 & -6 \end{bmatrix} = 10 \]  

\[ \alpha_{13} = (-1)^{1+3} \begin{bmatrix} 2 & -4 \\ -1 & 2 \end{bmatrix} = 0 \quad \alpha_{21} = (-1)^{2+1} \begin{bmatrix} 2 & -3 \\ 2 & -6 \end{bmatrix} = 6 \]

\[ \alpha_{22} = (-1)^{2+2} \begin{bmatrix} 1 & -3 \\ -1 & -6 \end{bmatrix} = -9 \quad \alpha_{23} = (-1)^{2+3} \begin{bmatrix} 1 & 2 \\ -1 & 2 \end{bmatrix} = -4 \]

\[ \alpha_{31} = (-1)^{3+1} \begin{bmatrix} 2 & -3 \\ -4 & 2 \end{bmatrix} = -8, \quad \alpha_{32} = (-1)^{3+2} \begin{bmatrix} 1 & -3 \\ 2 & 2 \end{bmatrix} = -8 \]

\[ \alpha_{33} = (-1)^{3+3} \begin{bmatrix} 1 & 2 \\ 2 & -4 \end{bmatrix} = -8 \]

It is useful to remember that the signs of the cofactors follow the pattern
that is, the cofactors on the diagonals have the same sign as their minors.

Let $A$ be a square matrix of any size; the value of the determinant of $A$ is the sum of the products obtained by multiplying each element of *any* row or *any* column by its cofactor.

**Example C.8**

Matrix $A$ is defined as

$$A = \begin{bmatrix} 1 & 2 & -3 \\ 2 & -4 & 2 \\ -1 & 2 & -6 \end{bmatrix}$$  \hspace{1cm} (C.25)

Compute the determinant of $A$ using the elements of the first row.

**Solution:**

$$\det A = 1 \left[ \begin{array}{cc} -4 & 2 \\ 2 & -6 \end{array} \right] - 2 \left[ \begin{array}{cc} 2 & 2 \\ -1 & -6 \end{array} \right] - 3 \left[ \begin{array}{cc} 2 & -4 \\ -1 & 2 \end{array} \right] = 1 \times 20 - 2 \times (-10) - 3 \times 0 = 40$$

Check with MATLAB:

```matlab
A = [1 2 -3; 2 -4 2; -1 2 -6]; det(A) \hspace{1cm} \% Define matrix A and compute \det A
ans =
   40
```

We must use the above procedure to find the determinant of a matrix $A$ of order 4 or higher. Thus, a fourth-order determinant can first be expressed as the sum of the products of the elements of its first row by its cofactor as shown below.
Minors and Cofactors

Determinants of order five or higher can be evaluated similarly.

**Example C.9**

Compute the value of the determinant of the matrix $A$ defined as

$$
A = \begin{bmatrix}
  a_{11} & a_{12} & a_{13} & a_{14} \\
  a_{21} & a_{22} & a_{23} & a_{24} \\
  a_{31} & a_{32} & a_{33} & a_{34} \\
  a_{41} & a_{42} & a_{43} & a_{44}
\end{bmatrix}
$$

Solution:

Using the above procedure, we will multiply each element of the first column by its cofactor. Then,

$$
A = 2 \begin{bmatrix}
  1 & 0 & -1 \\
  0 & 3 & -2 \\
  0 & 0 & 1
\end{bmatrix} - (-1) \begin{bmatrix}
  -1 & 0 & -3 \\
  0 & 3 & -2 \\
  0 & 0 & 1
\end{bmatrix} + 4 \begin{bmatrix}
  -1 & 0 & -3 \\
  1 & 0 & -1 \\
  0 & 0 & 1
\end{bmatrix} - (-3) \begin{bmatrix}
  -1 & 0 & -3 \\
  1 & 0 & -1 \\
  0 & 3 & -2
\end{bmatrix}
$$

Next, using the procedure of Example C.5 or Example C.8, we find

$$
[a] = 6, \quad [b] = -3, \quad [c] = 0, \quad [d] = -36
$$

and thus

$$
detA = [a] + [b] + [c] + [d] = 6 - 3 + 0 - 36 = -33
$$

We can verify our answer with MATLAB as follows:

```matlab
A=[2 -1 0 -3; -1 1 0 -1; 4 0 3 -2; -3 0 0 1]; delta = det(A)
delta = 
-33
```
Appendix C  Matrices and Determinants

Some useful properties of determinants are given below.

**Property 1:** *If all elements of one row or one column are zero, the determinant is zero.* An example of this is the determinant of the cofactor \([c]\) above.

**Property 2:** *If all the elements of one row or column are \(m\) times the corresponding elements of another row or column, the determinant is zero.* For example, if

\[
A = \begin{bmatrix}
2 & 4 & 1 \\
3 & 6 & 1 \\
1 & 2 & 1
\end{bmatrix}
\]  

(C.28)

then,

\[
detA = \begin{vmatrix}
2 & 4 & 1 \\
3 & 6 & 1 \\
1 & 2 & 1
\end{vmatrix} = 2 \cdot 4 + 6 \cdot 1 - 4 \cdot 2 = 12 + 4 + 6 - 4 - 12 = 0
\]  

(C.29)

Here, \(detA\) is zero because the second column in \(A\) is 2 times the first column.

Check with MATLAB:

\[
A=[2 \ 4 \ 1; \ 3 \ 6 \ 1; \ 1 \ 2 \ 1]; det(A)
\]

\[
\text{ans } = 0
\]

**Property 3:** *If two rows or two columns of a matrix are identical, the determinant is zero.* This follows from Property 2 with \(m = 1\).

C.6  Cramer’s Rule

Let us consider the systems of the three equations below

\[
a_{11}x + a_{12}y + a_{13}z = A
\]

\[
a_{21}x + a_{22}y + a_{23}z = B
\]

\[
a_{31}x + a_{32}y + a_{33}z = C
\]

(C.30)

and let

\[
\Delta = \begin{vmatrix}
a_{11} & a_{12} & a_{13} \\
a_{21} & a_{22} & a_{23} \\
a_{31} & a_{32} & a_{33}
\end{vmatrix}
\]

\[
D_1 = \begin{vmatrix}
a_{11} & a_{13} \\
B & a_{23} \\
C & a_{33}
\end{vmatrix}
\]

\[
D_2 = \begin{vmatrix}
a_{11} & A & a_{13} \\
B & a_{23} \\
C & a_{33}
\end{vmatrix}
\]

\[
D_3 = \begin{vmatrix}
a_{11} & a_{12} & a \n
C & a_{32} & C
\end{vmatrix}
\]
Cramer’s Rule

Cramer’s rule states that the unknowns \( x, y, \) and \( z \) can be found from the relations

\[
x = \frac{D_1}{\Delta} \quad y = \frac{D_2}{\Delta} \quad z = \frac{D_3}{\Delta}
\]

(C.31)

provided that the determinant \( \Delta \) (delta) is not zero.

We observe that the numerators of (C.31) are determinants that are formed from \( \Delta \) by the substitution of the known values \( A, B, \) and \( C, \) for the coefficients of the desired unknown.

Cramer’s rule applies to systems of two or more equations.

If (C.30) is a homogeneous set of equations, that is, if \( A = B = C = 0, \) then, \( D_1, D_2, \) and \( D_3 \) are all zero as we found in Property 1 above. Then, \( x = y = z = 0 \) also.

Example C.10

Use Cramer’s rule to find \( v_1, v_2, \) and \( v_3 \) if

\[
\begin{align*}
2v_1 - 5 - v_2 + 3v_3 &= 0 \\
-2v_3 - 3v_2 - 4v_1 &= 8 \\
v_2 + 3v_1 - 4 - v_3 &= 0
\end{align*}
\]

(C.32)

and verify your answers with MATLAB.

Solution:

Rearranging the unknowns \( v, \) and transferring known values to the right side, we get

\[
\begin{align*}
2v_1 - v_2 + 3v_3 &= 5 \\
-4v_1 - 3v_2 - 2v_3 &= 8 \\
3v_1 + v_2 - v_3 &= 4
\end{align*}
\]

(C.33)

Now, by Cramer’s rule,

\[
\Delta = \begin{vmatrix}
2 & -1 & 3 \\
-4 & -3 & -2 \\
3 & 1 & -1
\end{vmatrix}
\]

\[
D_1 = \begin{vmatrix}
5 & -1 & 3 \\
8 & -3 & -2 \\
4 & 1 & -1
\end{vmatrix}
\]

\[
\Delta = 2 - (-1) - 3 \\
\Delta = 6 + 6 - 12 + 27 + 4 + 4 = 35
\]

\[
D_1 = 5 - (-1) - 3 \\
D_1 = 15 + 8 + 24 + 36 + 10 - 8 = 85
\]
Then, using (C.31) we get

\[ D_2 = \begin{vmatrix} 2 & 5 & 3 \\ -4 & 8 & -2 \\ 3 & 4 & -1 \end{vmatrix} = 2 \begin{vmatrix} 8 & -2 \\ 4 & -1 \end{vmatrix} - 5 \begin{vmatrix} -4 & -2 \\ 3 & -1 \end{vmatrix} + 3 \begin{vmatrix} -4 & 8 \\ 3 & 4 \end{vmatrix} = -16 - 30 - 48 - 72 + 16 - 20 = -170 \]

\[ D_3 = \begin{vmatrix} 2 & -1 & 5 \\ -4 & -3 & 8 \\ 3 & 1 & 4 \end{vmatrix} = 2 \begin{vmatrix} -3 & 8 \\ 1 & 4 \end{vmatrix} - 1 \begin{vmatrix} -4 & 8 \\ 3 & 4 \end{vmatrix} + 5 \begin{vmatrix} -4 & -3 \\ 3 & 1 \end{vmatrix} = -24 - 24 - 20 + 45 - 16 - 16 = -55 \]

Then, using (C.31) we get

\[ x_1 = \frac{D_1}{\Delta} = \frac{85}{35} = \frac{17}{7} \quad x_2 = \frac{D_2}{\Delta} = \frac{-170}{35} = \frac{-34}{7} \quad x_3 = \frac{D_3}{\Delta} = \frac{55}{35} = \frac{-11}{7} \quad (C.34) \]

We will verify with MATLAB as follows.

% The following code will compute and display the values of v_1, v_2 and v_3.

\begin{verbatim}
format rat % Express answers in ratio form
B=[2 -1 3; -4 -3 -2; 3 1 -1]; % The elements of the determinant D of matrix B
delta=det(B); % Compute the determinant D of matrix B
d1=[5 -1 3; 8 -3 -2; 4 1 -1]; % The elements of D_1
detd1=det(d1); % Compute the determinant of D_1
d2=[2 5 3; -4 8 -2; 3 4 -1]; % The elements of D_2
detd2=det(d2); % Compute the determinant of D_2
d3=[2 -1 5; -4 -3 8; 3 1 4]; % The elements of D_3
detd3=det(d3); % Compute the determinant of D_3
v1=detd1/delta; % Compute the value of v_1
v2=detd2/delta; % Compute the value of v_2
v3=detd3/delta; % Compute the value of v_3

disp('v1=');disp(v1); % Display the value of v_1
disp('v2=');disp(v2); % Display the value of v_2
disp('v3=');disp(v3); % Display the value of v_3

v1 = 17/7
v2 = -34/7
v3 = -11/7
\end{verbatim}

These are the same values as in (C.34)
C.7 Gaussian Elimination Method

We can find the unknowns in a system of two or more equations also by the Gaussian elimination method. With this method, the objective is to eliminate one unknown at a time. This can be done by multiplying the terms of any of the equations of the system by a number such that we can add (or subtract) this equation to another equation in the system so that one of the unknowns will be eliminated. Then, by substitution to another equation with two unknowns, we can find the second unknown. Subsequently, substitution of the two values found can be made into an equation with three unknowns from which we can find the value of the third unknown. This procedure is repeated until all unknowns are found. This method is best illustrated with the following example which consists of the same equations as the previous example.

**Example C.11**

Use the Gaussian elimination method to find $v_1$, $v_2$, and $v_3$ of the system of equations

\[
\begin{align*}
2v_1 - v_2 + 3v_3 &= 5 \\
-4v_1 - 3v_2 - 2v_3 &= 8 \\
3v_1 + v_2 - v_3 &= 4
\end{align*}
\]  

(C.35)

**Solution:**

As a first step, we add the first equation of (C.35) with the third to eliminate the unknown $v_2$ and we obtain the following equation.

\[
5v_1 + 2v_3 = 9
\]  

(C.36)

Next, we multiply the third equation of (C.35) by 3, and we add it with the second to eliminate $v_2$. Then, we obtain the following equation.

\[
5v_1 - 5v_3 = 20
\]  

(C.37)

Subtraction of (C.37) from (C.36) yields

\[
7v_3 = -11 \text{ or } v_3 = \frac{-11}{7}
\]  

(C.38)

Now, we can find the unknown $v_1$ from either (C.36) or (C.37). By substitution of (C.38) into (C.36) we get

\[
5v_1 + 2 \cdot \left(\frac{-11}{7}\right) = 9 \text{ or } v_1 = \frac{17}{7}
\]  

(C.39)

Finally, we can find the last unknown $v_2$ from any of the three equations of (C.35). By substitution into the first equation we get
These are the same values as those we found in Example C.10.

The Gaussian elimination method works well if the coefficients of the unknowns are small integers, as in Example C.11. However, it becomes impractical if the coefficients are large or fractional numbers.

### C.8 The Adjoint of a Matrix

Let us assume that $A$ is an $n$ square matrix and $\alpha_{ij}$ is the cofactor of $a_{ij}$. Then the adjoint of $A$, denoted as $\text{adj}A$, is defined as the $n$ square matrix below.

$$\text{adj}A = \begin{bmatrix} \alpha_{11} & \alpha_{21} & \alpha_{31} & \cdots & \alpha_{n1} \\ \alpha_{12} & \alpha_{22} & \alpha_{32} & \cdots & \alpha_{n2} \\ \alpha_{13} & \alpha_{23} & \alpha_{33} & \cdots & \alpha_{n3} \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ \alpha_{1n} & \alpha_{2n} & \alpha_{3n} & \cdots & \alpha_{nn} \end{bmatrix}$$

(C.41)

We observe that the cofactors of the elements of the $i$th row (column) of $A$ are the elements of the $i$th column (row) of $\text{adj}A$.

#### Example C.12

Compute $\text{adj}A$ if Matrix $A$ is defined as

$$A = \begin{bmatrix} 1 & 2 & 3 \\ 1 & 3 & 4 \\ 1 & 4 & 3 \end{bmatrix}$$

(C.42)

Solution:

$$\text{adj}A = \begin{bmatrix} 3 & 4 \\ 4 & 3 \end{bmatrix} - \begin{bmatrix} 2 & 3 \\ 4 & 3 \end{bmatrix} - \begin{bmatrix} 2 & 3 \\ 3 & 4 \end{bmatrix} = \begin{bmatrix} -7 & 6 & -1 \\ 1 & 0 & -1 \\ 1 & -2 & 1 \end{bmatrix}$$
C.9 Singular and Non-Singular Matrices

An $n$ square matrix $A$ is called singular if $\det A = 0$; if $\det A \neq 0$, $A$ is called non-singular.

Example C.13

Matrix $A$ is defined as

$$A = \begin{bmatrix} 1 & 2 & 3 \\ 2 & 3 & 4 \\ 3 & 5 & 7 \end{bmatrix} \quad (C.43)$$

Determine whether this matrix is singular or non-singular.

Solution:

$$\det A = \begin{vmatrix} 1 & 2 & 3 \\ 2 & 3 & 4 \\ 3 & 5 & 7 \end{vmatrix} = 1 \begin{vmatrix} 3 & 4 \\ 5 & 7 \end{vmatrix} - 2 \begin{vmatrix} 2 & 4 \\ 3 & 7 \end{vmatrix} + 3 \begin{vmatrix} 2 & 3 \\ 3 & 5 \end{vmatrix} = 21 + 24 + 30 - 27 - 20 - 28 = 0$$

Therefore, matrix $A$ is singular.

C.10 The Inverse of a Matrix

If $A$ and $B$ are $n$ square matrices such that $AB = BA = I$, where $I$ is the identity matrix, $B$ is called the inverse of $A$, denoted as $B = A^{-1}$, and likewise, $A$ is called the inverse of $B$, that is, $A = B^{-1}$.

If a matrix $A$ is non-singular, we can compute its inverse $A^{-1}$ from the relation

$$A^{-1} = \frac{1}{\det A} \text{adj} A \quad (C.44)$$

Example C.14

Matrix $A$ is defined as

$$A = \begin{bmatrix} 1 & 2 & 3 \\ 1 & 3 & 4 \\ 1 & 4 & 3 \end{bmatrix} \quad (C.45)$$

Compute its inverse, that is, find $A^{-1}$.
Appendix C  Matrices and Determinants

Solution:

Here, \( \det A = 9 + 8 + 12 - 9 - 16 - 6 = -2 \), and since this is a non-zero value, it is possible to compute the inverse of \( A \) using (C.44).

From Example C.12,

\[
\text{adj} A = \begin{bmatrix}
-7 & 6 & -1 \\
1 & 0 & -1 \\
1 & -2 & 1
\end{bmatrix}
\]

Then,

\[
A^{-1} = \frac{1}{\det A} \text{adj} A = \frac{1}{-2} \begin{bmatrix}
-7 & 6 & -1 \\
1 & 0 & -1 \\
1 & -2 & 1
\end{bmatrix} = \begin{bmatrix}
3.5 & -3.0000 & 0.5000 \\
-0.5000 & 0.0000 & 0.5000 \\
-0.5000 & 1.0000 & -0.5000
\end{bmatrix}
\]

(C.46)

Check with MATLAB:

\[
\begin{align*}
A &= \begin{bmatrix} 1 & 2 & 3 \\ 1 & 3 & 4 \\ 1 & 4 & 3 \end{bmatrix}, \quad \text{inv} A = \text{inv}(A) \quad \% \text{Define matrix } A \text{ and compute its inverse} \\
A &= \\
\begin{bmatrix}
1 & 2 & 3 \\
1 & 3 & 4 \\
1 & 4 & 3
\end{bmatrix} \\
\text{inv} A &= \\
\begin{bmatrix}
3.5000 & -3.0000 & 0.5000 \\
-0.5000 & 0.0000 & 0.5000 \\
-0.5000 & 1.0000 & -0.5000
\end{bmatrix}
\end{align*}
\]

Multiplication of a matrix \( A \) by its inverse \( A^{-1} \) produces the identity matrix \( I \), that is,

\[
\begin{align*}
AA^{-1} &= I & \text{or} & & A^{-1}A &= I
\end{align*}
\]

(C.47)

Example C.15

Prove the validity of (C.47) for the Matrix \( A \) defined as

\[
A = \begin{bmatrix}
4 & 3 \\
2 & 2
\end{bmatrix}
\]

Proof:

\[
\begin{align*}
\det A &= 8 - 6 = 2 \quad \text{and} \quad \text{adj} A = \begin{bmatrix}
2 & -3 \\
-2 & 4
\end{bmatrix}
\end{align*}
\]
Then,

\[ A^{-1} = \frac{1}{\text{det}A} \text{adj}A = \frac{1}{2} \begin{bmatrix} 2 & -3 \\ -2 & 4 \end{bmatrix} = \begin{bmatrix} 1 & -3/2 \\ -1 & 2 \end{bmatrix} \]

and

\[ AA^{-1} = \begin{bmatrix} 4 & 3 \\ 2 & 2 \end{bmatrix} \begin{bmatrix} 1 & -3/2 \\ -1 & 2 \end{bmatrix} = \begin{bmatrix} 4-3 & -6+6 \\ 2-2 & -3+4 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = I \]

### C.11 Solution of Simultaneous Equations with Matrices

Consider the relation

\[ AX = B \]  \hspace{1cm} (C.48)

where \( A \) and \( B \) are matrices whose elements are known, and \( X \) is a matrix (a column vector) whose elements are the unknowns. We assume that \( A \) and \( X \) are conformable for multiplication. Multiplication of both sides of (C.48) by \( A^{-1} \) yields:

\[ A^{-1}AX = A^{-1}B = IX = A^{-1}B \]  \hspace{1cm} (C.49)

or

\[ X=A^{-1}B \]  \hspace{1cm} (C.50)

Therefore, we can use (C.50) to solve any set of simultaneous equations that have solutions. We will refer to this method as the inverse matrix method of solution of simultaneous equations.

#### Example C.16

For the system of the equations

\[
\begin{align*}
2x_1 + 3x_2 + x_3 &= 9 \\
x_1 + 2x_2 + 3x_3 &= 6 \\
3x_1 + x_2 + 2x_3 &= 8
\end{align*}
\]  \hspace{1cm} (C.51)

compute the unknowns \( x_1, x_2, \) and \( x_3 \) using the inverse matrix method.

**Solution:**

In matrix form, the given set of equations is \( AX = B \) where
Then,

\[ X = A^{-1}B \]  \hspace{1cm} \text{(C.53)}

or

\[
\begin{bmatrix}
    x_1 \\
    x_2 \\
    x_3
\end{bmatrix}
= 
\begin{bmatrix}
    2 & 3 & 1 \\
    1 & 2 & 3 \\
    3 & 1 & 2
\end{bmatrix}^{-1}
\begin{bmatrix}
    9 \\
    6 \\
    8
\end{bmatrix}
\] \hspace{1cm} \text{(C.54)}

Next, we find the determinant \( detA \), and the adjoint \( adjA \)

\[ detA = 18 \quad \text{and} \quad adjA = \begin{bmatrix}
    1 & -5 & 7 \\
    7 & 1 & -5 \\
    -5 & 7 & 1
\end{bmatrix} \]

Therefore,

\[ A^{-1} = \frac{1}{detA} adjA = \frac{1}{18} \begin{bmatrix}
    1 & -5 & 7 \\
    7 & 1 & -5 \\
    -5 & 7 & 1
\end{bmatrix} \]

and by (C.53) we obtain the solution as follows.

\[
\begin{bmatrix}
    x_1 \\
    x_2 \\
    x_3
\end{bmatrix}
= \frac{1}{18} \begin{bmatrix}
    1 & -5 & 7 \\
    7 & 1 & -5 \\
    -5 & 7 & 1
\end{bmatrix}
\begin{bmatrix}
    9 \\
    6 \\
    8
\end{bmatrix}
= \frac{1}{18} \begin{bmatrix}
    35 \\
    29 \\
    5
\end{bmatrix}
= \begin{bmatrix}
    1.94 \\
    1.61 \\
    0.28
\end{bmatrix} \] \hspace{1cm} \text{(C.55)}

To verify our results, we could use the MATLAB’s \texttt{inv(A)} function, and then multiply \( A^{-1} \) by \( B \). However, it is easier to use the \textit{matrix left division} operation \( X = A \backslash B \); this is MATLAB’s solution of \( A^{-1}B \) for the matrix equation \( A \cdot X = B \), where matrix \( X \) is the same size as matrix \( B \). For this example,

\[
A = \begin{bmatrix}
    2 & 3 & 1 \\
    1 & 2 & 3 \\
    3 & 1 & 2
\end{bmatrix}; \quad B = \begin{bmatrix}
    9 \\
    6 \\
    8
\end{bmatrix};
\]

\[
X = A \backslash B
\]
Example C.17

For the electric circuit of Figure C.1,

\begin{align*}
V = 100 \text{ v}
\end{align*}

the loop equations are

\begin{align*}
10I_1 - 9I_2 &= 100 \\
-9I_1 + 20I_2 - 9I_3 &= 0 \\
-9I_2 + 15I_3 &= 0
\end{align*}  \tag{C.56}

Use the inverse matrix method to compute the values of the currents $I_1$, $I_2$, and $I_3$.

Solution:

For this example, the matrix equation is $RI = V$ or $I = R^{-1}V$, where

\begin{align*}
R &= \begin{bmatrix}
10 & -9 & 0 \\
-9 & 20 & -9 \\
0 & -9 & 15
\end{bmatrix}, \quad V = \begin{bmatrix}
100 \\
0 \\
0
\end{bmatrix} \quad \text{and} \quad I = \begin{bmatrix}
I_1 \\
I_2 \\
I_3
\end{bmatrix}
\end{align*}

The next step is to find $R^{-1}$. This is found from the relation

\begin{align*}
R^{-1} = \frac{1}{\text{det}R} \text{adj}R
\end{align*}  \tag{C.57}

Therefore, we find the determinant and the adjoint of $R$. For this example, we find that

\begin{align*}
I = \begin{bmatrix}
1.9444 \\
1.6111 \\
0.2778
\end{bmatrix}
\end{align*}
Appendix C  Matrices and Determinants

\[ detR = 975, \quad \text{adj}R = \begin{bmatrix} 219 & 135 & 81 \\ 135 & 150 & 90 \\ 81 & 90 & 119 \end{bmatrix} \]  \hspace{1cm} (C.58)

Then,

\[ R^{-1} = \frac{1}{detR} \text{adj}R = \frac{1}{975} \begin{bmatrix} 219 & 135 & 81 \\ 135 & 150 & 90 \\ 81 & 90 & 119 \end{bmatrix} \]

and

\[ I = \begin{bmatrix} I_1 \\ I_2 \\ I_3 \end{bmatrix} = \frac{1}{975} \begin{bmatrix} 219 & 135 & 81 \\ 135 & 150 & 90 \\ 81 & 90 & 119 \end{bmatrix} \begin{bmatrix} 100 \\ 0 \\ 0 \end{bmatrix} = \frac{100}{975} \begin{bmatrix} 219 \\ 135 \\ 81 \end{bmatrix} = \begin{bmatrix} 22.46 \\ 13.85 \\ 8.31 \end{bmatrix} \]

Check with MATLAB:

\[ R=\begin{bmatrix} 10 & -9 & 0 \\ -9 & 20 & -9 \\ 0 & -9 & 15 \end{bmatrix}; \quad V=\begin{bmatrix} 100 \\ 0 \\ 0 \end{bmatrix}; \quad I=R\backslash V \]

\[ I = \begin{bmatrix} 22.4615 \\ 13.8462 \\ 8.3077 \end{bmatrix} \]

We can also use subscripts to address the individual elements of the matrix. Accordingly, the above code could also have been written as:

\[ R(1,1)=10; \quad R(1,2)=-9; \quad \% \text{No need to make entry for } A(1,3) \text{ since it is zero.} \]

\[ R(2,1)=-9; \quad R(2,2)=20; \quad R(2,3)=-9; \quad R(3,2)=-9; \quad R(3,3)=15; \quad V=\begin{bmatrix} 100 \\ 0 \\ 0 \end{bmatrix}; \quad I=R\backslash V \]

\[ I = \begin{bmatrix} 22.4615 \\ 13.8462 \\ 8.3077 \end{bmatrix} \]

Spreadsheets also have the capability of solving simultaneous equations using the inverse matrix method. For instance, we can use Microsoft Excel’s MINVERSE (Matrix Inversion) and MMULT (Matrix Multiplication) functions, to obtain the values of the three currents in Example C.17.

The procedure is as follows:

1. We start with a blank spreadsheet and in a block of cells, say B3:D5, we enter the elements of matrix \( R \) as shown in Figure C.2. Then, we enter the elements of matrix \( V \) in G3:G5.
2. Next, we compute and display the inverse of $R$, that is, $R^{-1}$. We choose B7:D9 for the elements of this inverted matrix. We format this block for number display with three decimal places. With this range highlighted and making sure that the cell marker is in B7, we type the formula 

=MIN INVERSE(B3:D5)

and we press the Ctrl-Shift-Enter keys simultaneously.

We observe that $R^{-1}$ appears in these cells.

3. Now, we choose the block of cells G7:G9 for the values of the current $I$. As before, we highlight them, and with the cell marker positioned in G7, we type the formula 

=MMULT(B7:D9,G3:G5)

and we press the Ctrl-Shift-Enter keys simultaneously. The values of $I$ then appear in G7:G9.

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<th>B</th>
<th>C</th>
<th>D</th>
<th>E</th>
<th>F</th>
<th>G</th>
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<td>0.138</td>
<td>0.083</td>
<td></td>
<td>22.462</td>
<td></td>
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</tr>
<tr>
<td>8</td>
<td>R$^{-1}$</td>
<td>0.138</td>
<td>0.154</td>
<td>0.092</td>
<td>I</td>
<td>13.846</td>
<td></td>
<td></td>
</tr>
<tr>
<td>9</td>
<td></td>
<td>0.083</td>
<td>0.092</td>
<td>0.122</td>
<td></td>
<td>8.3077</td>
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<tr>
<td>10</td>
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</tr>
</tbody>
</table>

*Figure C.2. Solution of Example C.17 with a spreadsheet*

**Example C.18**

For the phasor circuit of Figure C.18

*Figure C.3. Circuit for Example C.18*
the current $I_X$ can be found from the relation

$$I_X = \frac{V_1 - V_2}{R_3} \quad (C.59)$$

and the voltages $V_1$ and $V_2$ can be computed from the nodal equations

$$\frac{V_1 - 170 \angle 0^\circ}{85} + \frac{V_1 - V_2}{100} + \frac{V_1 - 0}{j200} = 0 \quad (C.60)$$

and

$$\frac{V_2 - 170 \angle 0^\circ}{-j100} + \frac{V_2 - V_1}{100} + \frac{V_2 - 0}{50} = 0 \quad (C.61)$$

Compute, and express the current $I_X$ in both rectangular and polar forms by first simplifying like terms, collecting, and then writing the above relations in matrix form as $YV = I$, where $Y = Admittance$, $V = Voltage$, and $I = Current$.

**Solution:**

The $Y$ matrix elements are the coefficients of $V_1$ and $V_2$. Simplifying and rearranging the nodal equations of (C.60) and (C.61), we get

$$(0.0218 - j0.005) V_1 - 0.01 V_2 = 2$$

$$-0.01 V_1 + (0.03 + j0.01) V_2 = j1.7 \quad (C.62)$$

Next, we write (C.62) in matrix form as

$$\begin{bmatrix} 0.0218 - j0.005 & -0.01 \\ -0.01 & 0.03 + j0.01 \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \end{bmatrix} = \begin{bmatrix} 2 \\ j1.7 \end{bmatrix} \quad (C.63)$$

where the matrices $Y$, $V$, and $I$ are as indicated.

We will use MATLAB to compute the voltages $V_1$ and $V_2$, and to do all other computations. The code is shown below.

```matlab
Y=[0.0218-0.005j -0.01; -0.01 0.03+0.01j]; I=[2; 1.7]; V=Y*I;% Define Y, I, and find V fprintf('
'); disp('V1 = '); disp(V(1)); disp('V2 = '); disp(V(2)); % Display values of V1 and V2
```

C-28
Solution of Simultaneous Equations with Matrices

\[ V_2 = 1.0490e+002 + 4.9448e+001i \]

\[ V_2 = 53.4162 + 55.3439i \]

Next, we find \( I_X \) from

\[ R_3 = 100; \quad I_X = (V(1) - V(2))/R_3 \quad \% \text{Compute the value of } I_X \]

\[ I_X = 0.5149 - 0.0590i \]

This is the rectangular form of \( I_X \). For the polar form we use

\[ \text{magIX} = \text{abs}(I_X) \quad \% \text{Compute the magnitude of } I_X \]

\[ \text{magIX} = 0.5183 \]

\[ \text{thetaIX} = \text{angle}(I_X) \times 180/\pi \quad \% \text{Compute angle theta in degrees} \]

\[ \text{thetaIX} = -6.5326 \]

Therefore, in polar form

\[ I_X = 0.518 \angle -6.53^\circ \]

Spreadsheets have limited capabilities with complex numbers, and thus we cannot use them to compute matrices that include complex numbers in their elements as in Example C.18
C.12 Exercises

For Problems 1 through 3 below, the matrices $A$, $B$, $C$, and $D$ are defined as:

\[
A = \begin{bmatrix} 1 & -1 & -4 \\ 5 & 7 & -2 \\ 3 & -5 & 6 \end{bmatrix}, \quad B = \begin{bmatrix} 5 & 9 & -3 \\ -2 & 8 & 2 \\ 7 & -4 & 6 \end{bmatrix}, \quad C = \begin{bmatrix} 4 & 6 \\ -3 & 8 \\ 5 & -2 \end{bmatrix}, \quad D = \begin{bmatrix} 1 & -2 & 3 \\ -3 & 6 & -4 \end{bmatrix}
\]

1. Perform the following computations, if possible. Verify your answers with MATLAB.
   a. $A + B$  
   b. $A + C$  
   c. $B + D$  
   d. $C + D$  
   e. $A - B$  
   f. $A - C$  
   g. $B - D$  
   h. $C - D$

2. Perform the following computations, if possible. Verify your answers with MATLAB.
   a. $A \cdot B$  
   b. $A \cdot C$  
   c. $B \cdot D$  
   d. $C \cdot D$  
   e. $B \cdot A$  
   f. $C \cdot A$  
   g. $D \cdot A$  
   h. $D \cdot C$

3. Perform the following computations, if possible. Verify your answers with MATLAB.
   a. $\det A$  
   b. $\det B$  
   c. $\det C$  
   d. $\det D$  
   e. $\det(A \cdot B)$  
   f. $\det(A \cdot C)$

4. Solve the following systems of equations using Cramer’s rule. Verify your answers with MATLAB.

   a. $x_1 - 2x_2 + x_3 = -4$  
   b. $-x_1 + 2x_2 - 3x_3 + 5x_4 = 14$  
   a. $-2x_1 + 3x_2 + x_3 = 9$  
   b. $x_1 + 3x_2 + 2x_3 - x_4 = 9$  
   a. $3x_1 + 4x_2 - 5x_3 = 0$  
   b. $3x_1 - 3x_2 + 2x_3 + 4x_4 = 19$  
   a. $4x_1 + 2x_2 + 5x_3 + x_4 = 27$  

5. Repeat Exercise 4 using the Gaussian elimination method.

6. Solve the following systems of equations using the inverse matrix method. Verify your answers with MATLAB.

   a. \[
   \begin{bmatrix} 1 & 3 & 4 \\ 3 & 1 & -2 \\ 2 & 3 & 5 \end{bmatrix} \cdot \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} = \begin{bmatrix} -3 \\ -2 \\ 0 \end{bmatrix}
   \]
   b. \[
   \begin{bmatrix} 2 & 4 & 3 & & -2 \\ 2 & -4 & 1 & & 3 \\ -1 & 3 & -4 & & 2 \\ 2 & -2 & 2 & & 1 \end{bmatrix} \cdot \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix} = \begin{bmatrix} 1 \\ 10 \\ -14 \\ 7 \end{bmatrix}
   \]
Appendix D

Constructing Semilog Plots with Microsoft Excel

This appendix contains instructions for constructing semilog plots with the Microsoft Excel spreadsheet. Semilog, short for semilogarithmic, paper is graph paper having one logarithmic and one linear scale. It is used in many scientific and engineering applications including frequency response illustrations and Bode Plots.

D.1 The Excel Spreadsheet Window

Figure D.1 shows the Excel spreadsheet workspace and identifies the different parts of the Excel window when we first start Excel.

![Figure D.1. The Excel Spreadsheet Workspace](image-url)
Figure D.2 shows that whenever a chart is selected, as shown by the visible handles around the selected chart, the Chart drop menu appears on the Menu bar and that the Chart toolbar now is visible. We can now use the Chart Objects Edit Box and Format Chart Area tools to edit our chart.

![Microsoft Excel - Print Screen](image)

**Figure D.2. The Excel Spreadsheet with Chart selected**

### D.2 Instructions for Constructing Semilog Plots

1. Start with a blank spreadsheet as shown in Figure D.1.
2. Click on ChartWizard.
3. Click on the X-Y (Scatter) Chart type under the Standard Types tab on the ChartWizard menu.
4. The Chart sub-type shows five different sub-types. Click on the upper right (the one showing two continuous curves without square points.)
5. Click on Next, Series tab, Add, Next.
6. Click on Gridlines tab and click on all square boxes under Value X-axis and Value Y-axis to place check marks on Major and Minor gridlines.

7. Click on Next, Finish, click on the Series 1 box to select it, and press the Delete key on the keyboard to delete it.

8. The plot area normally appears in gray color. To change it to white, first make sure that the chart is selected, that is, the handles (black squares) around the plot are visible. Point the mouse on the Chart Objects Edit Box tool (refer to Figure D.2), scroll down, click on the Plot Area, then click on the Format Plot Area (shown as Format Chart Area tool in Figure D.2).

9. The Area section on the Patterns tab shows several squares with different colors. Click on the white square, fifth row, right-most column, and click on OK to return to the Chart display. You will observe that the Plot Area has a white background.

10. Click anywhere near the x-axis (lowest horizontal line on the plot) and observe that the Chart Objects Edit Box now displays Value (X) axis. Click on the Format Chart Area tool which now displays Format Axis, click on the Scale tab and make the following entries:

   Minimum: 1     Maximum: 100000     Major Unit: 10     Minor Unit: 10

   Make sure that the squares to the left of these values are not checked.
   
   Click on Logarithmic scale to place a check mark, and click on OK to return to the plot.

11. Click anywhere near the y-axis (left-most vertical line on the plot) and observe that the Chart Objects Edit Box now displays Value (Y) axis. Click on the Format Chart Area tool which now displays Format Axis, click on the Scale tab and make the following entries:

   Minimum: –80     Maximum: 80     Major Unit: 20     Minor Unit: 20

   Make sure that the squares to the left of these values are not checked. Also, make sure that the Logarithmic scale is not checked. Check on OK to return to the plot.

12. You will observe that the x-axis values appear at the middle of the plot. To move them below the plot, click on Format Chart Area tool, click on the Patterns tab, click on Tick mark labels (lower right section), and click on OK to return to the plot area.

13. To expand the plot so that it will look more useful and presentable, make sure that the chart is selected (the handles are visible). This is done by clicking anywhere in the chart area. Bring the mouse close to the lower center handle until a bidirectional arrow appears and stretch downwards. Repeat with the right center handle to stretch the plot to the right. Alternately, you may bring the mouse near the lower right handle and stretch the plot diagonally.

14. You may wish to display the x-axis values in exponential (scientific) format. To do that, click anywhere near the x-axis (zero point), and observe that the Chart Objects Edit Box now displays Value (X) axis. Click on the Format Chart Area tool which now displays Format Axis, click on the Number tab and under Category click on Scientific with zero decimal places.
15. If you wish to enter title and labels for the x- and y-axes, with the chart selected, click on Chart (on the Menu bar), click on chart Options, and on the Titles tab enter the Title and the x- and y-axis labels. Remember that the Chart drop menu on the Menu bar and the Chart toolbar are hidden when the chart is deselected.

16. With the values used for this example, your semilog plot should look like the one below.
This appendix discusses magnitude and frequency scaling procedures that allow us to transform circuits that contain passive devices with unrealistic values to equivalent circuits with realistic values.

**E.1 Magnitude Scaling**

*Magnitude scaling* is the process by which the impedance of a two terminal network is changed by a factor $k_m$ which is a real positive number greater or smaller than unity.

If we increase the input impedance by a factor $k_m$, we must increase the impedance of each device of the network by the same factor. Thus, if a network consists of $R$, $L$, and $C$ devices and we wish to scale this network by this factor, the magnitude scaling process entails the following transformations where the subscript $m$ denotes magnitude scaling.

\[
\begin{align*}
R_m &\rightarrow k_m R \\
L_m &\rightarrow k_m L \\
C_m &\rightarrow \frac{C}{k_m}
\end{align*}
\]  
(E.1)

These transformations are consistent with the time-domain to frequency domain transformations

\[
\begin{align*}
R &\rightarrow R \\
L &\rightarrow j\omega L \\
C &\rightarrow \frac{1}{j\omega C}
\end{align*}
\]  
(E.2)

and the $t$-domain to $s$-domain transformations

\[
\begin{align*}
R &\rightarrow R \\
L &\rightarrow sL \\
C &\rightarrow \frac{1}{sC}
\end{align*}
\]  
(E.3)

**E.2 Frequency Scaling**

*Frequency scaling* is the process in which we change the values of the network devices so that at the new frequency the impedance of each device has the same value as at the original frequency. The fre-
Appendix E Scaling

Frequency scaling factor is denoted as \( k_f \). This factor is also a real positive number and can be greater or smaller than unity.

The resistance value is independent of the frequency. However, the complex impedance of any inductor is \( sL \), and in order to maintain the same impedance at a frequency \( k_f \) times as great, we must replace the inductor value by another which is equal to \( L/k_f \). Similarly, a capacitor with value \( C \) must be replaced with another having a capacitance value equal to \( C/k_f \). For frequency scaling then, the following transformations are necessary where the subscript \( f \) denotes magnitude scaling.

\[
\begin{align*}
R_f & \rightarrow R \\
L_f & \rightarrow \frac{L}{k_f} \\
C_f & \rightarrow \frac{C}{k_f}
\end{align*}
\]

(E.4)

A circuit can be scaled simultaneously in both magnitude and frequency using the scales values below where the subscript \( mf \) denotes simultaneous magnitude and frequency scaling.

\[
\begin{align*}
R_{mf} & \rightarrow k_m R \\
L_{mf} & \rightarrow \frac{k_m L}{k_f} \\
C_{mf} & \rightarrow \frac{1}{k_m k_f} C
\end{align*}
\]

(E.5)

Example E.1

For the network of Figure E.1 compute

\[
Z \begin{array}{c}
\rightarrow \hspace{0.5cm} R \\
\Updownarrow \hspace{0.5cm} L_2.5 \ \Omega \\
\Updownarrow \hspace{0.5cm} 0.5 \ \text{H} \\
\hspace{0.5cm} C \hspace{0.5cm} \left\{ \hspace{0.5cm} 2 \ \text{F} \right. \\
\end{array}
\]

Figure E.1. Network for Example E.1

a. the resonant frequency \( \omega_0 \).

b. the maximum impedance \( Z_{max} \).

c. the quality factor \( Q_{0p} \).

d. the bandwidth \( BW \).
e. the magnitude of the input impedance $Z$, and using MATLAB sketch it as a function of frequency.

f. Scale this circuit so that the impedance will have a maximum value of $5 \, \text{K}\Omega$ at a resonant frequency of $5 \times 10^6 \, \text{rad/s}$

**Solution:**

a. The resonant frequency of the given circuit is

$$\omega_0 = \frac{1}{\sqrt{LC}} = 1 \, \text{rad/s}$$

and thus the circuit is parallel resonant.

b. The impedance is maximum at parallel resonance. Therefore,

$$Z_{\text{max}} = 2.5 \, \Omega$$

c. The quality factor at parallel resonance is

$$Q_{0P} = \frac{\omega_0 C}{G} = \omega_0 CR = 1 \times 2 \times 2.5 = 5$$

d. The bandwidth of this circuit is

$$BW = \frac{\omega_0}{Q_{0P}} = \frac{1}{5} = 0.2$$

e. The magnitude of the input impedance versus radian frequency $\omega$ is shown in Figure E.2 and was generated with the MATLAB code below.

```matlab
w=0.01: 0.005: 5; R=2.5; G=1/R; C=2; L=0.5; Y=G+j.*(w.*C/(w.*L));
magY=abs(Y); magZ=1./magY; plot(w,magZ); grid
```
f. Using (E.1), we get

\[ k_m = \frac{R_m}{R} = \frac{5000}{2.5} = 2000 \]

Then,

\[ L_m = k_m L = 2000 \times 0.5 = 1000 \text{ H} \]

and

\[ C_m = \frac{C}{k_m} = \frac{2}{2000} = 10^{-3} \text{ F} \]

After being scaled in magnitude by the factor \( k_m = 2000 \), the network constants are as shown in Figure E.3, and the plot is shown in Figure E.4.

![Figure E.2. Plot for Example E.1](image)

Figure E.3. The network of Figure E.2 scaled by the factor \( k_m = 2000 \)

The final step is to scale the above circuit to \( 5 \times 10^6 \text{ rad/s} \). Using (E.4), we get:

\[ R_f = R = 5 \text{ k}\Omega \]

\[ L_f = L/k_f = \frac{1000}{(5 \times 10^6)} = 200 \mu\text{H} \]
The resonant frequency of the scaled circuit is

$$\omega_0 = \frac{1}{\sqrt{LC}} = \frac{1}{\sqrt{0.2 \times 10^{-2} \times 0.2 \times 10^{-9}}} = \frac{1}{0.2 \times 10^{-6}} = 5 \times 10^6 \text{ rad/s}$$

and thus the circuit is parallel resonant at this frequency.

The impedance is maximum at parallel resonance. Therefore,

$$Z_{\text{max}} = 5 \text{ K}\Omega$$
Figure E.6. Plot for Example E.1 scaled to its final form

The quality factor at parallel resonance is

\[ Q_{0P} = \frac{\omega_0 C}{G} = \omega_0 CR = 5 \times 10^6 \times 2 \times 10^{-10} \times 5 \times 10^3 = 5 \]

and the bandwidth is

\[ BW = \frac{\omega_0}{Q_{0P}} = \frac{5 \times 10^6}{5} = 10^6 \]

The values of the circuit devices could have been obtained also by direct application of (E.5), that is,

\[ R_{mf} \rightarrow k_m R \]
\[ L_{mf} \rightarrow \frac{k_m}{k_f} L \]
\[ C_{mf} \rightarrow \frac{k_m}{k_f} C \]

\[ R_{mf} = k_m R = 2000 \times 2.5 = 5 \, \text{K}\Omega \]

\[ L_{mf} = \frac{k_m}{k_f} L = \frac{2000}{5 \times 10^6} \times 0.5 = 200 \, \mu\text{H} \]

\[ C_{mf} = \frac{1}{k_m k_f} C = \frac{1}{2 \times 10^3 \times 5 \times 10^6} \times 2 = 200 \, \text{pF} \]

and these values are the same as obtained before.
Example E.2

A series RLC circuit has resistance $R = 1 \ \Omega$, inductance $L = 1 \ H$, and capacitance $C = 1 \ F$. Use scaling to compute the new values of $R$ and $L$ which will result in a circuit with the same quality factor $Q_{OS}$, resonant frequency at $500 \ Hz$ and the new value of the capacitor to be $2 \ \mu F$.

Solution:

The resonant frequency of the circuit before scaling is

$$\omega_0 = \frac{1}{\sqrt{LC}} = 1 \ \text{rad/s}$$

and we want the resonant frequency of the scaled circuit to be $500 \ Hz$ or $2\pi \times 500 = 3142 \ \text{rad/s}$.

Therefore, the frequency scaling factor must be

$$k_f = \frac{3142}{1} = 3142$$

Now, we must compute the magnitude scale factor, and since we want the capacitor value to be $2 \ \mu F$, we use (E.5), that is,

$$C_{mf} = \frac{1}{k_m k_f} C$$

or

$$k_m = \frac{C}{k_f C_{mf}} = \frac{1}{3142 \times 2 \times 10^{-6}} = 159$$

Then, the scaled values for the resistance and inductance are

$$R_m = k_m R = 159 \times 1 = 159 \ \Omega$$

and

$$L_{mf} = \frac{k_m}{k_f} L = \frac{159}{3142} \times 1 = 50.6 \ \text{mH}$$
E.3 Exercises

1. A series resonant circuit has a bandwidth of 100 rad/s, \( Q_{0s} = 20 \) and \( C = 50 \mu F \). Compute the new resonant frequency and inductance if the circuit is scaled
   a. in magnitude by a factor of 5
   b. in frequency by a factor of 5
   c. in both magnitude and frequency by factors of 5

2. A scaled parallel resonant circuit consists of \( R = 4 \, K\Omega \), \( L = 0.1 \, H \), and \( C = 0.3 \, \mu F \). Compute \( k_m \) and \( k_f \) if the original circuit had the following values before scaling.
   a. \( R = 10 \, \Omega \) and \( L = 1 \, H \)
   b. \( R = 10 \, \Omega \) and \( C = 5 \, F \)
   c. \( L = 1 \, H \) and \( C = 5 \, F \)
E.4 Solutions to the Exercises

1. a. It is given that \( BW = \omega_0/Q_{OS} = 100 \) and \( Q_{OS} = 20 \); then,
\[
\omega_0 = BW \cdot Q_{OS} = 100 \times 20 = 2000 \text{ rad/s}
\]
Since \( \omega_0^2 = 1/LC \), \( L_{OLD} = 1/\omega_0^2C = 1/(4 \times 10^6 \times 50 \times 10^{-6}) = 5 \text{ mH} \), and with \( k_m = 5 \),
\[
L_{NEW} = k_mL_{OLD} = 5 \times 5 \text{ mH} = 25 \text{ mH} \text{. Also, } C_{NEW} = C_{OLD}/k_m = 50 \times 10^{-6}/5 = 10 \text{ µF}
\]
and \( \omega_{0_{NEW}}^2 = 1/L_{NEW}C_{NEW} = 1/(25 \times 10^{-3} \times 10 \times 10^{-6}) = 10^8/25 \text{ or } \omega_{0_{NEW}} = 2000 \text{ r/s} \)

b. It is given that \( C_{OLD} = 50 \times 10^{-6} \) and from (a) \( L_{OLD} = 5 \text{ mH} \). Then, with \( k_f = 5 \),
\[
L_{NEW} = L_{OLD}/k_f = 5 \times 10^{-3}/5 = 1 \text{ mH} \text{. Also, } C_{NEW} = C_{OLD}/k_f = 50 \times 10^{-6}/5 = 10 \text{ µF}
\]
and \( \omega_{0_{NEW}}^2 = 1/L_{NEW}C_{NEW} = 1/(10^{-3} \times 10 \times 10^{-6}) = 10^8 \text{ or } \omega_{0_{NEW}} = 10000 \text{ r/s} \)

c. \( L_{OLD} = 5 \text{ mH} \) and \( C_{OLD} = 50 \times 10^{-6} \). Then, from (E.5)
\[
L_{NEW} = (k_m/k_f) \cdot L_{OLD} = (5/5) \cdot 5 \text{ mH} = 5 \text{ mH} \text{. Also from (E.5)}
\]
\[
C_{NEW} = (1/(k_mk_f)) \cdot C_{OLD} = 50 \text{ µF}/(5 \times 5) = 2 \text{ µF} \text{ and}
\]
\[
\omega_{0_{NEW}}^2 = 1/L_{NEW}C_{NEW} = 1/(5 \times 10^{-3} \times 2 \times 10^{-6}) = 10^8 \text{ or } \omega_{0_{NEW}} = 10000 \text{ r/s} \)

2. a. From (E.1), \( k_m = R_{NEW}/R_{OLD} = 4000/10 = 400 \) and from (E.5)
\[
k_f = (L_{OLD}/L_{NEW}) \cdot k_m = (1/0.1) \times 400 = 4000
\]

b. From (a) \( k_m = 400 \) and from (E.5),
\[
k_f = (1/k_m) \cdot (C_{OLD}/C_{NEW}) = (1/400) \cdot (5/0.3 \times 10^{-6}) = 41677
\]
c. From (E.5) \( k_f/k_m = L_{OLD}/L_{NEW} = 1/0.1 = 10 \) and thus \( k_f = 10k_m \) (1)
\[
\text{Also from (E.5), } k_m \cdot k_f = C_{OLD}/C_{NEW} = 5/0.3 \times 10^{-6} = 5 \times 10^6/0.3 \text{ (2)}
\]
Substitution of (1) into (2) yields \( 10k_m \cdot k_m = 5 \times 10^6/0.3 \), \( k_m^2 = 5 \times 10^6/3 \), or \( k_m = 1291 \),
and from (1) \( k_f = 1291 \times 10 = 12910 \)
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